

**UiO** • **Department of Informatics**  
University of Oslo

# **IN5450 Diffraction: from waves to signal processing**

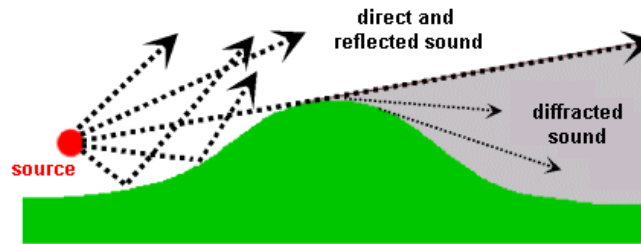
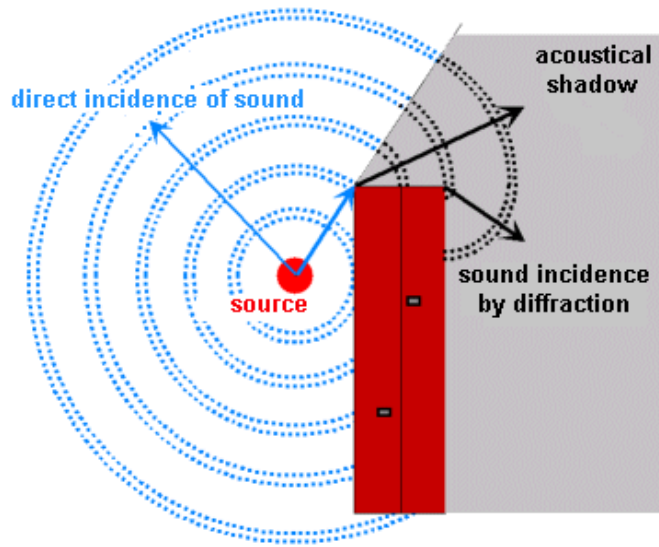
Sverre Holm



# Diffraction

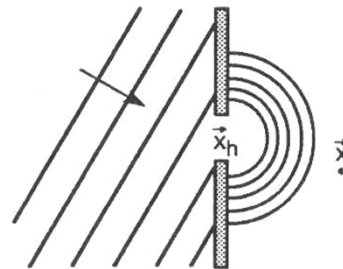
- Ray theory: Geometrical model of optics
- High-frequency – small wavelength model
- Diffraction:
  - Wavelength comparable to structure size
  - Edges of shadows are not perfectly sharp
  - Can hear around corners
- This course: only consequences of diffraction
- Johnson & Dudgeon, Array Signal Processing, Concepts and Techniques, ch. 2

# Diffraction – (spredning)



# Huygens' principle

- Christian Huygens, NL, 1629-1695
- Each point on a travelling wavefront can be considered as a secondary source of spherical radiation
- Also a model for an oscillating piston = acoustic source



**Figure 2.13** A wave is shown impinging on a hole in a planar screen. The Rayleigh-Sommerfeld diffraction formula tells us what the wavefield at the point  $\vec{x}$  is in terms of the wavefield at the aperture.

## Mathematical physics of diffraction

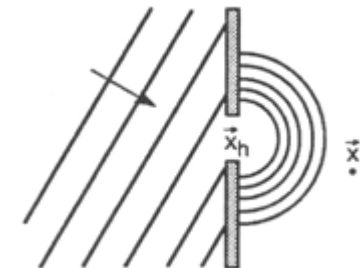
- Augustin Jean Fresnel (F) 1788 – 1827
- Gustav Robert Kirchhoff (D) 1824 – 1887
- Lord Rayleigh, John William Strutt (GB) 1842 – 1919, Nobel prize physics, 1904.
- Arnold Johannes Wilhelm Sommerfeld (D) 1868 – 1951
- Joseph von Fraunhofer (D) 1787 - 1826

# Diffraction: deviation from geometrical model

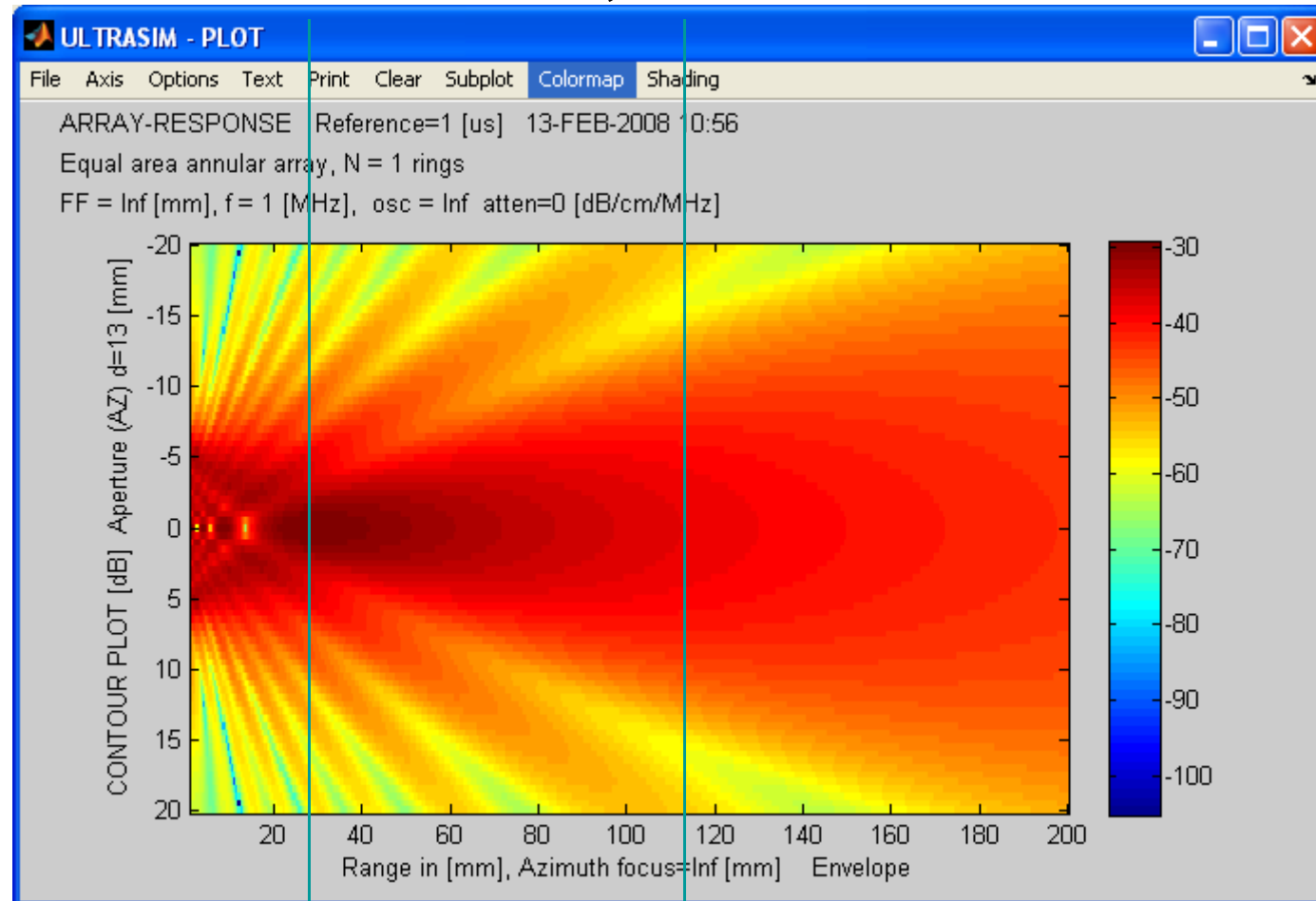
- Rayleigh-Sommerfeld diffraction formula from a hole with aperture A:

$$s(\vec{x}) = \frac{1}{j\lambda} \int \int_A s(\vec{x}_h) \frac{\exp\{jkr\}}{r} \cos\theta dA$$

- Wave at  $x$ : superposition of fields from the hole, due to linearity of wave equation
- Weighted by spherical spreading function  $e^{jkr}/r$
- Also weighted by  $1/\lambda$
- Obliquity factor  $\cos\theta$
- Phase shift of  $\pi/2$  due to  $1/j$



# 1 MHz 13 mm, unfocused xdcr



Olympus-Panametrics  
A303S  
(in our lab)

Simulation:  
[http://www.ifi.uio.no/  
~ultrasim](http://www.ifi.uio.no/~ultrasim)

$D^2/4\lambda$

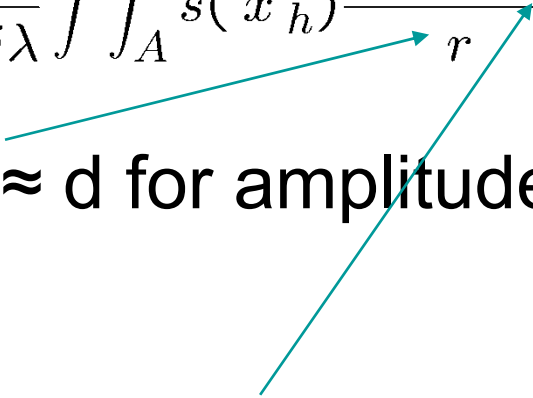
$D^2/\lambda=113 \text{ mm}$

## Two approximations

- Fresnel, nearfield, (but not quite near)
- Fraunhofer, farfield
- Leads to
  - important estimates for nearfield – farfield transition distance
  - Fourier relationship between aperture excitation and field



# 1. Fresnel approximation

$$s(\vec{x}) = \frac{1}{j\lambda} \iint_A s(\vec{x}_h) \frac{\exp\{jkr\}}{r} \cos\theta dA$$


- $\cos\theta \approx 1$ ,  $r \approx d$  for amplitude (not for phase!)
- Phase:
  - spherical surfaces  $\approx$  quadratic
  - parabolic approximation

# Fresnel derivation

- Point in the hole  $(\tilde{x}, \tilde{y}, 0)$ , in observation plane  $(\mathbf{x}, \mathbf{y}, \mathbf{z}=\mathbf{d})$
- Distance:  $r = [(x - \tilde{x})^2 + (y - \tilde{y})^2 + d^2]^{1/2}$

$$r = d \left[ 1 + \frac{(x - \tilde{x})^2 + (y - \tilde{y})^2}{d^2} \right]^{1/2}$$

- Approximate  $(1+x)^{1/2} \approx 1+x/2$ , i.e. small  $x/d \Leftrightarrow$  small angles

$$r \approx d + \frac{(x - \tilde{x})^2 + (y - \tilde{y})^2}{2d}$$

- Use the above expression for the phase and  $r=d$  for the amplitude in the Rayleigh-Sommerfeld integral

# Fresnel approximation

$$s(x, y) = \frac{\exp\{jkd\}}{j\lambda d} \int \int_A s(\tilde{x}, \tilde{y}) \exp\left\{\frac{jk[(x - \tilde{x})^2 + (y - \tilde{y})^2]}{2d}\right\} d\tilde{x}d\tilde{y}$$

- Nearfield approximation & within  $\approx 15^\circ$  of z-axis
- Also called paraxial approximation
- **2D convolution** between field in hole and  $h(x,y)$ :

$$h(x, y) = \frac{\exp\{jkd\}}{j\lambda d} \exp\left\{\frac{jk(x^2 + y^2)}{2d}\right\}$$

- This is a quadratic phase function = the phase shift that a secondary wave encounters during propagation

## 2. Fraunhofer approximation

- Expand phase term of Fresnel approximation

$$s(x, y) = \frac{\exp\{jkd\}}{j\lambda d} \iint_A s(\tilde{x}, \tilde{y}) \exp\left\{\frac{jk[(x - \tilde{x})^2 + (y - \tilde{y})^2]}{2d}\right\} d\tilde{x}d\tilde{y}$$

- and neglect quadratic phase term variation over hole

$$(x - \tilde{x})^2 + (y - \tilde{y})^2 = x^2 + y^2 - 2x\tilde{x} - 2y\tilde{y} + \tilde{x}^2 + \tilde{y}^2 \approx x^2 + y^2 - 2x\tilde{x} - 2y\tilde{y}$$

- If  $D = \max$  linear dimension of hole, this is equivalent to assuming ( $d = \text{dist. from source}$ ):

$$\frac{\tilde{x}^2}{2d} \leq \frac{(D/2)^2}{2d} \ll \lambda/2 \Rightarrow d \gg \frac{D^2}{4\lambda} \quad \text{Fresnel limit}$$

## Fraunhofer approximation

$$s(x, y) = \frac{\exp\{jkd\}}{j\lambda d} \exp\left\{\frac{jk(x^2 + y^2)}{2d}\right\} \iint_A s(\tilde{x}, \tilde{y}) \exp\left\{\frac{jk(x\tilde{x} + y\tilde{y})}{d}\right\} d\tilde{x}d\tilde{y}$$

- Far-field approximation: valid far away from hole
- $s(x, y) =$  **2D Fourier transform** of field in hole

## Linking physics & signal processing

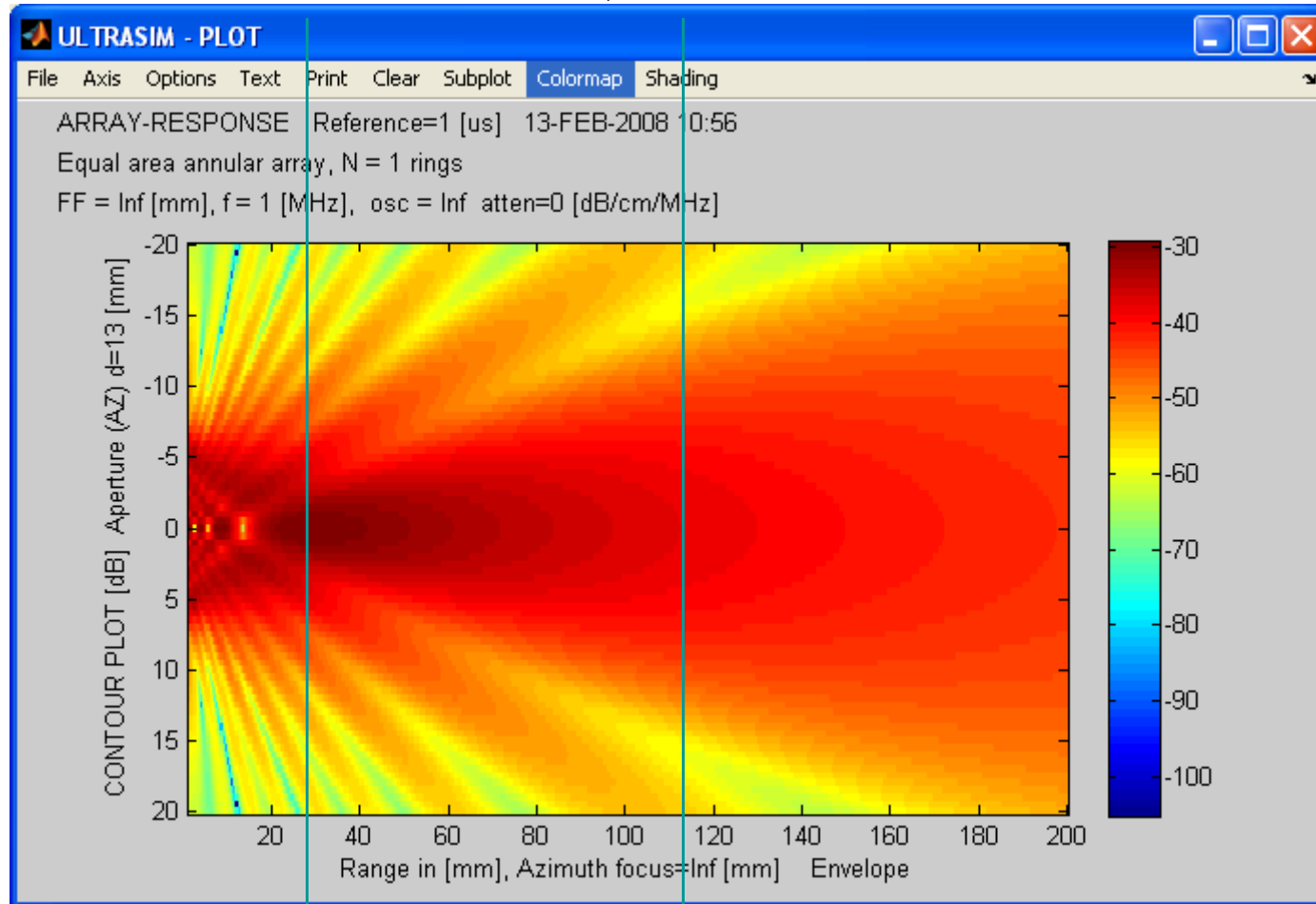
- Fourier transform relationship is a very important result
- *Link* between the physics and the signal processing!
- Basis for simplified expressions like angular resolution  $\approx \lambda/D$  etc
- Small hole leads to wide beam and vice versa just like a short time-function has a wide spectrum etc

## Nearfield-farfield limit

Not a clear transition, several limits are used, in increasing size:

- $d_F = D^2/4\lambda$  : Fresnel limit
- $d = \pi r^2/\lambda = \pi D^2/4\lambda$  : Diffraction limit
- $d = D^2/\lambda$  : max path length difference  $\lambda/8$
- $d_R = 2D^2/\lambda$  : Rayleigh dist:  $\Delta \text{ path} = \lambda/16$
- Proportional to  $D^2/\lambda$ , multiplied by 0.25, 0.79, 1, or 2

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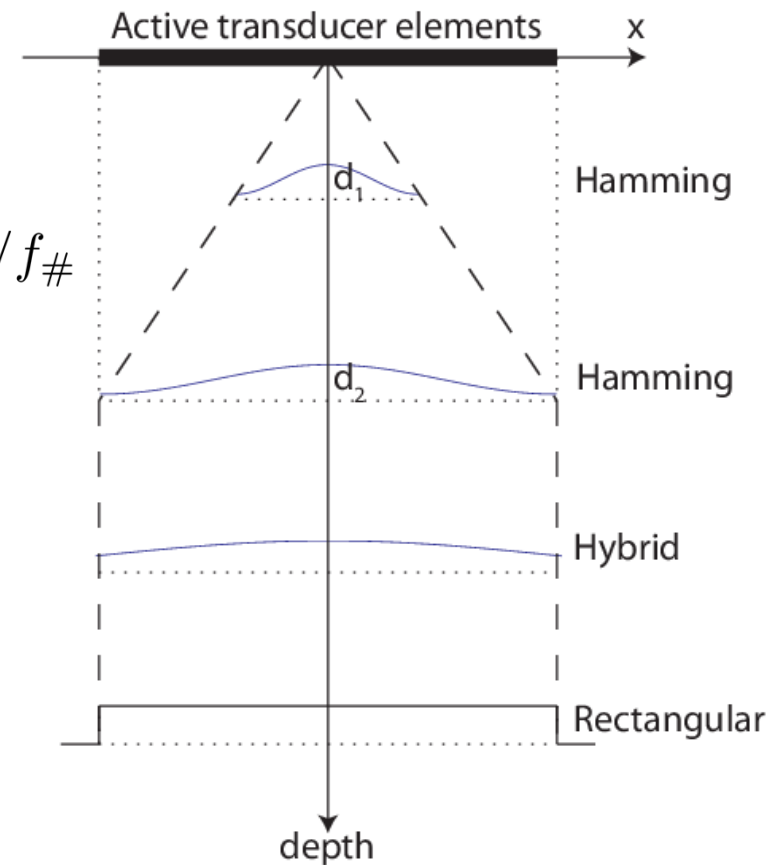


# Expanding aperture

- f-number = distance/aperture:

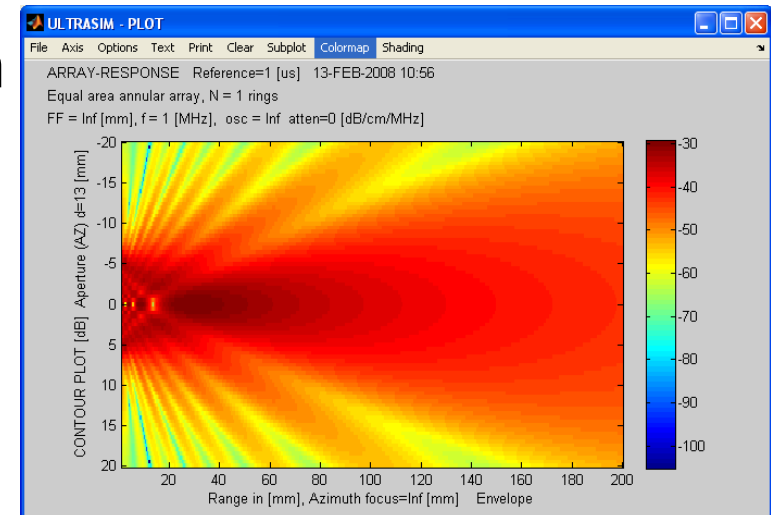
$$f_{\#} = d/D \Leftrightarrow \tan \theta/2 = d/(D/2) = 0.5/f_{\#}$$

- Fresnel theory:  
<15 deg  $\Leftrightarrow$   $f_{\#} > 3.8$
- Rule-of-thumb  $f_{\#} > 1-2$   
 $\Leftrightarrow$  53...28 deg
- Ultrasound imaging



## Field vs distance

1. Near: Collimated beam
  - Beamwidth  $\propto$  source width
  - Laser, x-ray
2. Mid: Near/far transition
3. Far: Farfield
  - Fourier relationship
  - Constant angle
  - Sonar, radar, ...



Near

Far

## Example of collimated beam: Laser pointer

- Aperture:  $D = 1\text{mm}(?)$ , could be 2
- Wavelength (red):  $\lambda = 650\text{ nm}$
- Near/far-transition:
  - $D^2/4\lambda = (1 \dots 2\text{e-}3)^2/650\text{e-}9 \heartsuit 1.5 \dots 6\text{ m}$
  - $D^2/\lambda \heartsuit 6 \dots 24\text{ m}$

# Array Processing Implications

- Diffraction means that opaque objects located between the source and the array can induce complicated wavefields
  - Scattering theory:
    - Acoustics: Schools of fish
    - Electromagnetics: rain drops
    - Complicated, but important to understand

## Norsk terminologi

- Bølgeligningen
- Planbølger, sfæriske bølger
- Propagerende bølger, bølgetall
- Sinking/sakking:
- Dispersjon  $\vec{\alpha}$
- Attenuasjon eller demping
- Refraksjon
- Ikke-linearitet
- Diffraksjon; nærfelt, fjernfelt
- Gruppeantenne (= array)

Kilde: Bl.a. J. M. Hovem: ``Marin akustikk'', NTNU, 1999