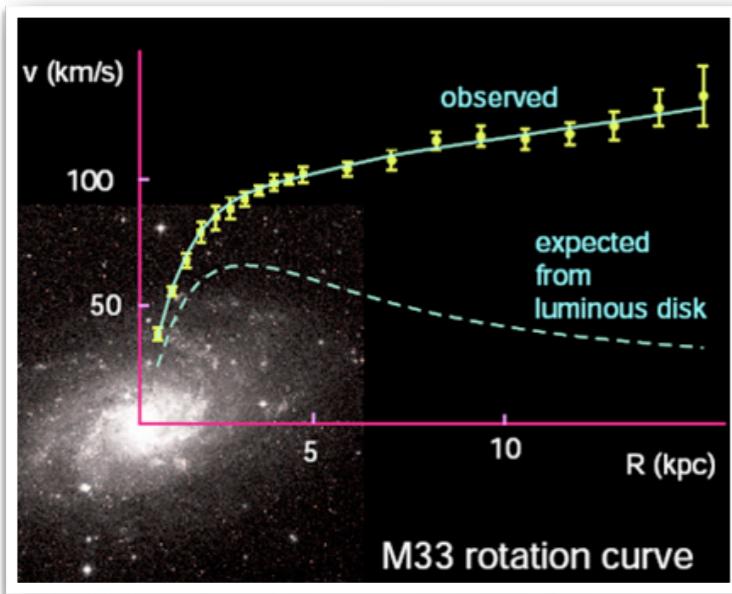
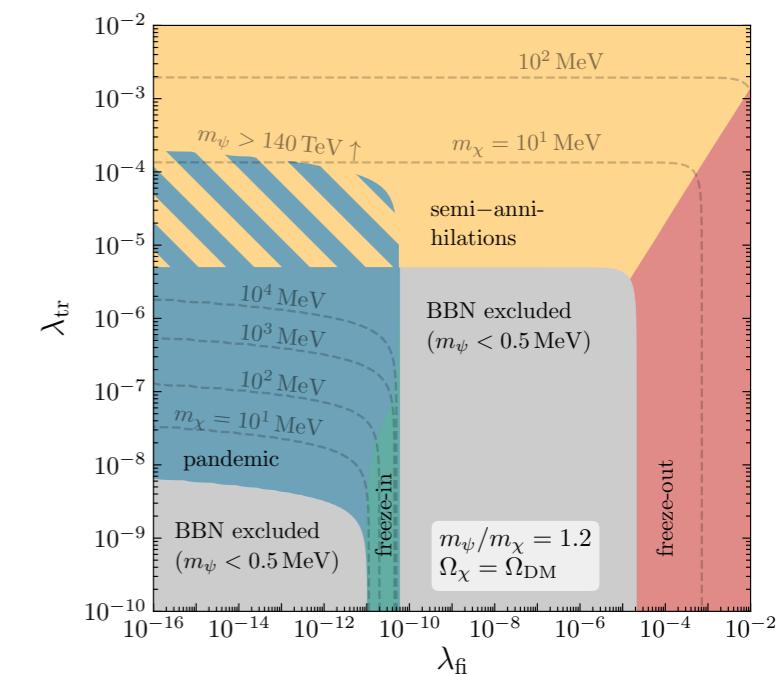
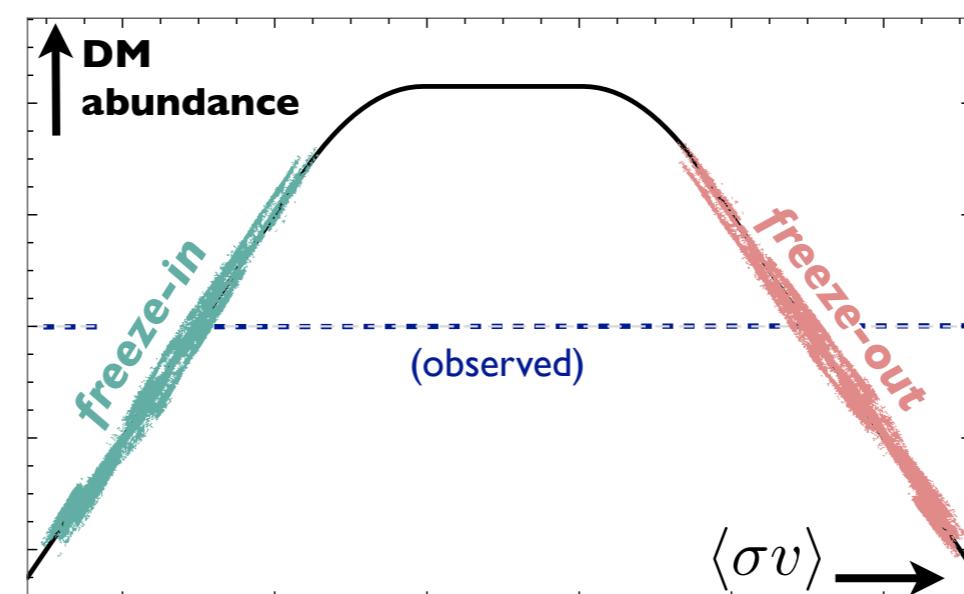


Thermal production of dark matter in the early universe

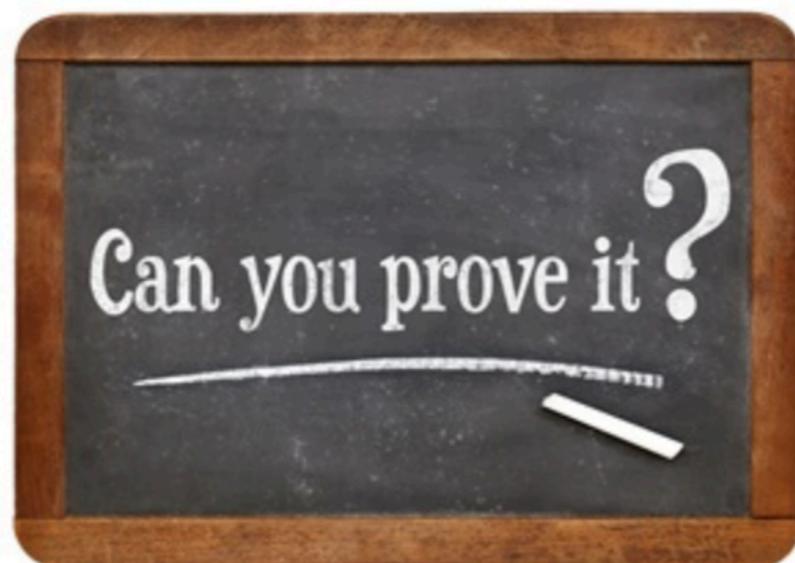
Torsten Bringmann



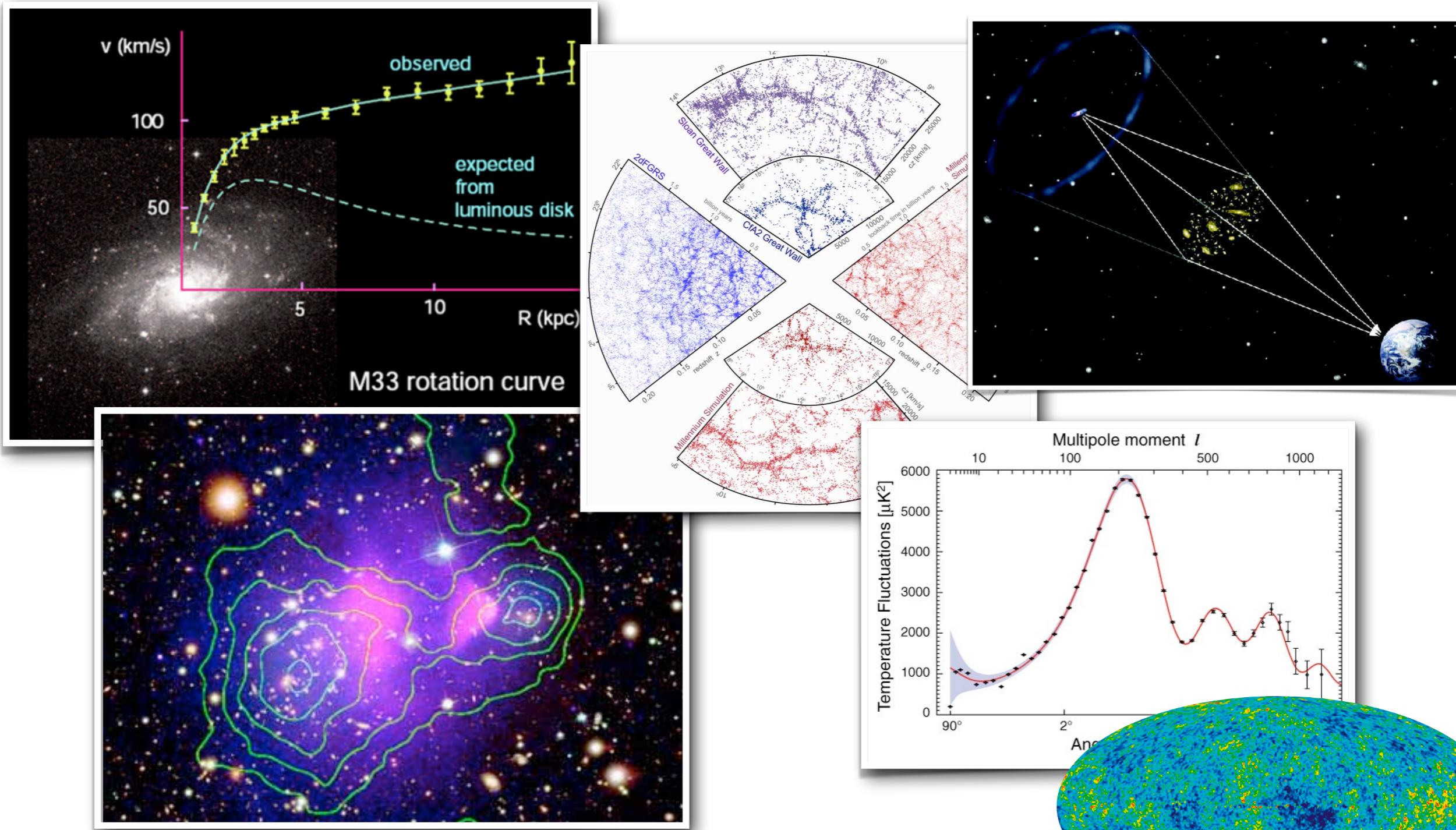
evidence



Part I

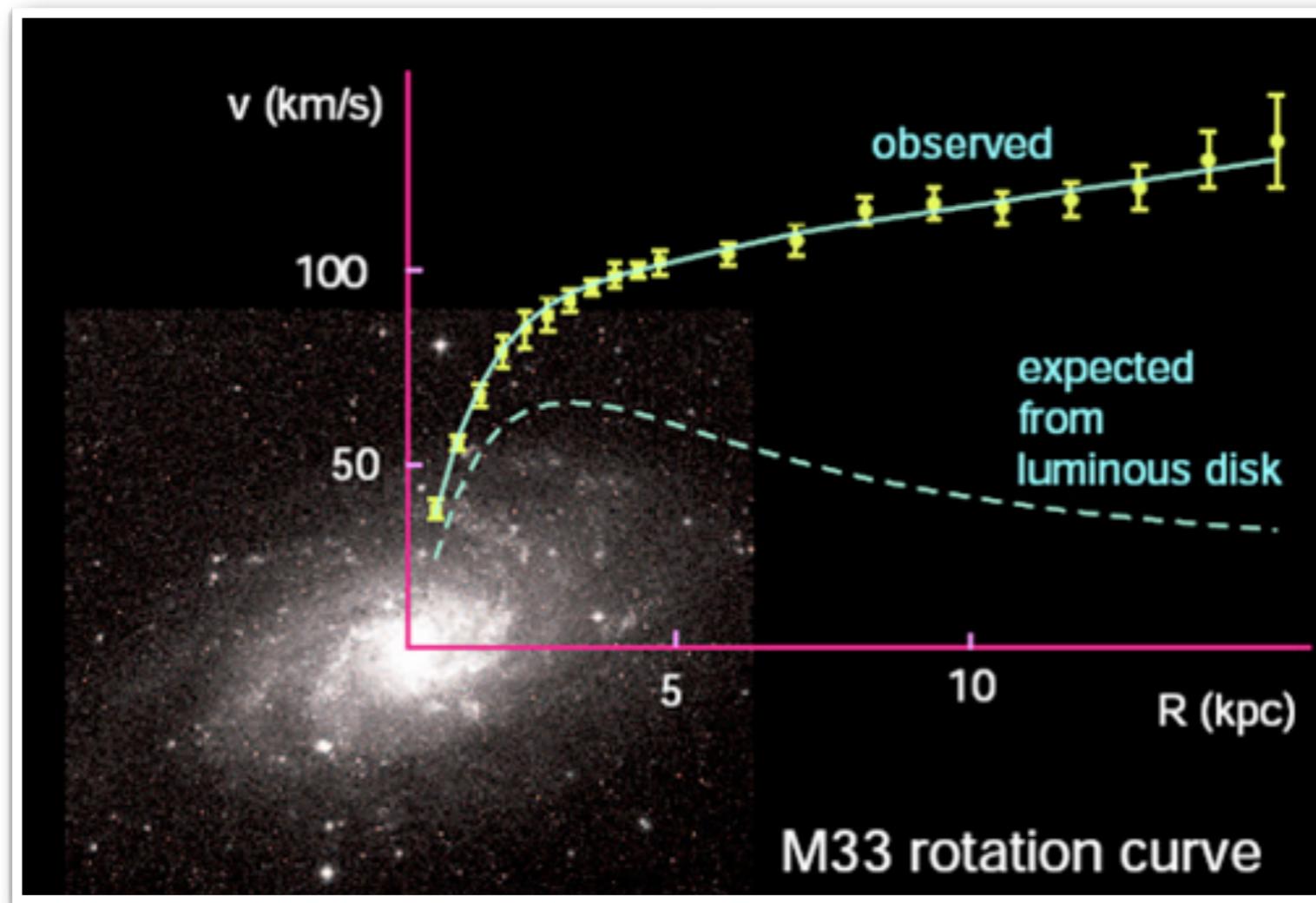


Dark matter all around



→ overwhelming evidence on all scales!

Galactic scales



Newton:

$$G_N m_{\odot} \frac{M(r < R)}{R^2} = m_{\odot} \frac{v^2}{R}$$

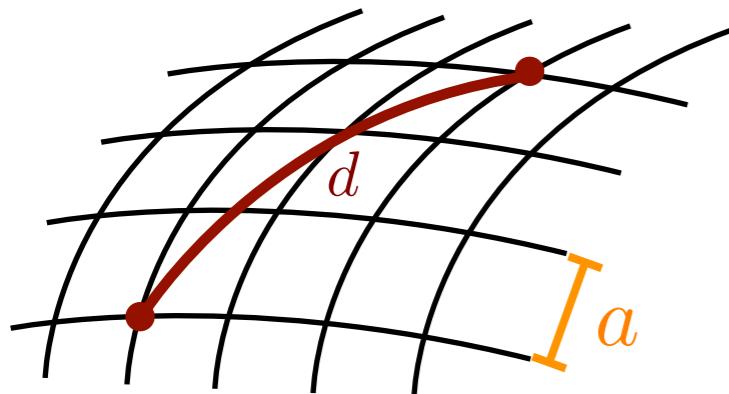
‘missing’ mass



- Rotation curves no longer main argument for existence of dark matter !!!
 - observed rotation curves rather diverse
 - other potential explanations (for this particular discrepancy)

Cosmological scales

Image credit: Jimmy Harris



- Friedman equations fix $a(t)$:

$$H^2 \equiv \left(\frac{\dot{a}}{a} \right)^2 = \frac{8\pi G}{3} \rho$$
$$\ddot{a} = -\frac{4\pi G}{3} (\rho + 3p)$$

→ **background evolution**

homogeneity + isotropy add tiny initial **perturbations** to background evolution

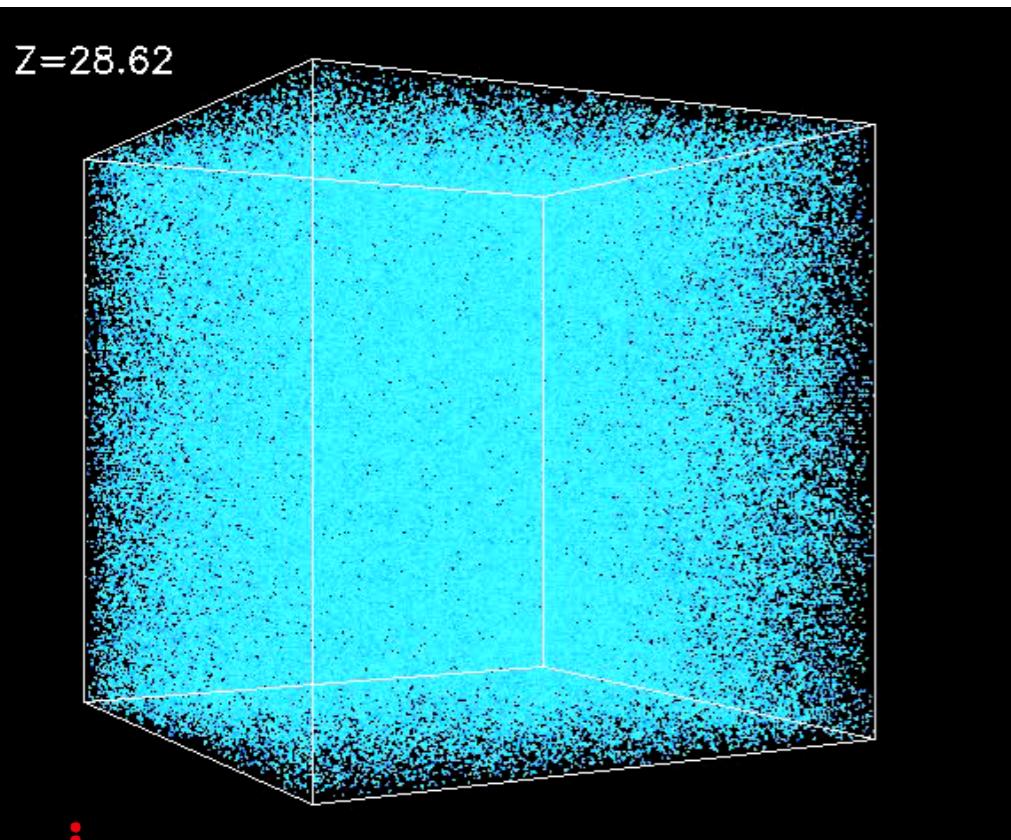
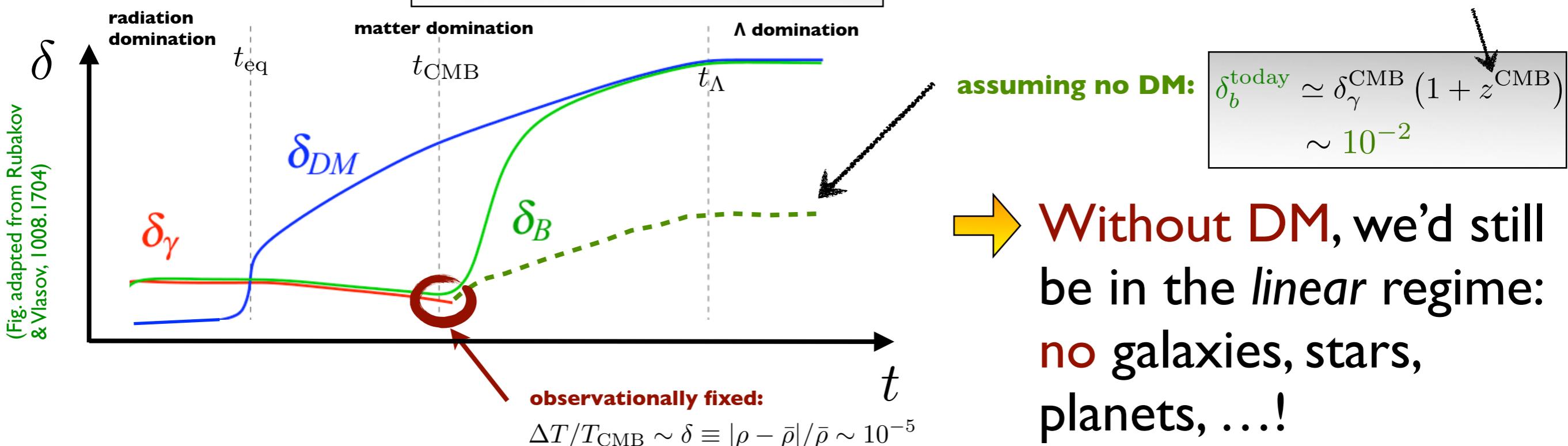
- Gravitational clustering (in linear regime) and **collapse** (non-linear)

→ **Strong impact of dark matter**

Cosmological scales

- Linear gravity

$$\rho(t, \mathbf{x}) = \bar{\rho}(t)[1 + \delta(t, \mathbf{x})]$$



- Non-linear evolution
 - Need simulations
 - Dark matter required to reach ~perfect agreement with observations (at large scales)

From evidence to precision

- DM is a **crucial ingredient** of cosmological SM!

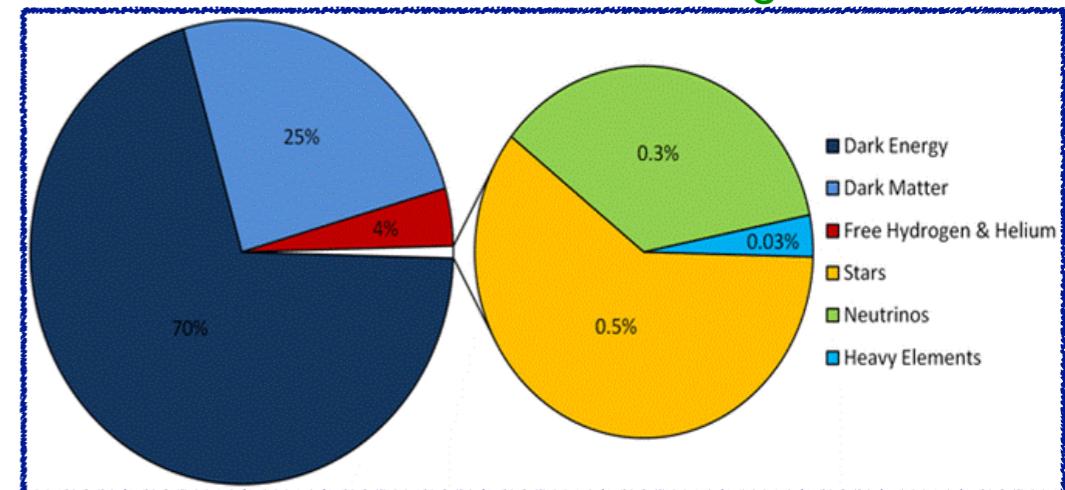
- **constant** co-moving energy **density**
- **only gravitational** interactions
- cold + dissipation-less

→ $\Omega_{\text{CDM}} h^2 = 0.1188 \pm 0.0010$

Ade+ [Planck Coll.], A&A '16

Percent-level measurements of a single parameter!

Image credit: KIAS



- DM **conversion** into (in)visible energy?

- E.g. decays, late-time annihilation, coalescing PBHs, ...

→ Ω_{CDM} decrease of **up to 10%** possible during matter domination!

(*model-independent*; NB: much more allowed during RD)

TB, Kahlhoefer, Schmidt-Hoberg & Walia, PRD '18

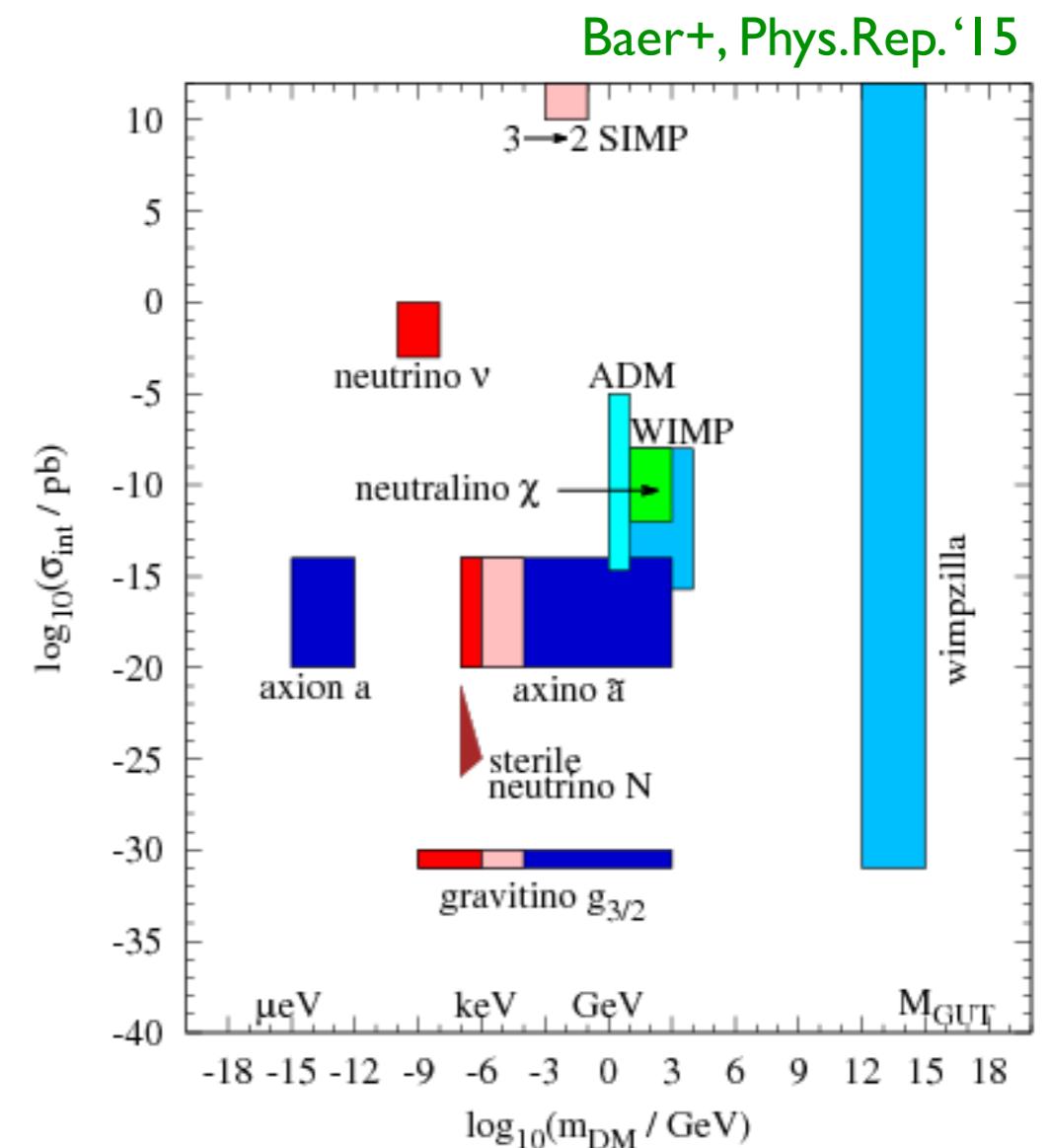
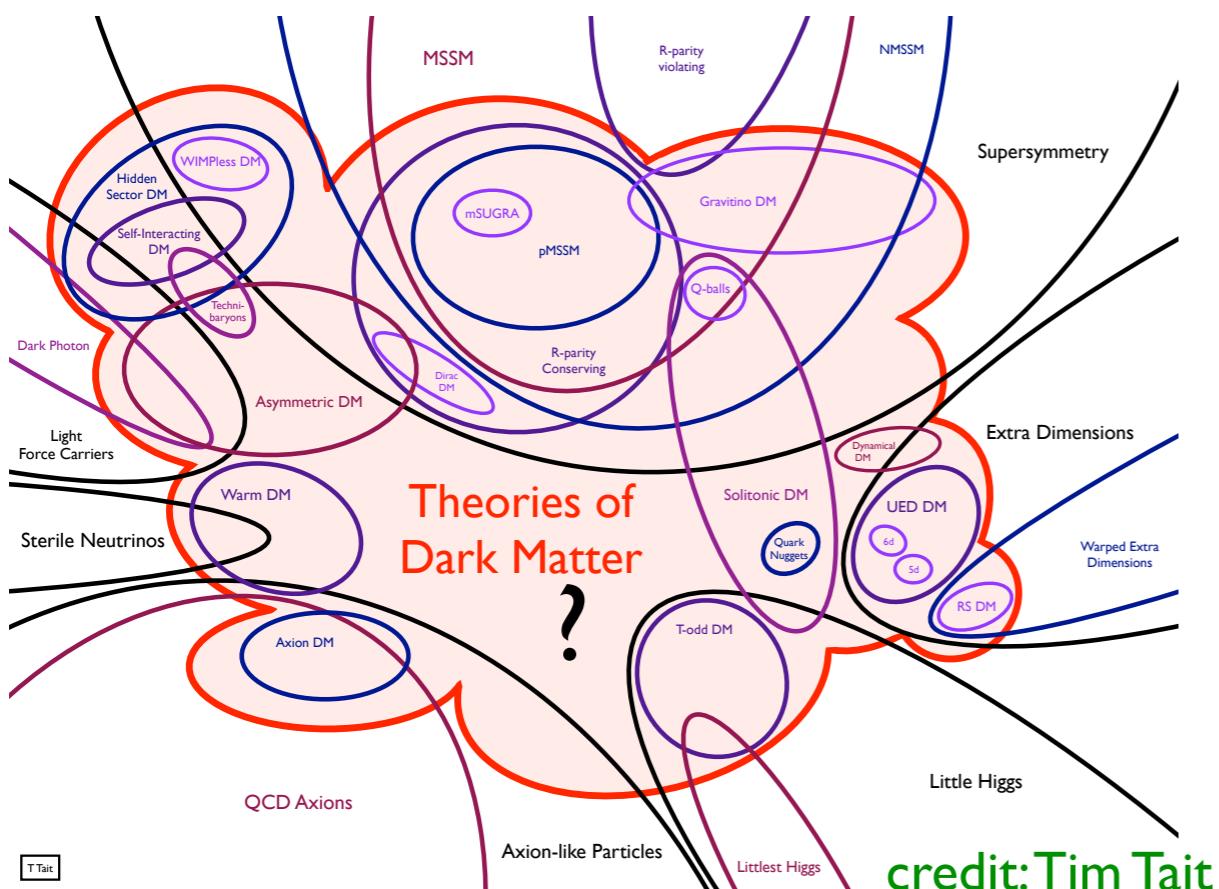
- Q: Can't we explain *all* this also by **modified gravity**?

- A: **No!** [though definitely yes for *selected* observations]



Candidates

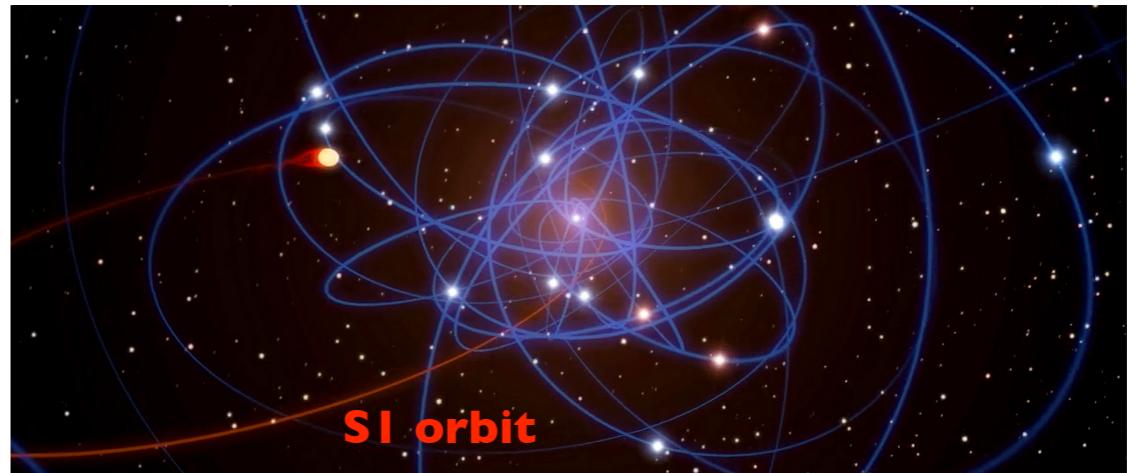
- Existence of (particle) DM = evidence for BSM physics!
- + rather good handle on what it is **not**
- Unfortunately, this still leaves too many options...



Black holes (I)

- Wouldn't (super-)**solar** mass **black holes** be an “obvious” / “conventional” candidate?

[#  2017, 2020]



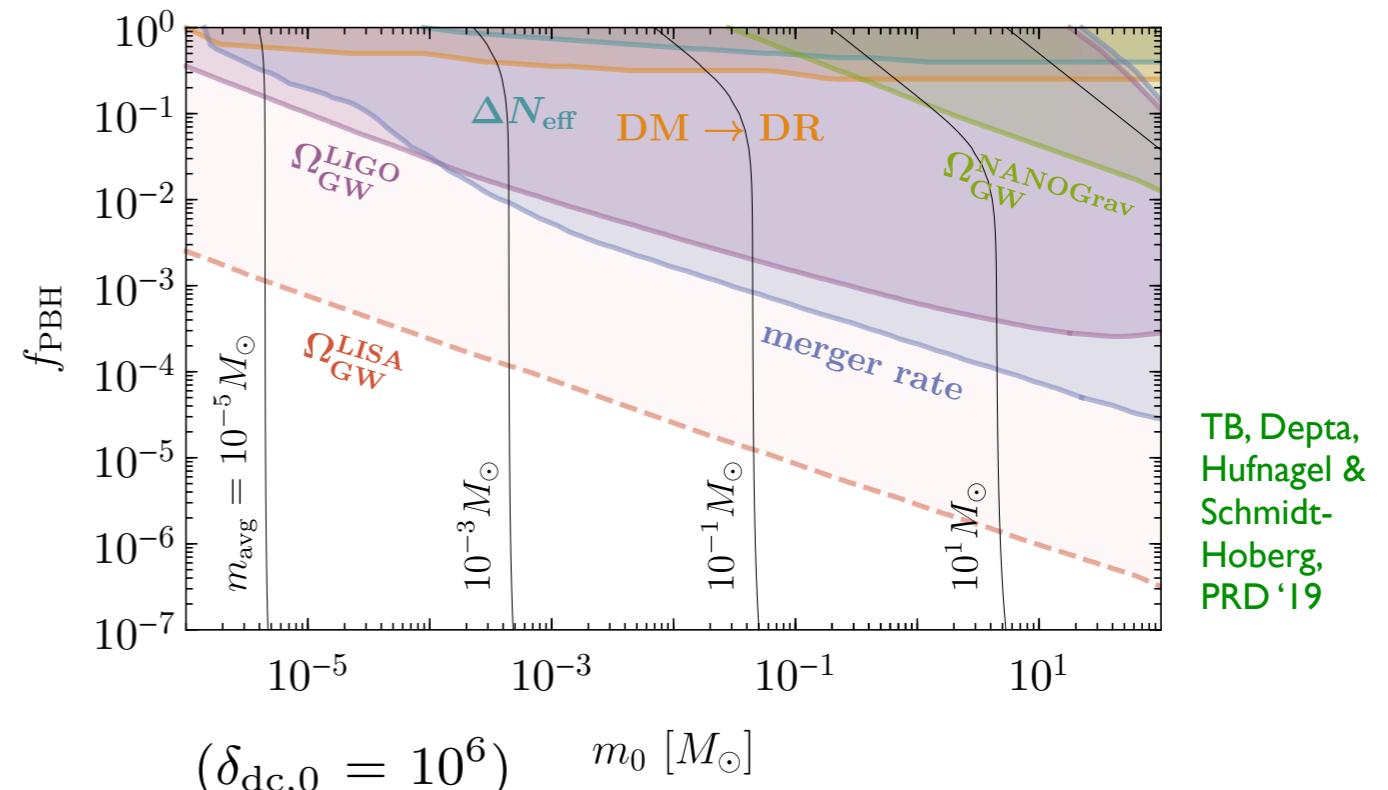
- Strongly constrained by micro-lensing and CMB!
→ Black holes can only be a **sub-dominant** DM component

overview:
Carr, Kohri, Sendouda & Yokoyama, 2002.12778

- Conclusion does not change for large black hole **clustering**...

c.f. García-Bellido & Clesse, PDU '18

...

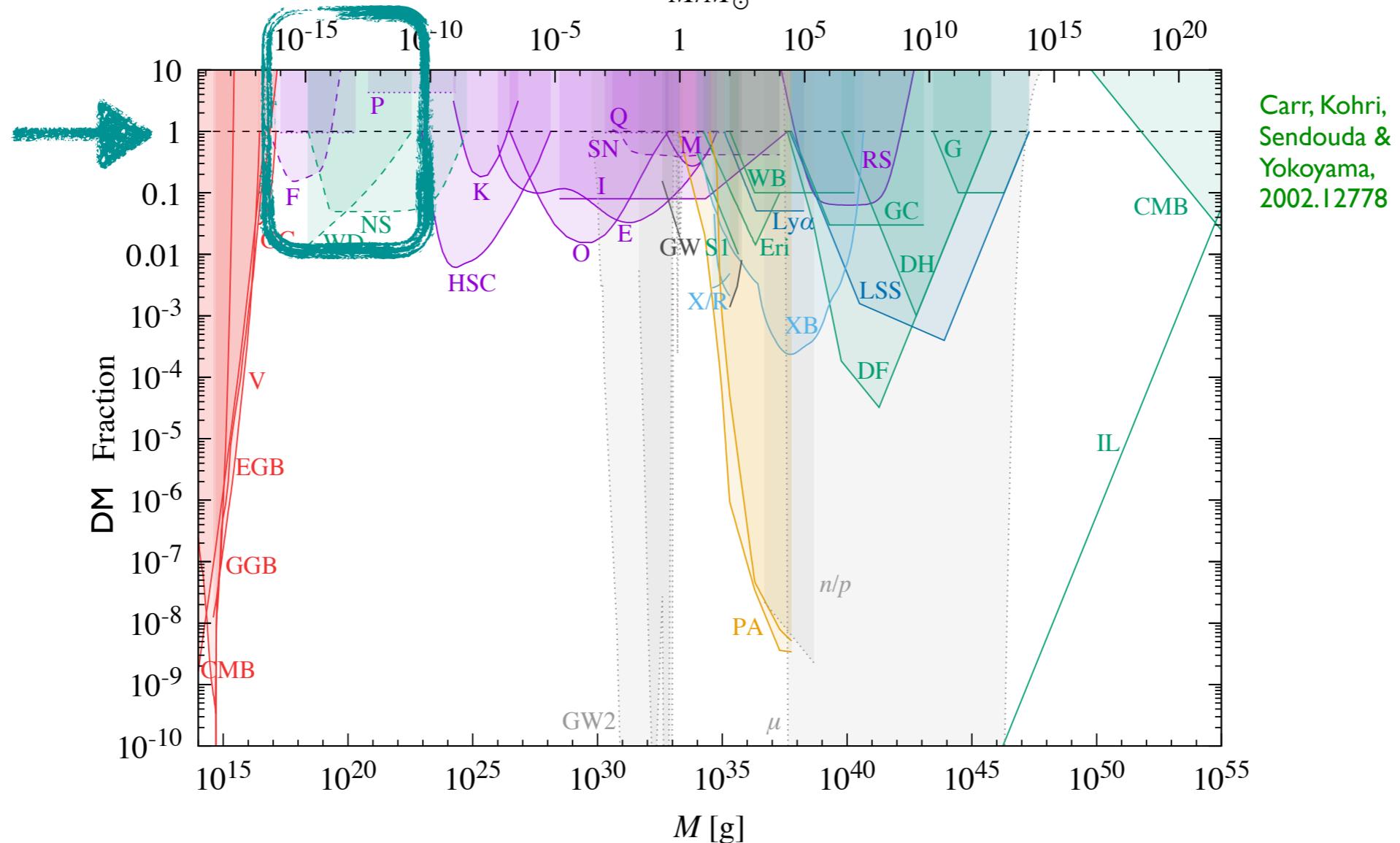


Black holes (II)

- Primordial black holes can be much smaller

→ open mass window for $f_{\text{PBH}} \sim 1$!

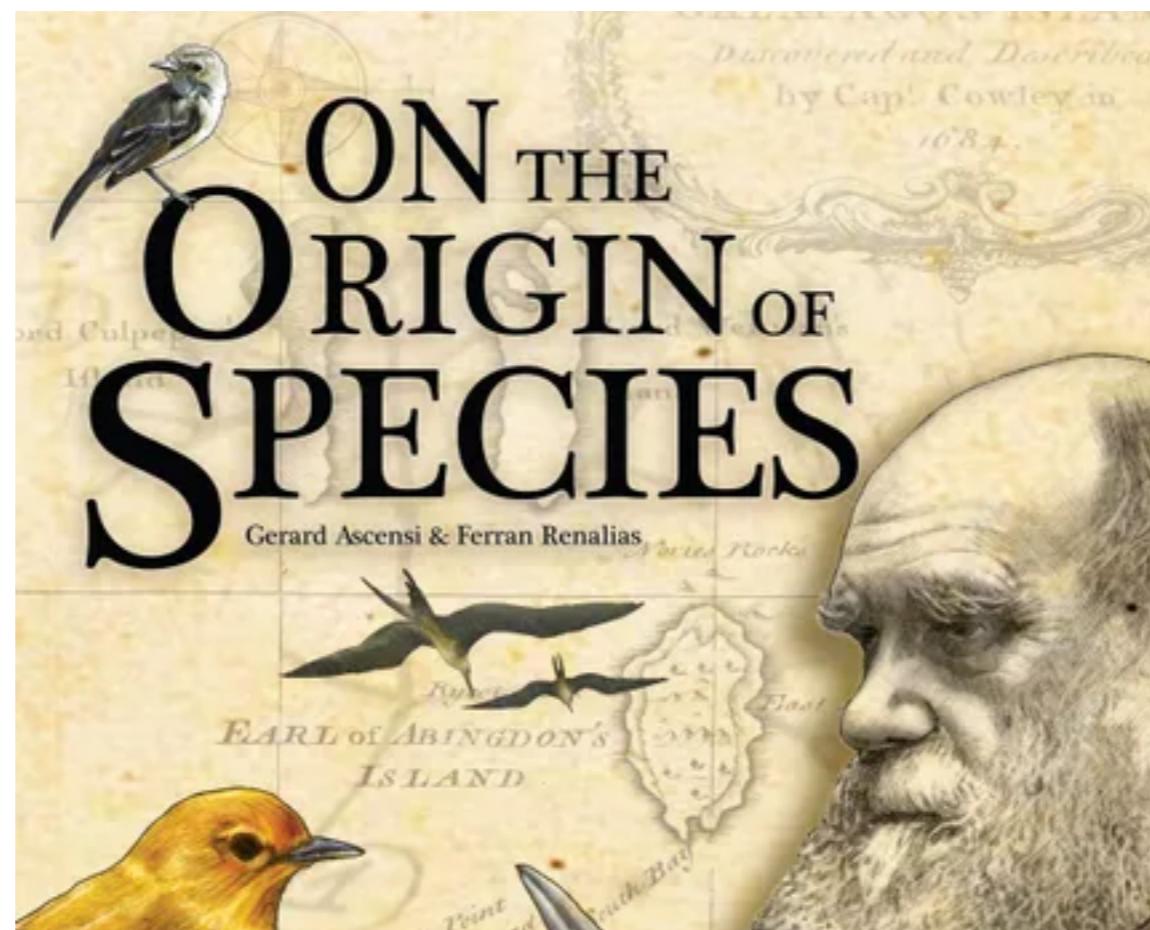
Katz, Kopp & Sibiryakov, JCAP '18
 M/M_\odot



- But this would also **not** be “SM physics” ... !
- formation (+ requirement of $f_{\text{PBH}} \sim 1$) requires BSM physics

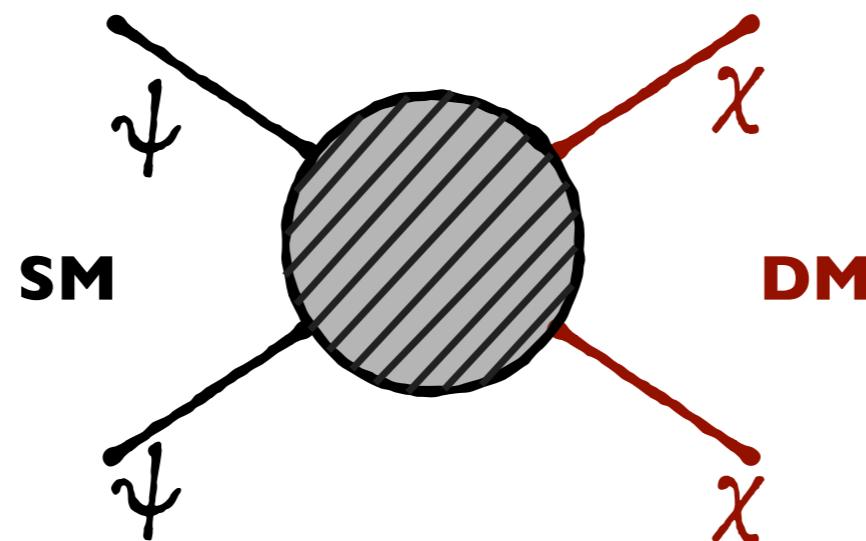


Part II



The origin of dark matter

- Existence of (particle) DM = **evidence** for BSM physics
- Any convincing model for dark matter must include a **production mechanism** that can explain the observed abundance!
- Let's postulate some interaction with the primordial heat bath:
 - [Z_2 symmetry not strictly necessary, but automatically guarantees stability of DM]



Boltzmann equation

- # Evolution of DM phase-space density:

$$L[f_\chi] = C[f_\chi]$$

 “ $= \frac{df_\chi}{dt}$ ” = $E (\partial_t - H \mathbf{p} \cdot \nabla_{\mathbf{p}}) f_\chi$ **for FRW metric**

- # Collision term:

$$C[f_\chi] = \frac{1}{2g_\chi} \int \frac{d^3\tilde{p}}{(2\pi)^3 2\tilde{E}} \int \frac{d^3k}{(2\pi)^3 2\omega} \int \frac{d^3\tilde{k}}{(2\pi)^3 2\tilde{\omega}} (2\pi)^4 \delta^{(4)}(\tilde{p} + p - \tilde{k} - k) \\ \times \left[|\mathcal{M}|_{\bar{\chi}\chi \leftarrow \bar{f}f}^2 f_\psi(\omega) f_\psi(\tilde{\omega}) - |\mathcal{M}|_{\bar{\chi}\chi \rightarrow \bar{f}f}^2 f_\chi(E) f_\chi(\tilde{E}) \{1 \pm f_\psi(\omega)\} \{1 \pm f_\psi(\tilde{\omega})\} \right]$$

- **Detailed balance:** ‘production’ = ‘annihilation’ *in equilibrium*

~~~ allows to re-write everything in terms of a *would-be equilibrium* population of DM

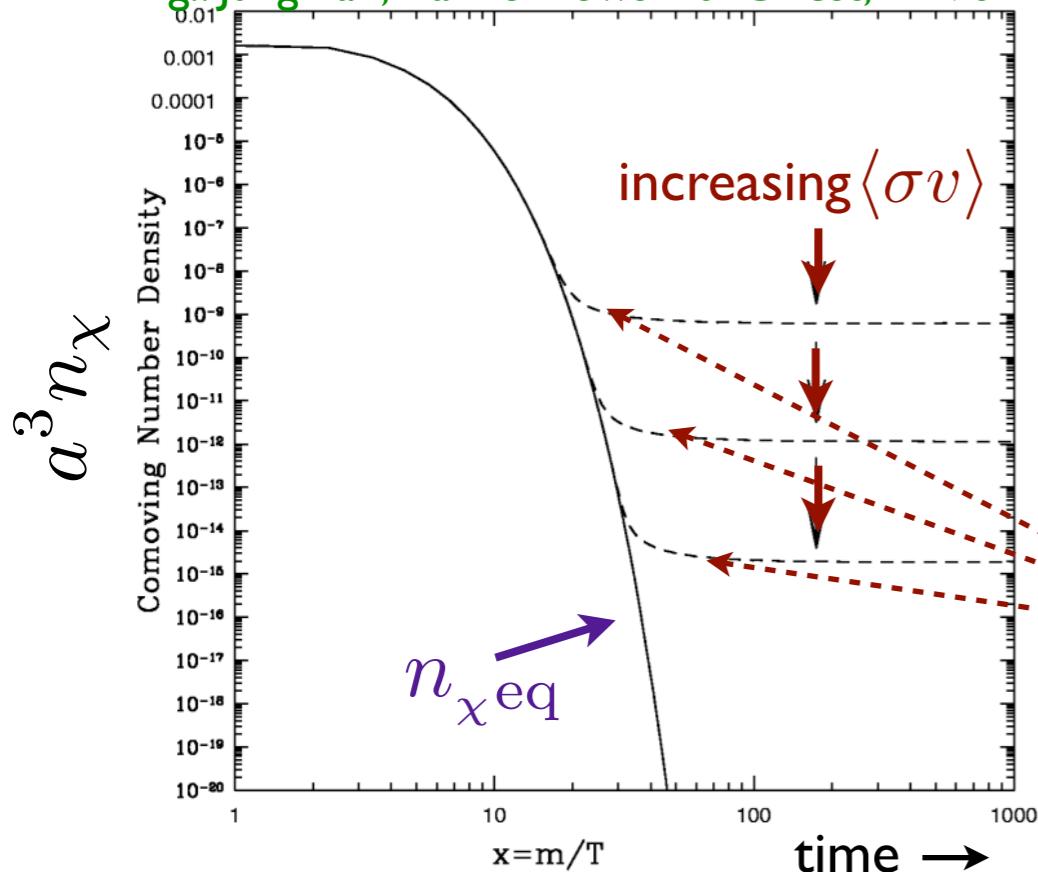
- # Pauli suppression / Bose enhancement

↪ not relevant for production of non-relativistic **DM** [energy conservation!]

# Weakly Interacting Massive Particles

- well-motivated from particle physics [SUSY, EDs, ...]
- thermal production in early universe:

Fig.: Jungman, Kamionkowski & Griest, PR'96



$$\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle \sigma v \rangle (n_\chi^2 - n_{\chi}^{\text{eq}})$$

$\langle \sigma v \rangle$ :  $\chi\chi \rightarrow \text{SM SM}$   
(thermal average)



“Freeze-out” when annihilation  
rate falls behind expansion rate

→ Relic density (today):  $\Omega_\chi h^2 \sim \frac{3 \cdot 10^{-27} \text{cm}^3/\text{s}}{\langle \sigma v \rangle} \sim \mathcal{O}(0.1)$

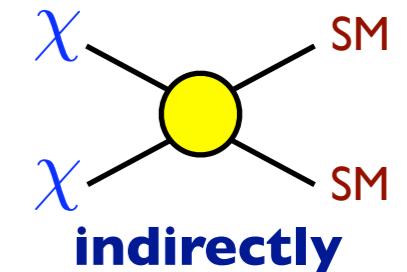
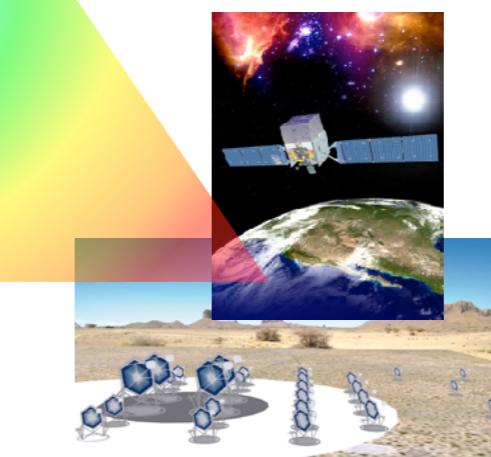
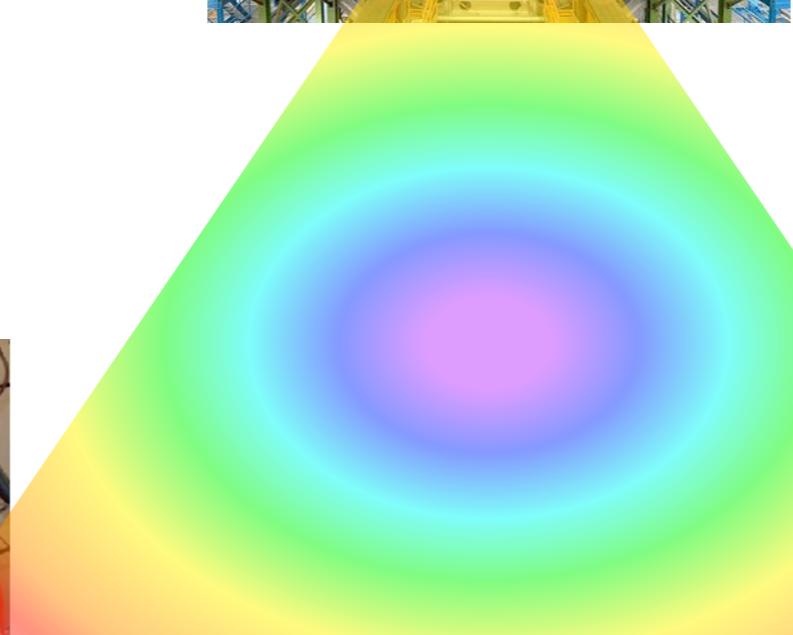
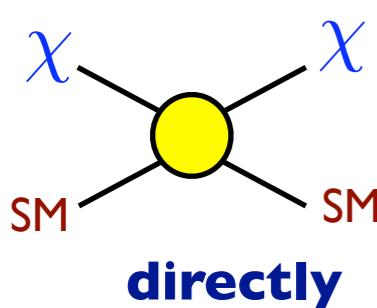
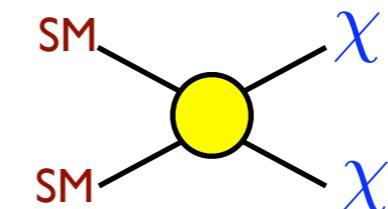
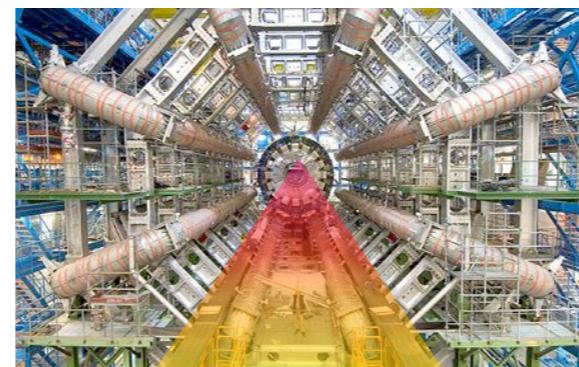
= a ‘miracle’ ?

for weak-scale  
interactions!

# WIMP DM is a *predictive* scenario

- Same interaction can be probed **today**, in multiple ways:

**at colliders**

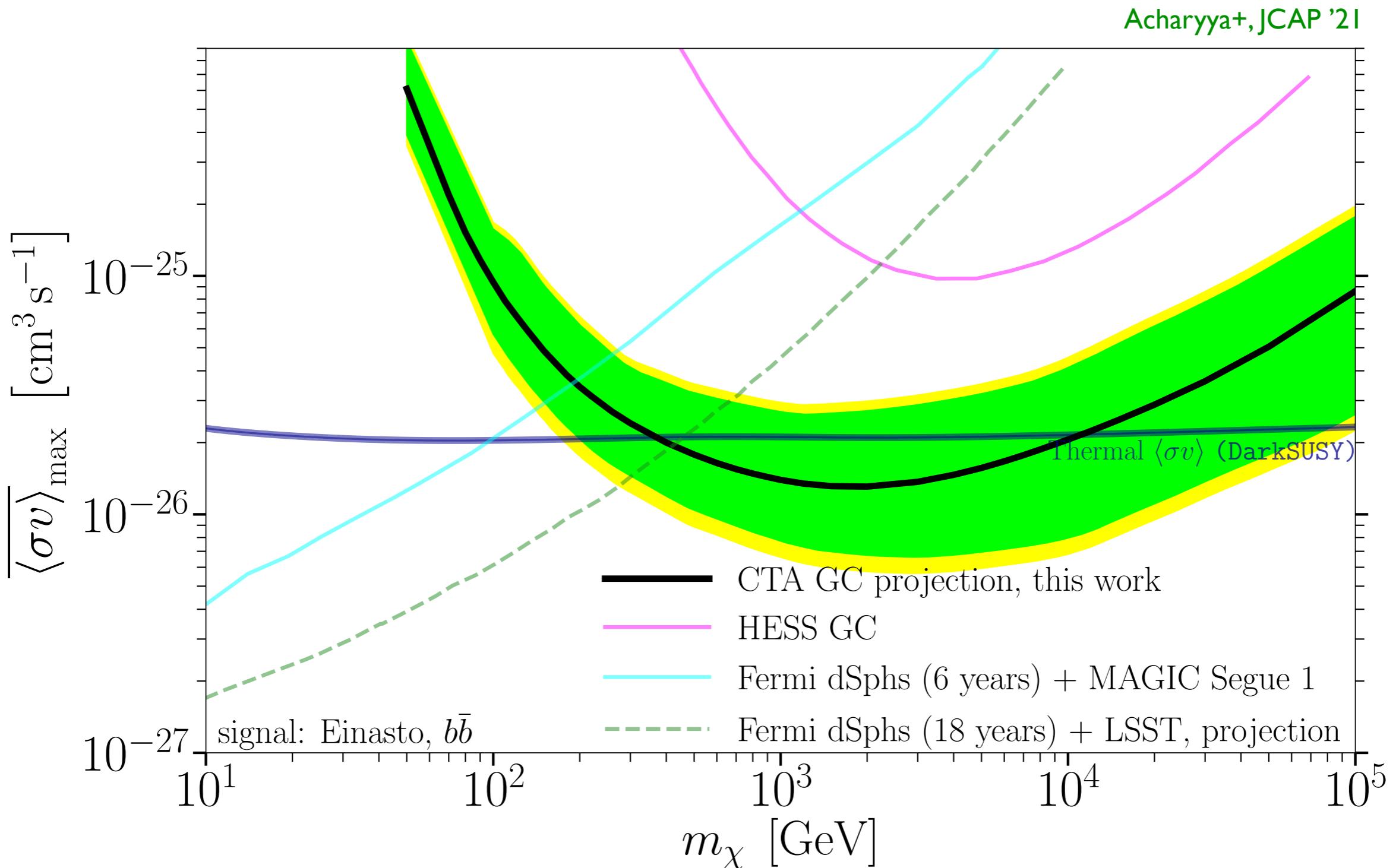


**indirectly**

- WIMP DM is seriously **pressured**,  
but certainly not (yet) '**dead**' !

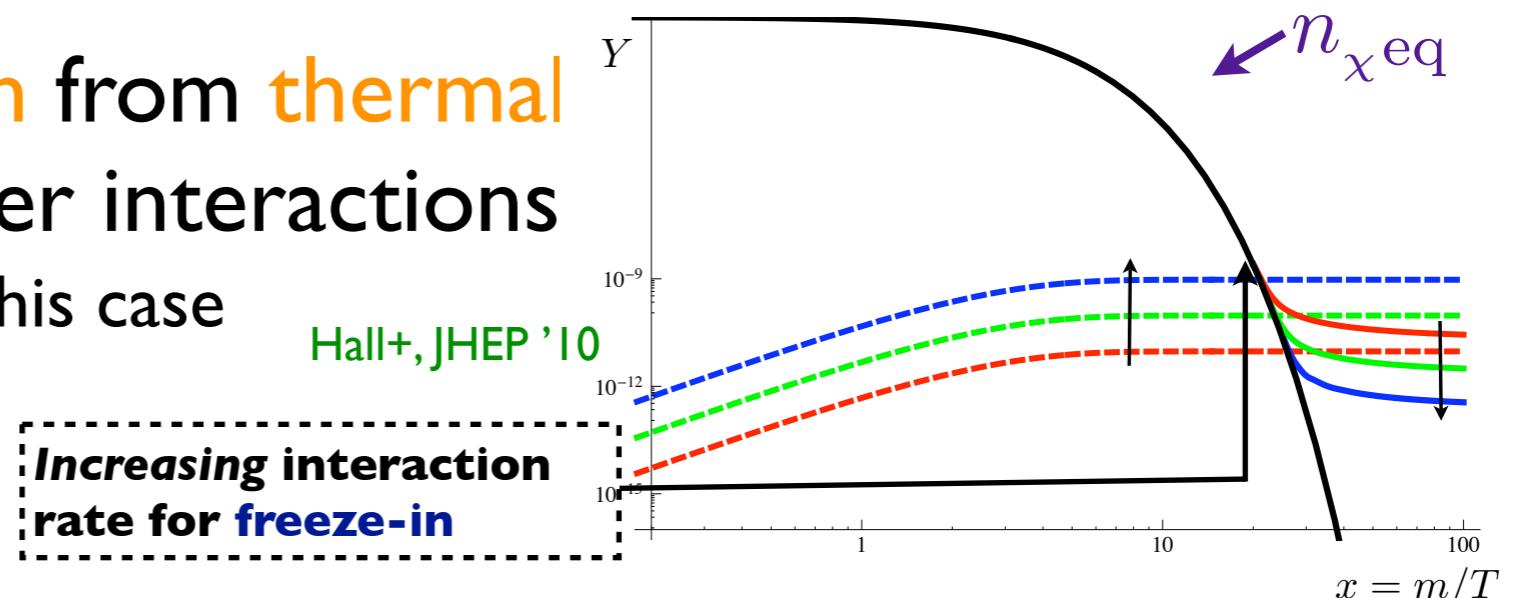
Arcadi+, EPJC '18  
Athron+, 2106.02056  
(+ many more)

# Example: indirect detection constraints

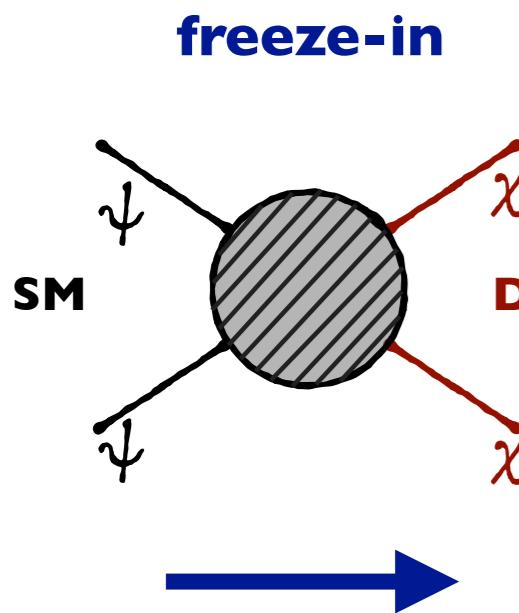


# Feably Interacting Massive Particles

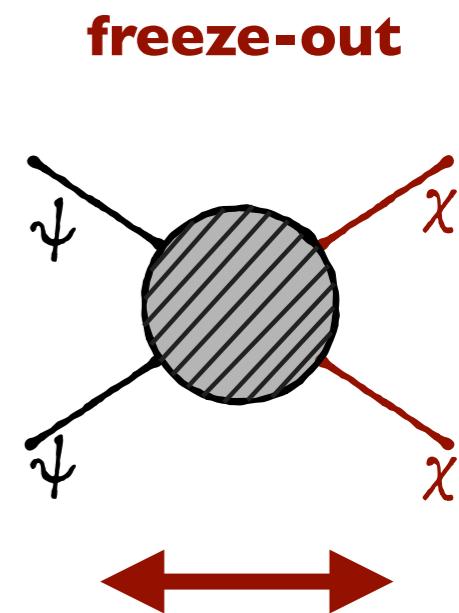
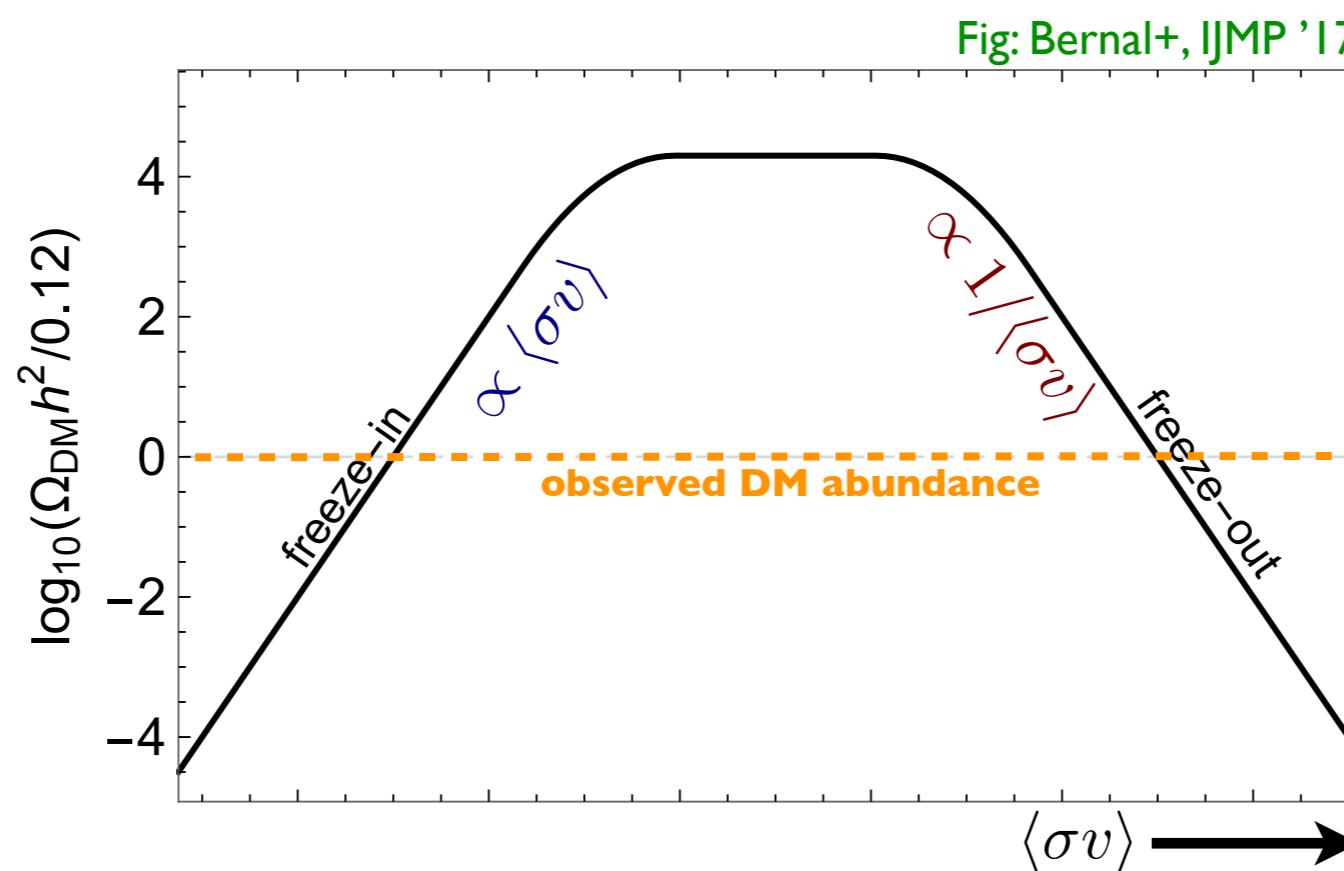
- alternative production from thermal bath with much smaller interactions
- DM never equilibrates in this case



- Smooth transition between two regimes:



*depends on initial conditions*



*insensitive to initial conditions*

# Collision term for FIMPs

$\rightarrow C[f_\chi] = \langle \sigma v \rangle_{\chi\chi \rightarrow \psi\psi} (n_\chi^{\text{MB}})^2$

annihilation of would-be MB population  
 $\rightarrow$  **actual** (in eq)

- Only 2 integrals after exploiting spherical symmetry:  
 $[\gamma : \text{Lorenz boost to CMS}; \quad \tilde{s} = s/(4m_\chi^2)]$

$$\langle \sigma v \rangle_{\chi\chi \rightarrow \psi\psi} = \frac{8x^2}{K_2^2(x)} \int_1^\infty d\tilde{s} \, \tilde{s} \, (\tilde{s} - 1) \int_1^\infty d\gamma \, \sqrt{\gamma^2 - 1} e^{-2\sqrt{\tilde{s}}x\gamma} \sigma_{\chi\chi \rightarrow \psi\psi}(s, \cancel{\gamma})$$

$\rightarrow K_1(2\sqrt{\tilde{s}}x)/(2\sqrt{\tilde{s}}x)$  ✓

TB, Heeba, Kahlhoefer & Vangnnes,  
JHEP '22  
(see also Lebedev & Toma, PLB '19  
Arcadi+, JHEP '19)

$$\sigma_{\chi\chi \rightarrow \psi\psi}(p, \tilde{p}) = \frac{(2\pi)^4}{4N_\psi E\tilde{E}v_{M\emptyset l}} \int \frac{d^3k}{(2\pi)^3 2\omega} \int \frac{d^3\tilde{k}}{(2\pi)^3 2\tilde{\omega}} \delta^{(4)}(\tilde{p} + p - \tilde{k} - k) |\overline{\mathcal{M}}|^2 \{1 \pm f_\psi(\omega)\} \{1 \pm f_\psi(\tilde{\omega})\}$$

- In *this* formulation, direct analogy with **WIMP** case!

- ➔ Can recycle sophisticated numerical tools for thermal averages
  - ➊ Easier to estimate higher-order corrections

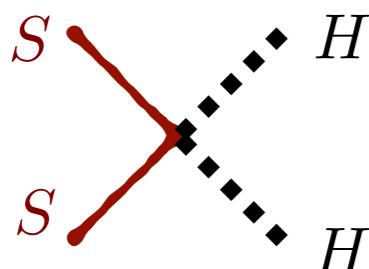


# Case study

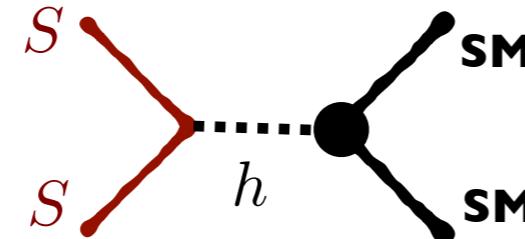
- Scalar Singlet DM

$$\mathcal{L} = \frac{1}{2}\partial_\mu S\partial^\mu S + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{2}\lambda_{hs} S^2 |H|^2 + \frac{1}{4}\lambda_s S^4$$

- before EWSB:



- after EWSB:



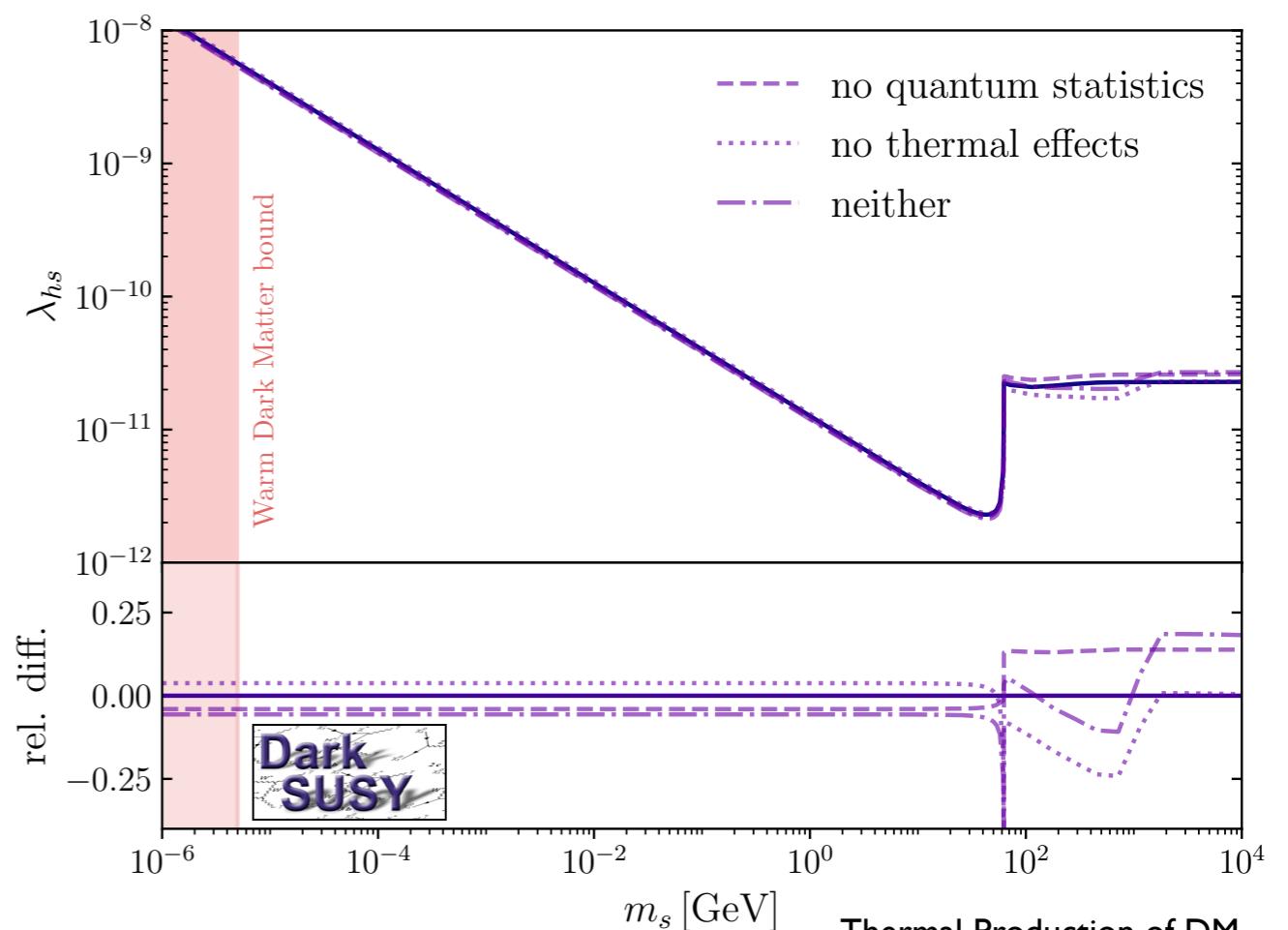
Silveira & Zee, PLB '85

production

- Freeze-in coupling required for **correct** relic density

TB, Heeba, Kahlhoefer  
& Vangnes, JHEP '22

- unlike for (non-rel.) freeze-out, **quantum statistics** do affect final result [@  $\mathcal{O}(10\%)$ ] !



# DarkSUSY

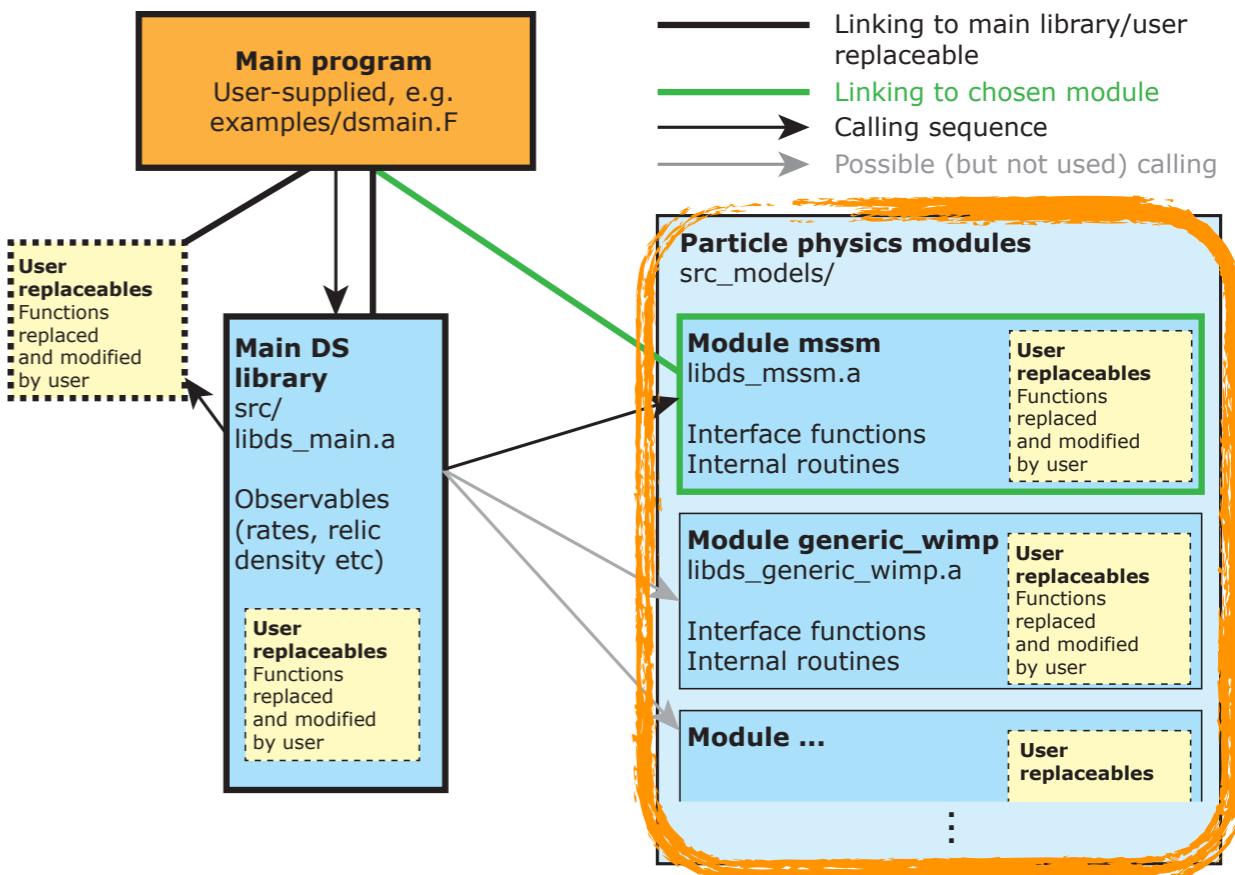


TB, Edsjö, Gondolo,  
Ullio & Bergström,  
JCAP '18

[http://  
darksusy.hepforge.org](http://darksusy.hepforge.org)

**Since version 6:  
no longer restricted to  
supersymmetric DM !**

- Numerical package to calculate ‘all’ DM related quantities:
  - relic density + kinetic decoupling (**also for  $T_{\text{dark}} \neq T_{\text{photon}}$** )
  - generic SUSY models + laboratory constraints implemented
  - cosmic ray propagation
  - particle yields for generic DM annihilation or decay
  - indirect detection rates: gammas, positrons, antiprotons, neutrinos
  - direct detection rates
  - ...

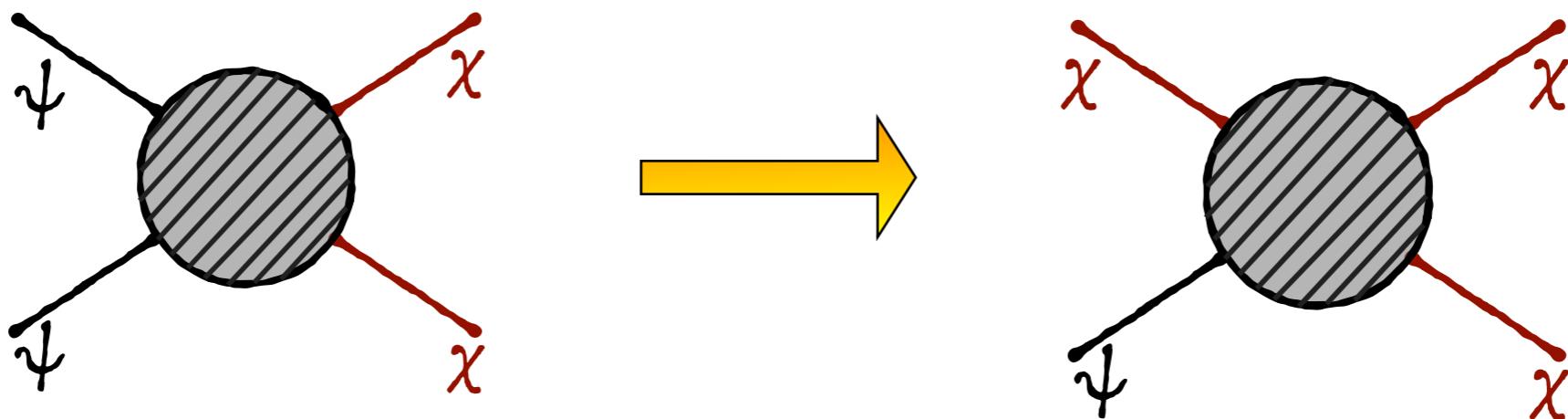


**since 6.1: DM self-interactions**

**since 6.2: ‘reverse’ direct detection**  
(also  $Q^2$ -dependent scattering!)

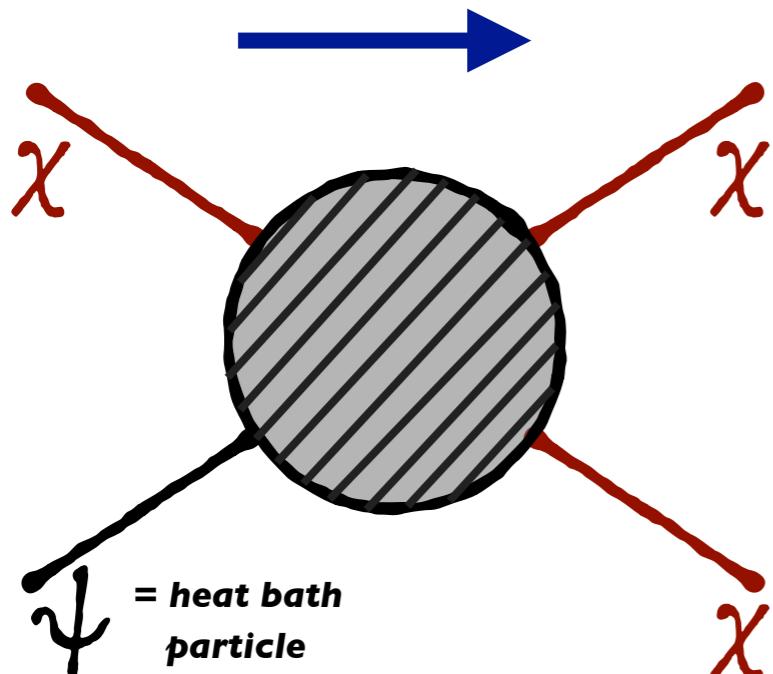
**since 6.3: freeze-in**

# Part III



# A new production mechanism

- ‘Pandemic’ dark matter



TB, Depta, Hufnagel, Rudermann  
& Schmidt-Hoberg, 2103.16572

Hryczuk & Laletin, 2104.05684

$$\dot{n}_\chi + 3H n_\chi = n_\chi n_\psi^{\text{eq}} \langle \sigma v \rangle$$

[for  $n_\chi \ll n_\psi^{\text{eq}}$ ]

→ reproduction number, or ‘R-value’:

- The ‘SIR’ compartmental model

*A Contribution to the Mathematical Theory of Epidemics.*

By W. O. KERMACK and A. G. MCKENDRICK.

(Communicated by Sir Gilbert Walker, F.R.S.—Received May 13, 1927.)

$S$  # susceptible individuals

$I$  # infected individuals

# recovered ( $R = \text{tot} - S - I$ )

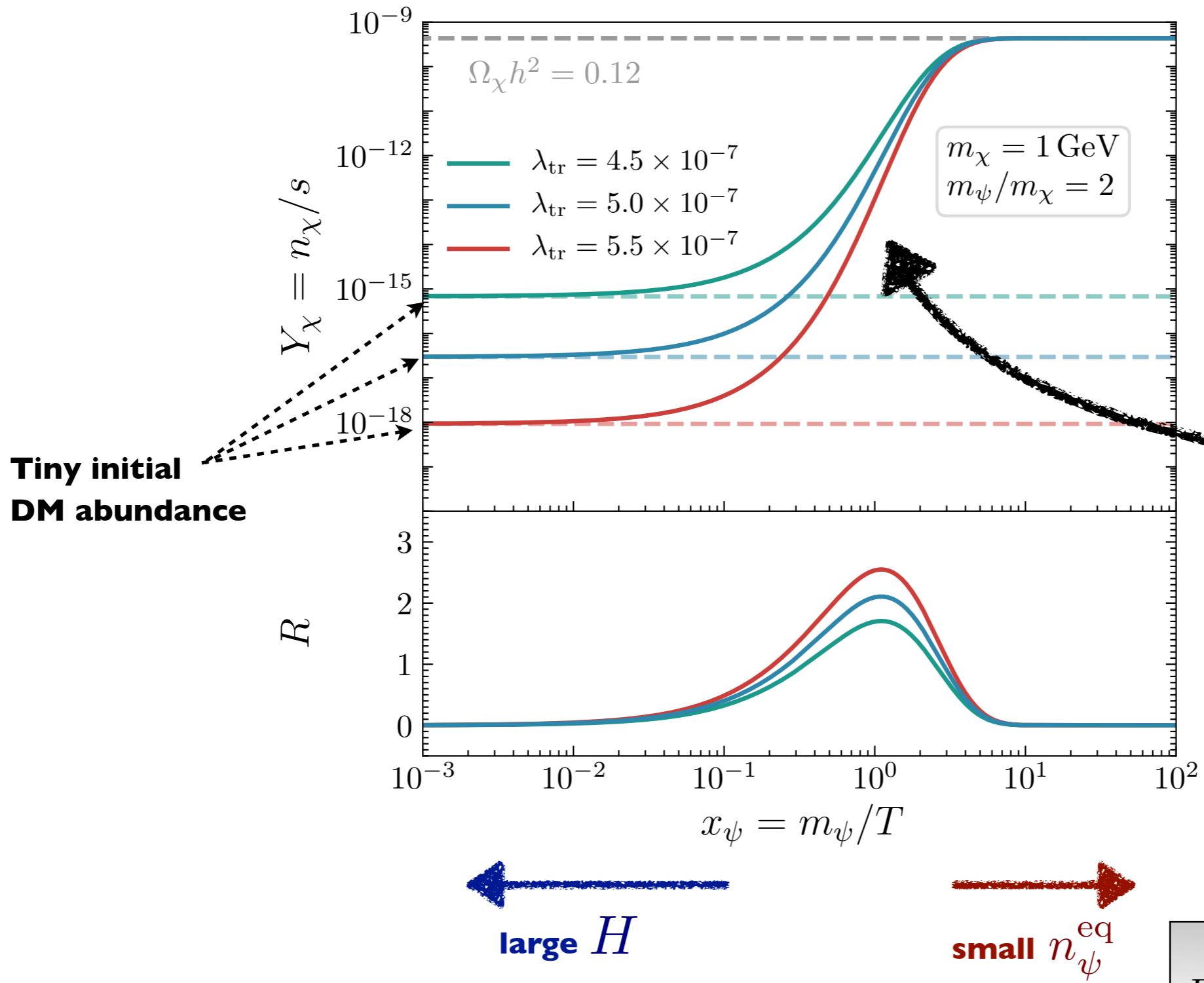
$\beta$  infection rate

$\gamma$  recovery rate

$$\dot{I} = \beta S I - \gamma I$$

$$R \equiv \frac{\beta S}{\gamma} = \frac{n_\psi^{\text{eq}} \langle \sigma v \rangle}{3H}$$

# Exponential DM production



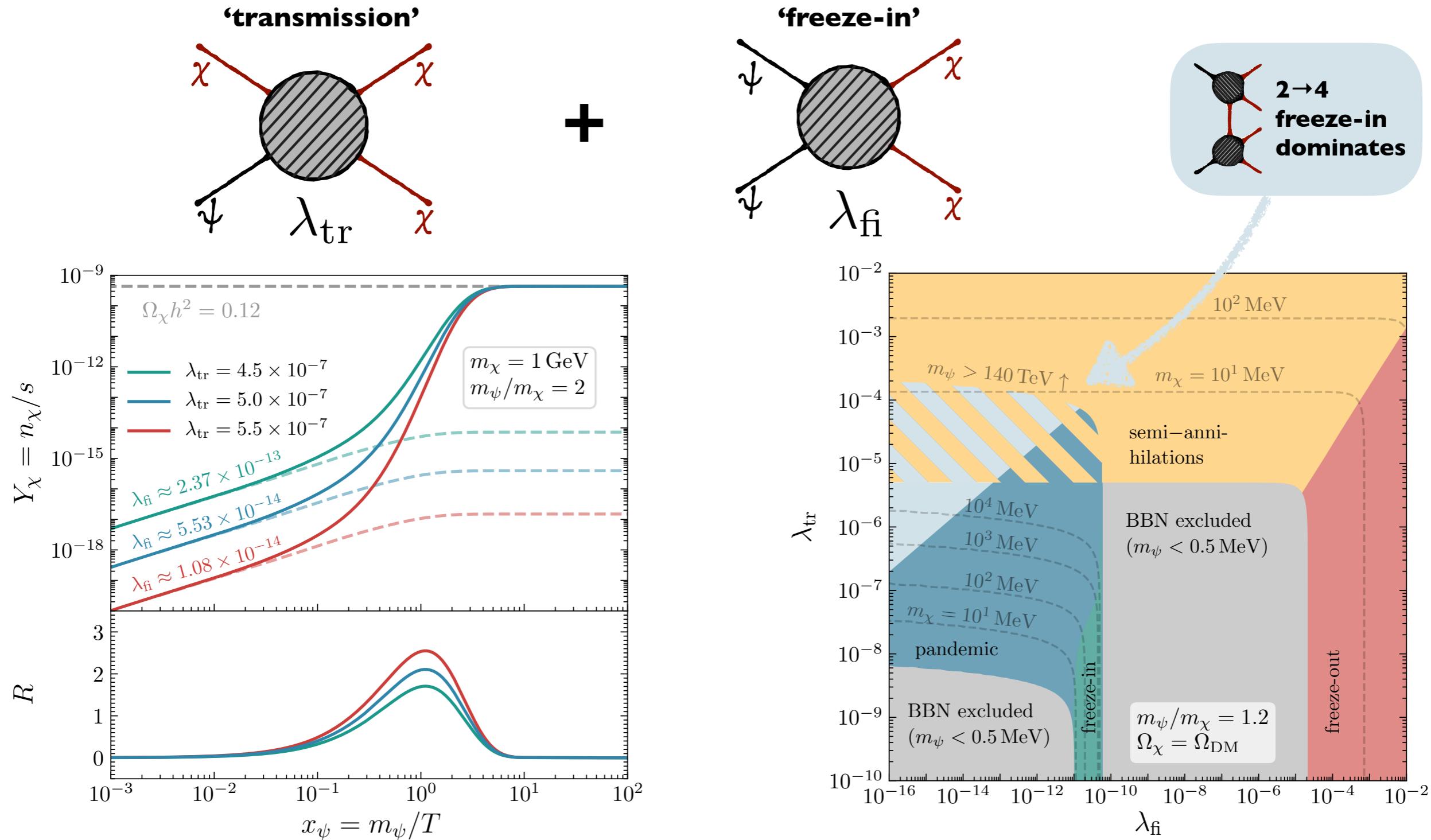
$$\mathcal{L} \supset (\lambda_{\text{tr}}/3!) \psi \chi^3$$

$= -i \lambda_{\text{tr}}$

**toy model**

$$R \equiv \frac{\beta S}{\gamma} = \frac{n_\psi^{\text{eq}} \langle \sigma v \rangle}{3H}$$

# Adding freeze-in production

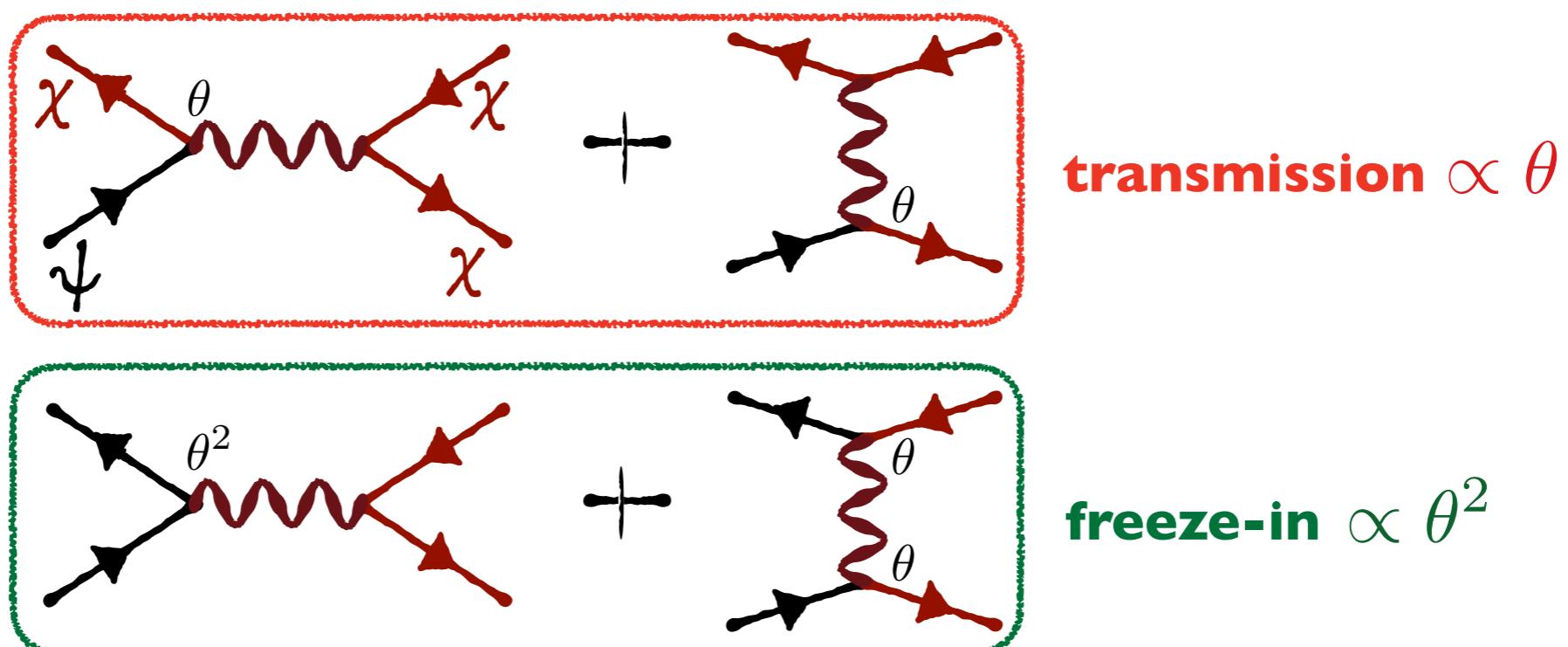


→ ‘Pandemic’ production is a very generic mechanism for the genesis of DM!

# Signals ?

- Necessarily model-dependent
  - ‘Pandemic DM’ describes a **class** of models, just like ‘WIMP’ does
- Q: Is there a *generic* way to realize  $\langle\sigma v\rangle_{\text{fi}} \ll \langle\sigma v\rangle_{\text{tr}}$  ?
- A: yes — just add a **dark sector mediator** and **mass mixing!**

$$\mathcal{L} \supset -\delta m (\bar{\psi}\chi + \bar{\chi}\psi) - g\bar{\chi}\not{V}\chi \quad \rightarrow \quad \mathcal{L} \supset -g[\bar{\chi}\not{V}\chi + \theta(\bar{\psi}\not{V}\chi + \bar{\chi}\not{V}\psi) + \theta^2\bar{\psi}\not{V}\psi]$$



# Sterile neutrinos

- A **right-handed** neutrino would be **neutral** under *all* SM gauge forces

- An **excellent, well-motivated** dark matter **candidate**

- Production by SM processes: **oscillations** with active neutrinos, combined with NC and CC scatterings

## SM fermions

|          |                                      |                                    |                                    |
|----------|--------------------------------------|------------------------------------|------------------------------------|
| mass →   | 2.4 MeV                              | 1.27 GeV                           | 171.2 GeV                          |
| charge → | $\frac{2}{3}$                        | $\frac{2}{3}$                      | $\frac{2}{3}$                      |
| name →   | u<br>Left<br>up                      | c<br>Left<br>charm                 | t<br>Left<br>top                   |
| Quarks   | d<br>Left<br>down                    | s<br>Left<br>strange               | b<br>Left<br>bottom                |
| Leptons  | $\nu_e$<br>0 eV<br>electron neutrino | $\nu_\mu$<br>0 eV<br>muon neutrino | $\nu_\tau$<br>0 eV<br>tau neutrino |
|          | e<br>Left<br>electron                | $\mu$<br>Left<br>muon              | $\tau$<br>Left<br>tau              |

Dodelson &  
Widrow, PRL '94

- Unfortunately, this scenario is **ruled out** by observations...

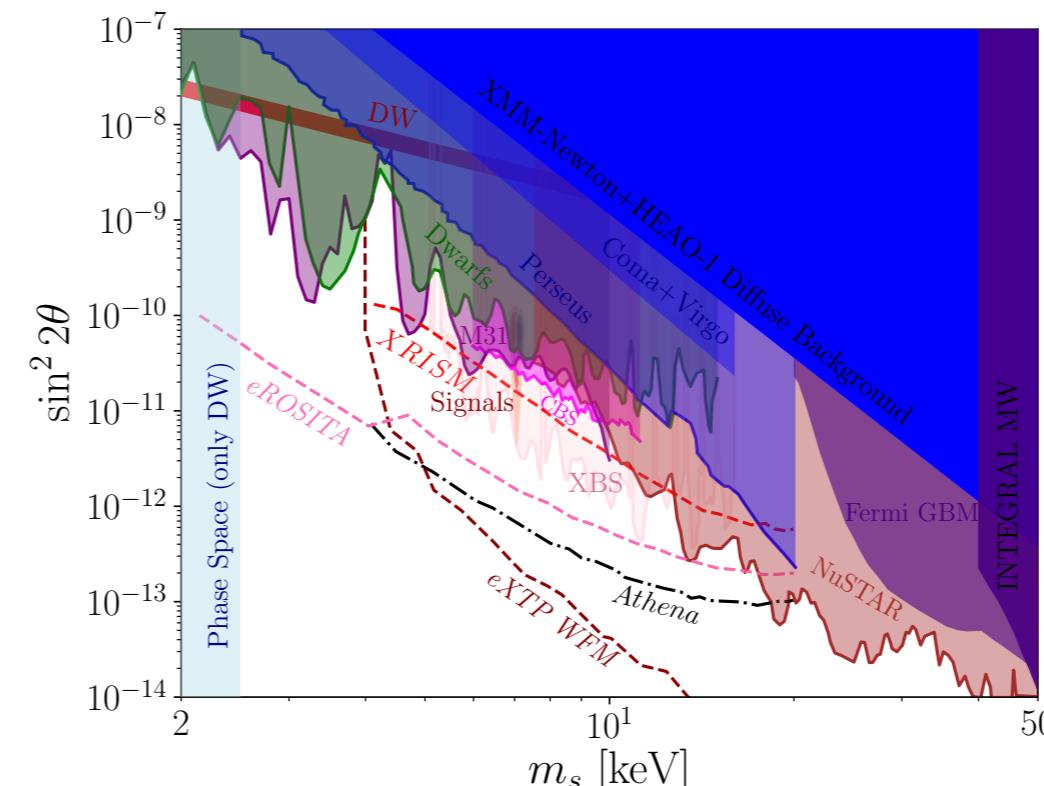
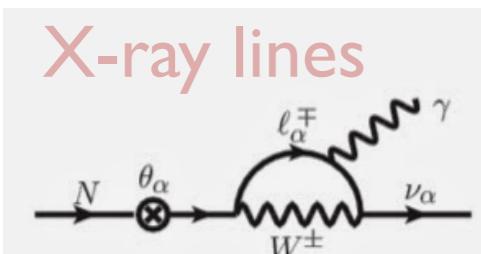


Fig.: Abazajian+, 2203.7377

# Interacting sterile neutrinos

TB, Depta, Hufnagel, Kersten, Ruderman & Schmidt-Hoberg, arXiv:2206.10630

- Let's add a **scalar  $\phi$**  that only couples to the sterile neutrinos

$$\mathcal{L} \supset \frac{y}{2} \phi \bar{\nu}_s \nu_s \rightarrow \frac{y}{2} \phi [\sin^2 \theta \bar{\nu}_\alpha \nu_\alpha - \sin \theta \cos \theta (\bar{\nu}_\alpha \nu_s + \bar{\nu}_s \nu_\alpha) + \cos^2 \theta \bar{\nu}_s \nu_s]$$

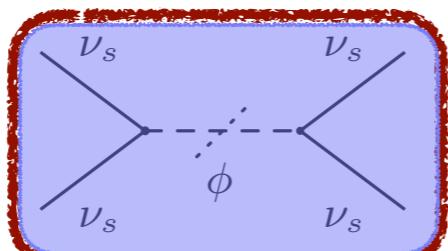
$$m_\phi > 2m_s$$

- Early times ( $\sim$ QCD PT): standard **DW** production

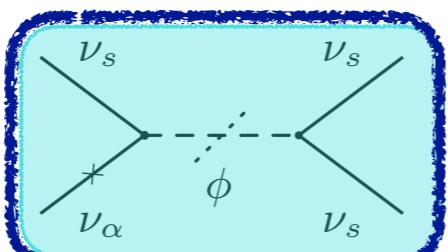
- Evolution afterwards:

**solid**: benchmark point with large  $\theta$ , small  $y$   
**dashed**: benchmark point with small  $\theta$ , large  $y$

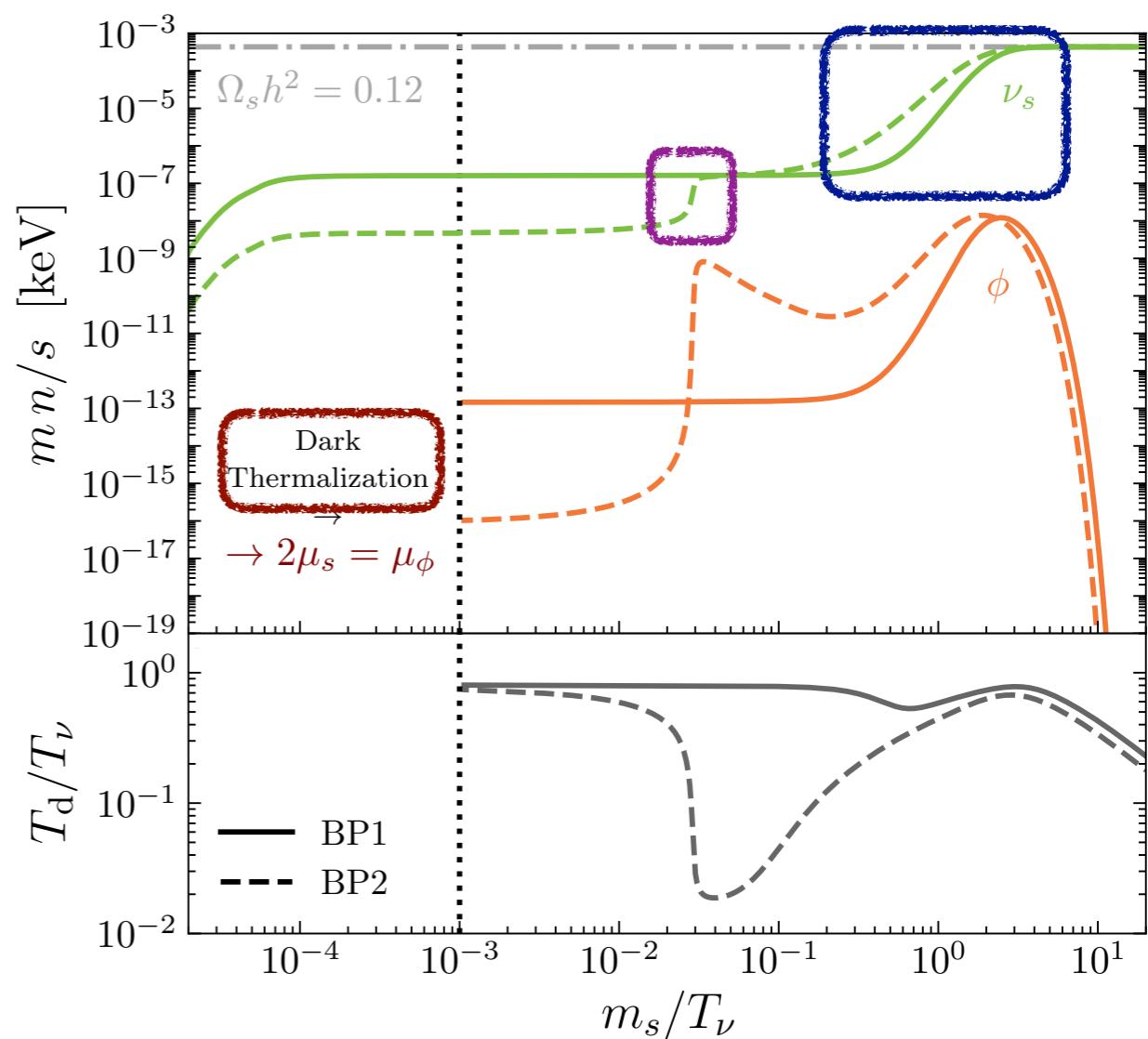
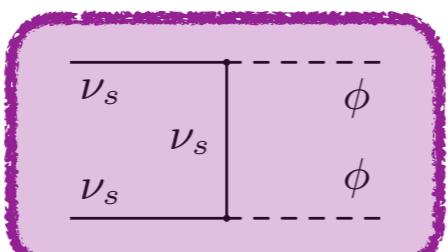
- Thermalization in dark sector



- Exponential growth



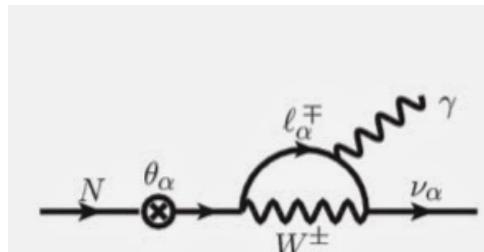
- Reproductive freeze-in



# Sterile neutrinos... revived !

TB, Depta, Hufnagel, Kersten, Ruderman & Schmidt-Hoberg, arXiv:2206.10630

- Observational constraints
- (Standard) X-ray lines



- $\nu_s$  self-interactions

$$\sigma_T/m_s \lesssim 1 \text{ cm}^2/\text{g}$$

cf. Tulin & Yu, PR '18

maybe 0.1 possible... (?)

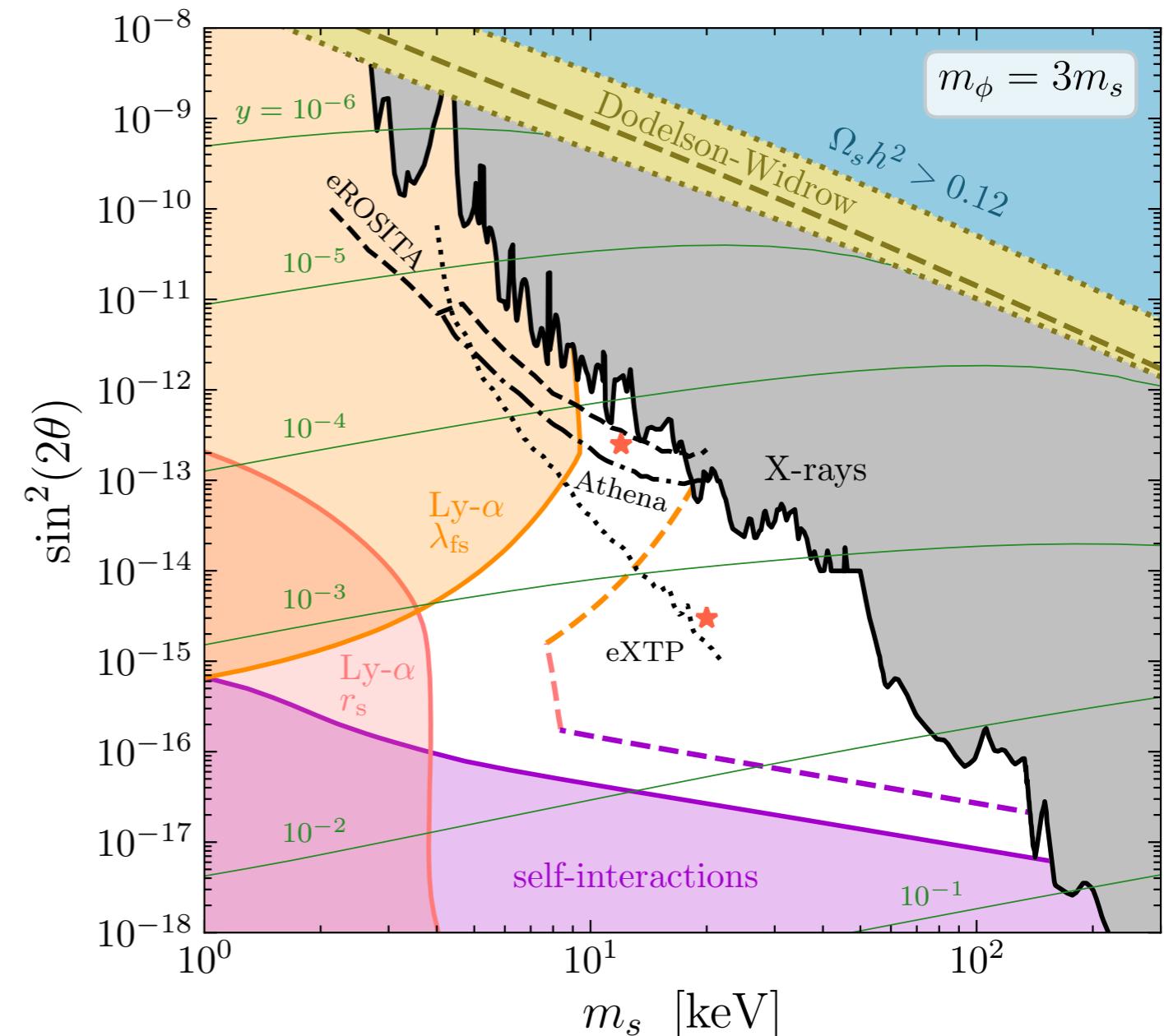
- Lyman-alpha

recast  $m_{\text{WDM}} > 1.9 \text{ keV}$  to  
Garzilli+, MNRAS '21

$$\lambda_{\text{FS}} < 0.24 \text{ Mpc}$$
$$r_s < 0.36 \text{ Mpc}$$

maybe  $m_{\text{WDM}} > 5.3 \text{ keV}$  possible... (?)

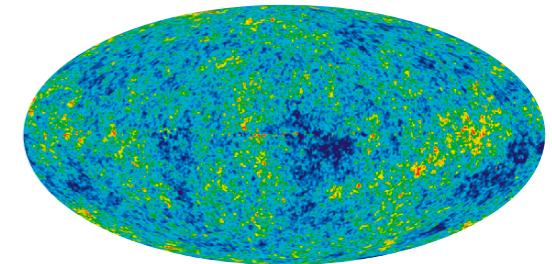
Palanque-Delabrouille+, JCAP '20



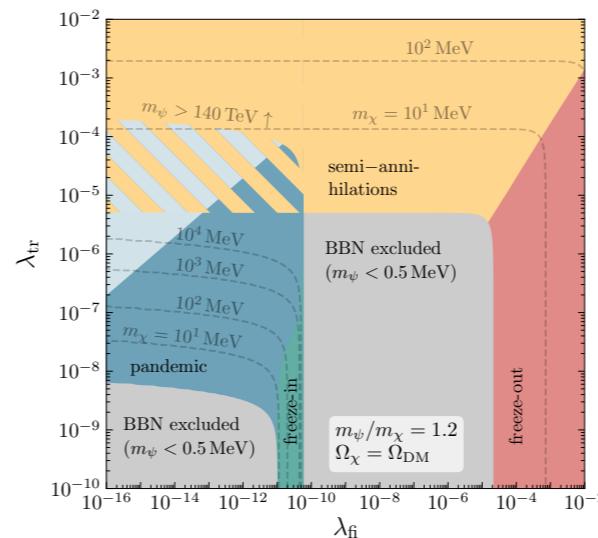
- New parameter space
- Bounded from above and below
- Significant parts in observational reach

# Conclusions

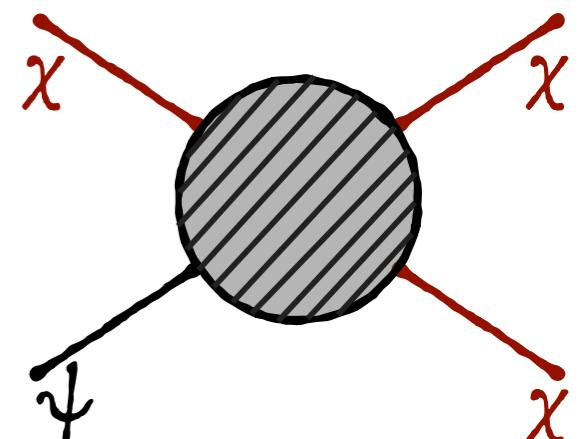
- Dark matter = evidence for BSM physics



- There are a handful of generic ways of DM genesis from primordial heat bath



- Pandemic dark matter is a novel such mechanism. Mathematical analogy to spread of diseases works almost scarily well [see arXiv v1 for exploiting most of it 😊]



## Thanks for your attention!