The Higgs and Cosmology

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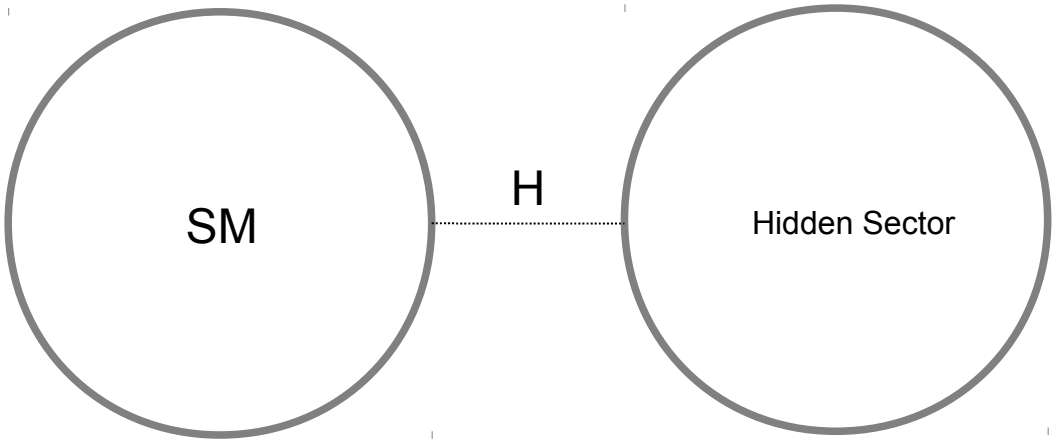
- the Higgs and the hidden sector
- the Higgs and dark matter
- the Higgs and inflation

The Higgs and the hidden sector



Lowest order operators ("Higgs Portal") :

$\overline{H}H\, S^2 \quad + \quad \ldots \qquad\qquad$ (scalar)

$\overline{H}H\, V_\mu V^\mu \; + \quad \ldots \qquad\qquad$ (vector)

$\overline{H}H\, \bar{\chi}\chi / \Lambda \; + \quad \ldots \qquad\qquad$ (fermion)

"Portal" due to Patt,Wilczek'06 (earlier : Silveira,Zee'85; Shabinger,Wells'05;...)

Special role of the Higgs :

Silveira, Zee '85
Veltman, Yndurain '89
...

$|H|^2$ = the only gauge and Lorentz-inv. dim-2 operator

$$L = a\,|H|^2\,S^2 + b\,|H|^2\,S$$

(S = "hidden" scalar)

b=0 (S has hidden charge):

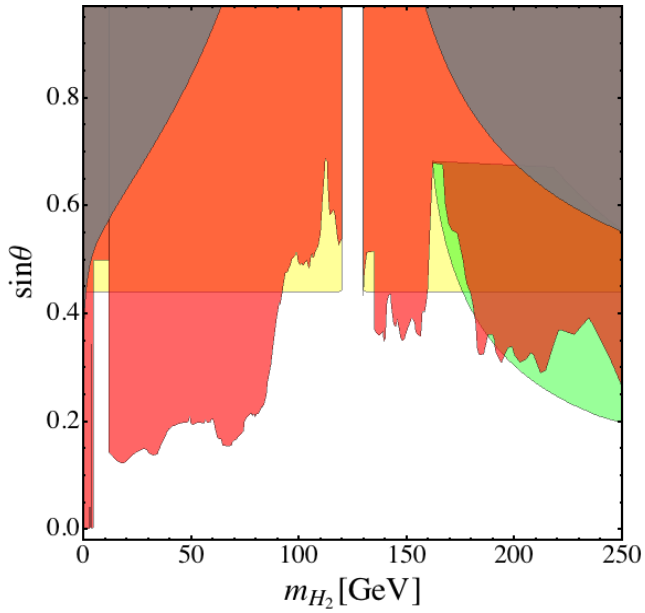
$$L = a\,|H|^2\,S^2$$

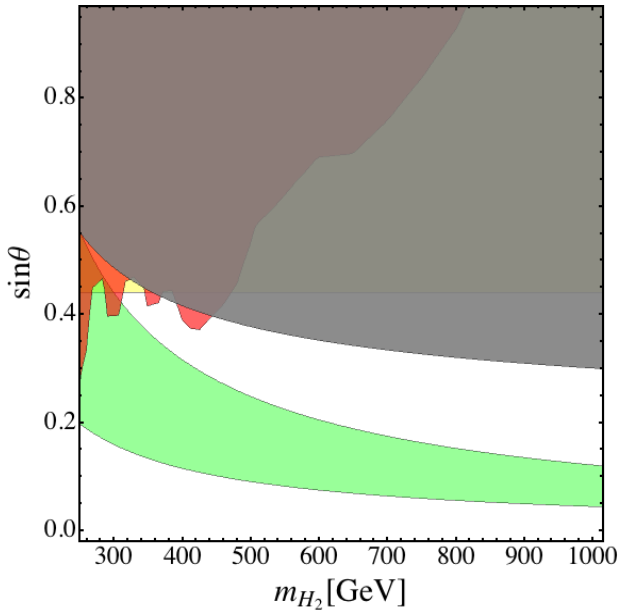
"S" is stable and couples weakly to SM --> DARK MATTER (?)

Constraints on Higgs-singlet mixing :

H_1 = 125 GeV Higgs ; H_2 = extra Higgs ; θ = mixing angle

Falkowski, Gross, OL '15

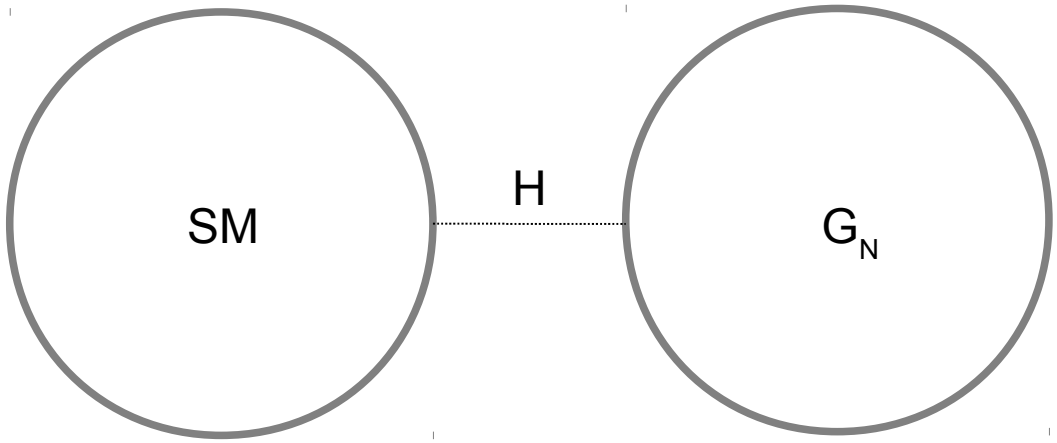




Grey = EW precision data , Yellow = LHC Higgs couplings , Reddish = B-physics, LEP, LHC

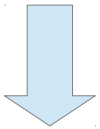
Green (optional) = Higgs potential stability/perturbativity up to M_{Pl}

The Higgs and dark matter



$V \sim \bar{H}H\,\bar{S}S$ ⇨ H-S mixing ⇨ **h couples to G_N**

Lie groups possess discrete symmetries



gauge fields as dark matter

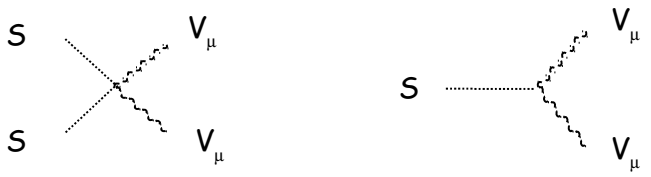
E.g. U(1) : $A_\mu \rightarrow -A_\mu$

Lee, OL, Mambrini '11

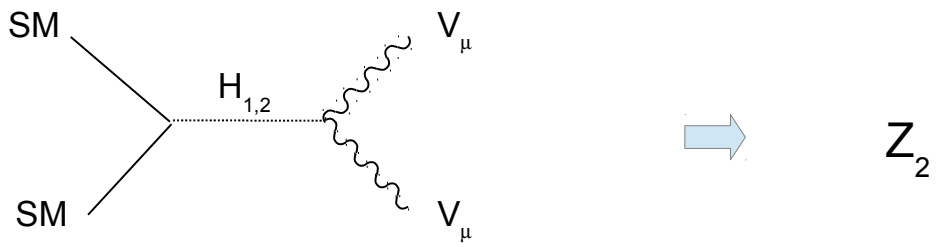
Higgs mechanism in the hidden sector :

$$L = -1/4\, F_{\mu\nu} F^{\mu\nu} + D_\mu S^* D^\mu S - V(S) + \lambda/4\, \bar{H}H\, S^*S$$

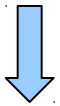
$$S \longrightarrow VEV$$



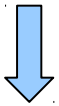
SM couplings:



gauge invariance (+ minimal field content)



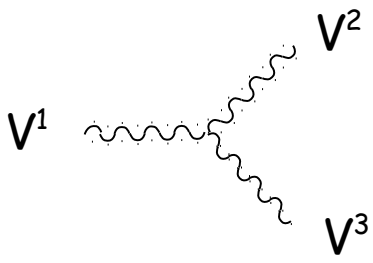
Z_2



gauge fields are natural DM candidates

Hambye '08
Gross,OL, Mambrini '15

Non-abelian case:



$\Rightarrow$ Z_2 odd

But there are 2 Z_2's:

$$V^{1,2} \longrightarrow -V^{1,2} \quad , \quad V^3 \longrightarrow V^3$$

$$V^{1,3} \longrightarrow -V^{1,3} \quad , \quad V^2 \longrightarrow V^2$$

$\Rightarrow$ V^a = stable

hidden Higgs = $\begin{pmatrix} 0 \\ v \end{pmatrix}$

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}; \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}; \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

Z_2 : reflects real generators

complex (or "charge") conjugation = outer automorphism

Z_2' : reflects off-diagonal generators with non-zero elements in the first row

gauge transformation

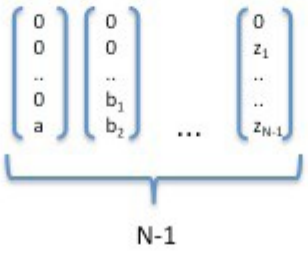
General SU(N) case

$$\left[T^a, T^b\right] = i\, f^{abc}\, T^c$$

Z_2 : reflects real generators

Z_2' : reflects non-Cartan generators with non-zero first row (Pauli-like basis)

Higgsing:



$Z_2 \times Z_2'$ preserved if CP is conserved

Arcadi,Gross, OL,
Pokorski,Toma '16

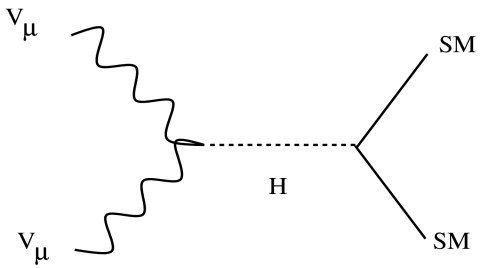
minimal SU(N): $Z_2 \times Z_2$

two stable components

E.g. spin1 + spin 0 DM

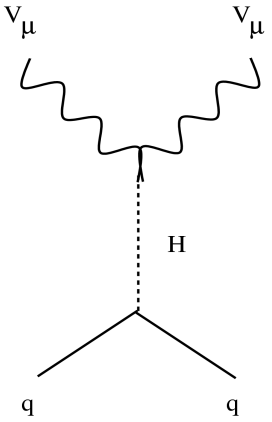
DM phenomenology :

annihilation



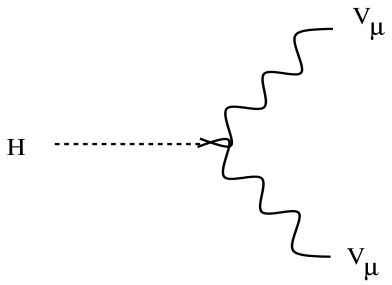
$< \sigma v >$

DM-nucleon scattering

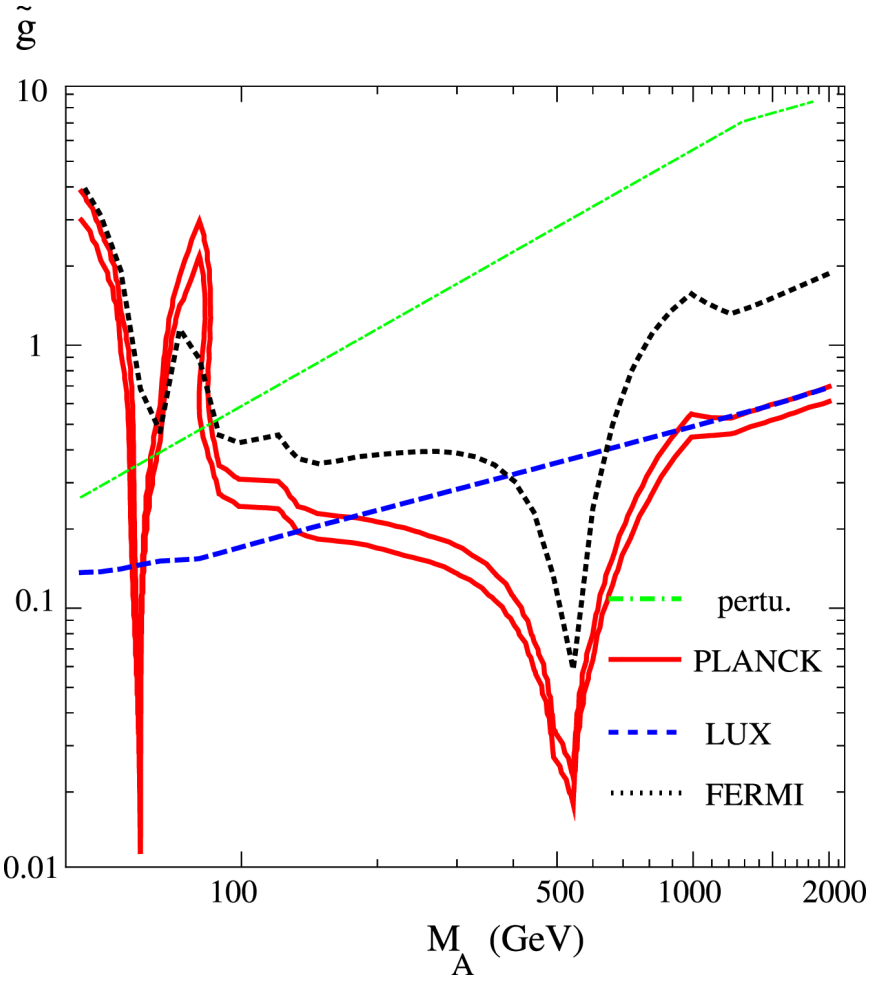


σ^{SI}_{S-P}

invisible Higgs decay



Γ^{inv}_H


$\tilde{g}$
10
1
0.1
0.01
100
500
1000
2000
M_A (GeV)
pertu.
PLANCK
LUX
FERMI

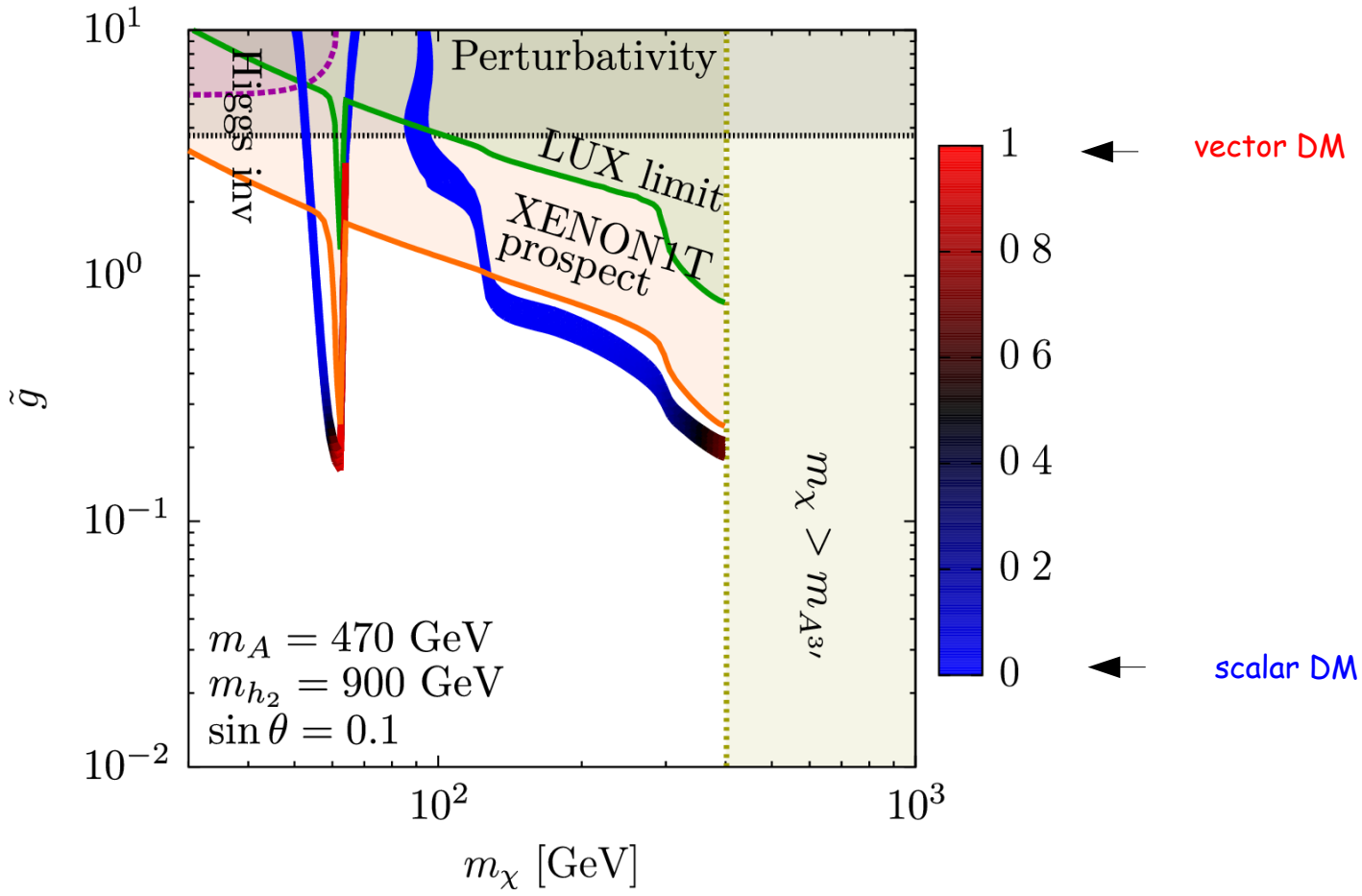
U(1)

sin θ = 0.3

mixed spin DM

SU(3)

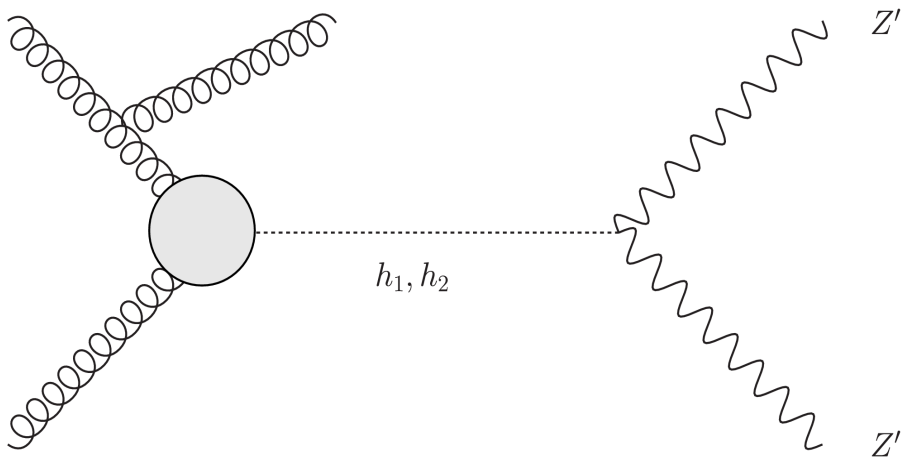
sin θ = 0.1



Kim, OL ,Schmeier '15

Higgsophilic gauge bosons at the LHC:

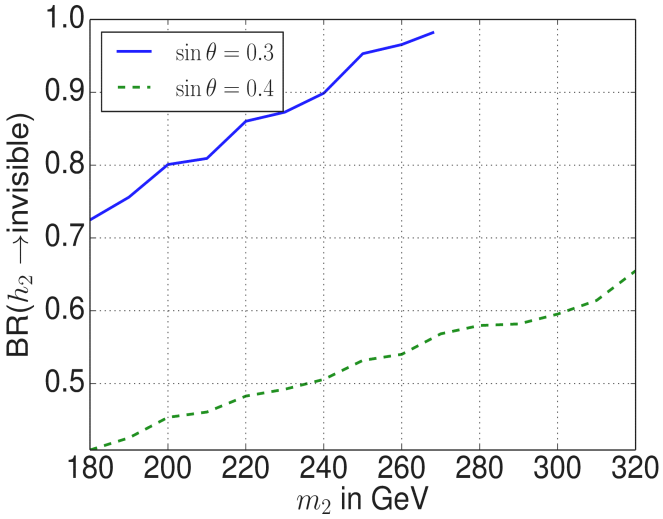
monojets (or VBF) + missing E_T

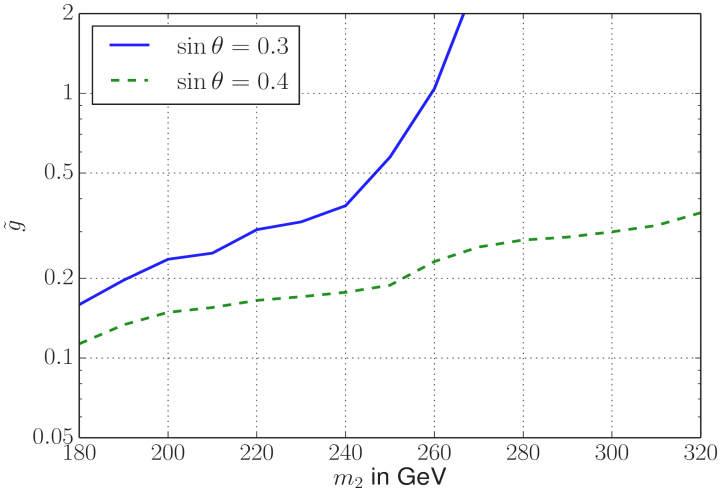


efficient for $m_2 > 2\, m_{Z'}$

Kim, OL ,Schmeier '15

Monojet results with 600 fb^{-1} :





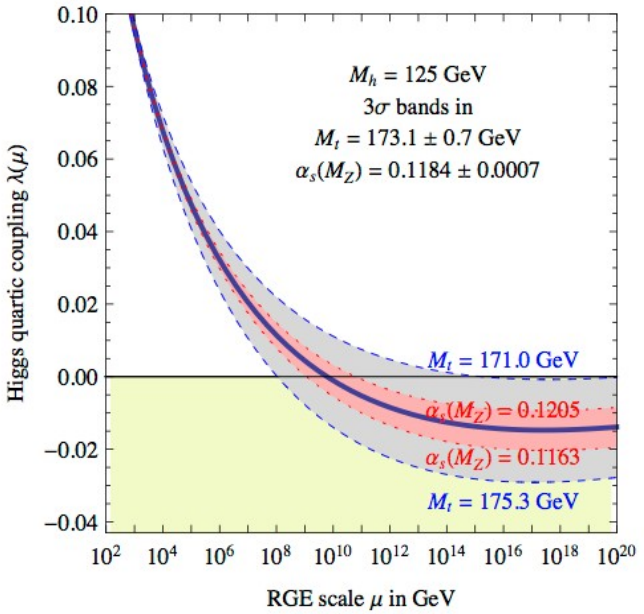
+ similar results for VBF with already 300 fb^{-1}

The Higgs and inflation

Buttazzo et al.'13

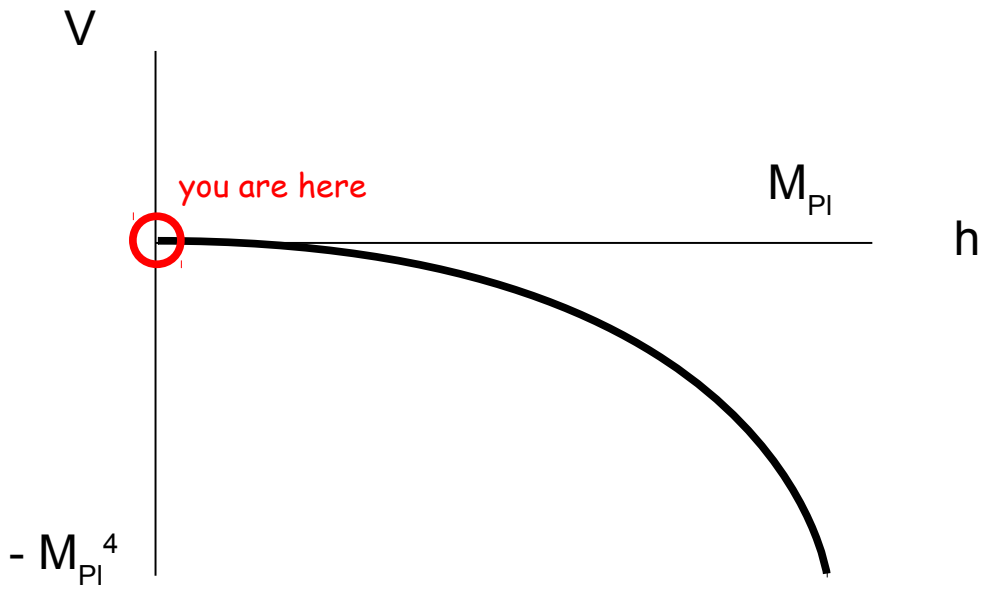
SM stability bound:

$$m_h > (129.6 \pm 1.5)\ \text{GeV}$$



(not settled : Alekhin et al. '12
Bezrukov et al. '12)

$h >> \Lambda \sim 10^{10}\,\mathrm{GeV}$ ⇨ $V \sim \frac{1}{4}\,\lambda(h)\,h^4$, $\lambda(h) < 0$



$\Lambda = 10^{-8}\,M_{Pl}$, barrier $= 10^{-32}\;M_{Pl}^4$

Problems :

- how did the Universe end up at h ~ 0 ?

- why did it stay there during inflation ?

Solutions :

- modify the Higgs potential during inflation

- just modify the Higgs potential

Solution 1:

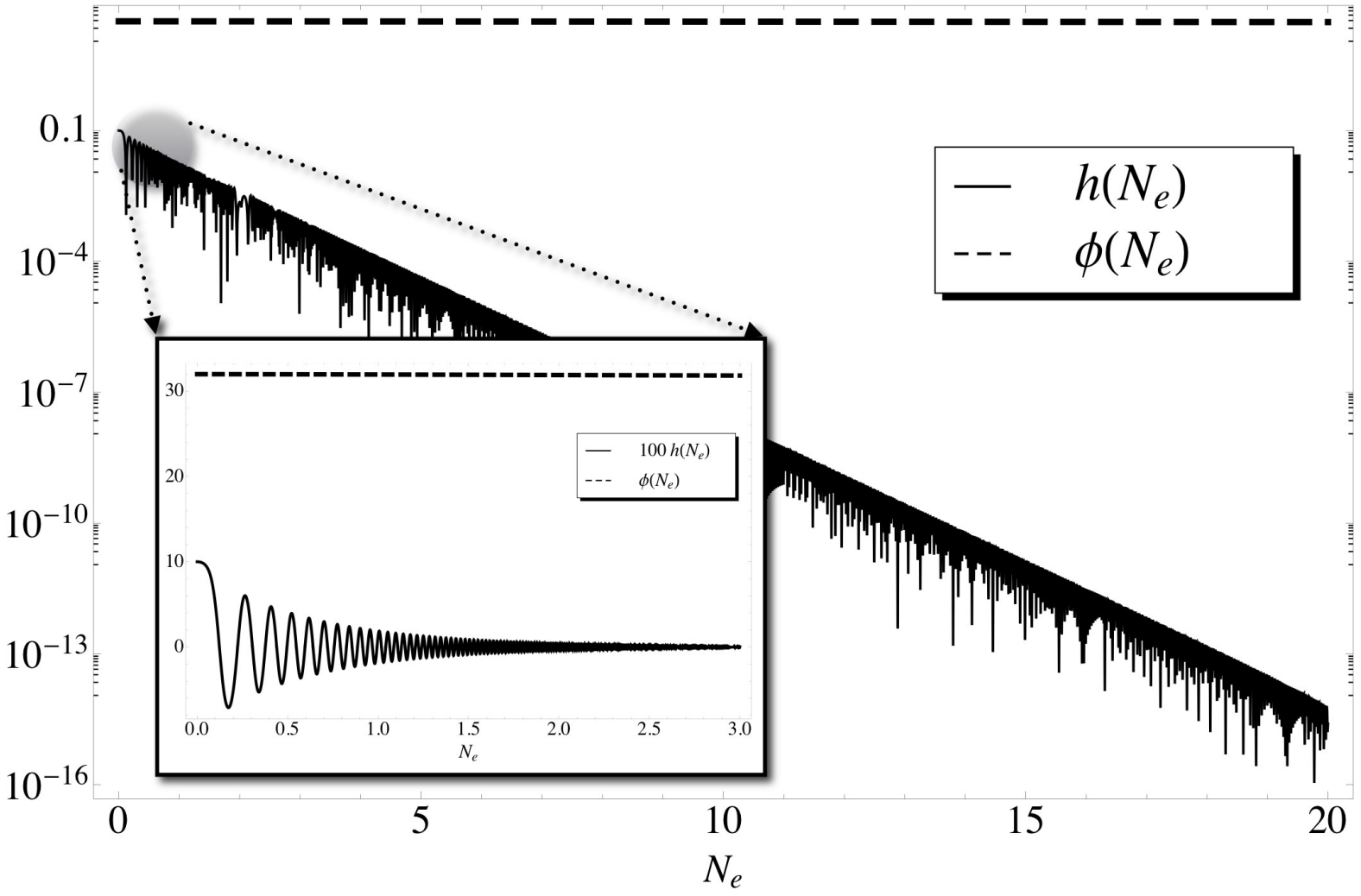
Higgs-inflaton coupling $\Delta V = \frac{1}{2} \lambda_{h\phi} h^2 \phi^2$ ("Higgs portal" coupling)

$\Delta V + V_{Higgs} > 0$ ⇨ $\phi_0 \sim 20\ M_{Pl}$, $\lambda_{h\phi} \sim 10^{-6}$

Large effective mass term ⇨ $h(t) \sim h(0) \exp(-3/2\ Ht)$

Higgs field is driven to zero during inflation !

Higgs/inflaton evolution (in M_{Pl}):



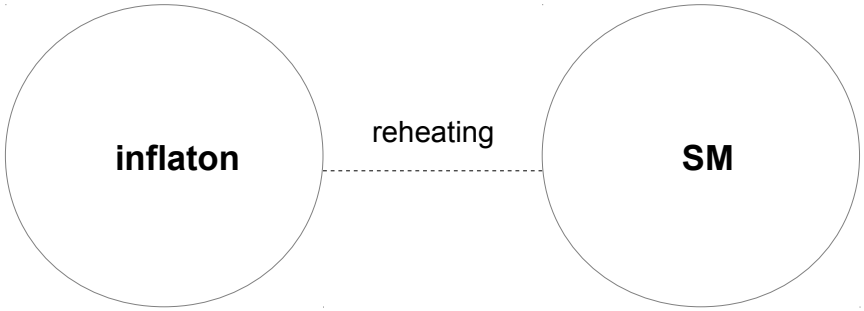
Gross, OL, Zatta '15

The Higgs-inflaton coupling is

- generated radiatively

- required by renormalizability

- important if greater than $\mathbf{10^{-10}}$

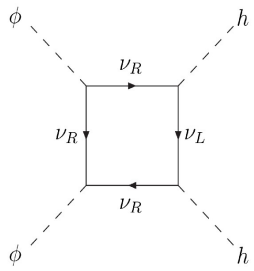


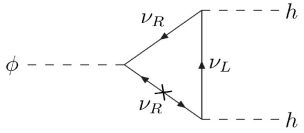
Gross, OL, Zatta '15

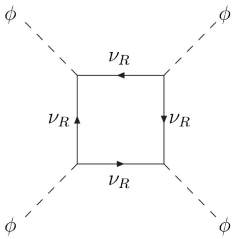
Example 1:

reheating through RH neutrinos

$$\Delta L = \lambda \, \phi \, \nu_R \, \nu_R$$







divergent, renormalize at M_{Pl} : $\lambda_i (M_{Pl}) = 0$

$\Rightarrow$ $\lambda_{h\phi}(H) < 2 \times 10^{-7}$

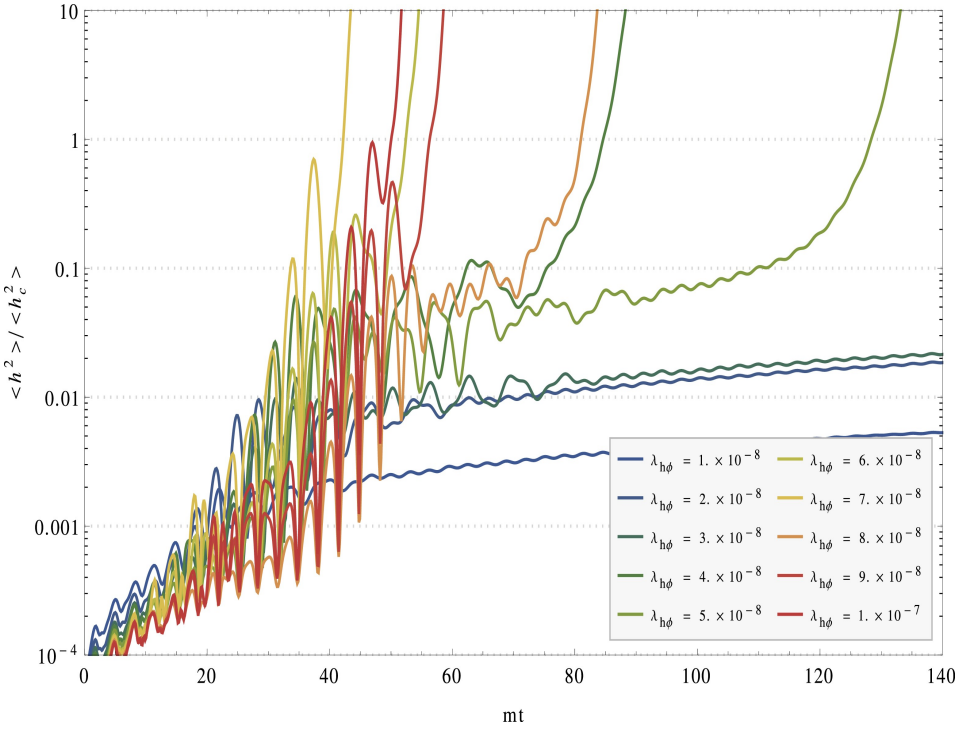
The Higgs-inflaton coupling is essential for the Higgs evolution

$$m_h^2 \sim \lambda_{h\phi}\, \phi^2 > H^2$$

Enqvist, Karciauskas, OL, Rusak, Zatta '16

Higgs vacuum destabilization through preheating:

$$V \sim \lambda\ h^2\,\phi^2 + \sigma\ h^2\,\phi$$



Conclusion

- Higgs sector is special
- key to the hidden sector / DM / inflation
- need precise Higgs data