

UiO Department of Physics

The Faculty of Mathematics and Natural Sciences

" Bayesian inference of Self-Interactions in the Dark Matter Halos of Milky Way Dwarfs "



ALSO SUPPORTED BY:

NPFIavour The Flavour of New Physics



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Nov 29 2017, Oslo, Norway





We live in the "Cosmological Precision Era"

—> compelling questions intimately related to **NP** !

- E.g.: $\Omega_{DM} \sim 5 \times \Omega_b$, $O(\Omega_{DM+b}) \sim O(\Omega_{\Lambda})$... just parametric coincidences ?
 - Why today anti-matter/matter << 1?
 - The (small) Cosmological Constant: a (huge) hierarchy problem ?
 - Origin about non-baryonic matter?





Available mass windowDM self-interaction range m_{DM} σ_{DM}/m_{DM} $10^{-23}eV$ $10^{67}eV$ $\mathcal{O}(M_{Pl}^{-2}m_{DM}^{-1})$ $\mathcal{O}(cm^2g^{-1})$



STATE-OF-THE-ART CDM SIMULATIONS ABLE TO PROVIDE ACCURATE SNAPSHOTS OF KNOWN OBSERVED GALAXIES.

N-body simulations implementing CDM idea + dedicated recipes for baryonic physics reproduce realistic galaxy morphologies.

Cold Dark Matter (CDM)

A cosmologically cold and collisionless particle species beyond the Standard Model well describes all observations we have on large-scale structures of the Universe.

CDM nicely match observations on typical scales \geq O(100) kpc!



Illustris Project, Vogelsberger et al. (2014)



HOWEVER ...

Standalone CDM framework FAILS IN REPRODUCING MEASURED ROTATION CURVES for many spiral galaxies.

Universal CDM-only prediction: Navarro-Frenk-White profile.

Central density & scale radius from N-body are correlated! -> mass-concentration relation

-> " CORE VS CUSP " PROBLEM ... TO GET AN IDEA:





4. The diversity problem

Monthly Notices of the ROYAL ASTRONOMICAL SOCIETY MNRAS **452**, 3650–3665 (2015)

The unexpected diversity of dwarf galaxy^{*} rotation curves

*) not dSph!

Kyle A. Oman,¹* Julio F. Navarro,¹[†] Azadeh Fattahi,¹ Carlos S. Frenk,² Till Sawala,² Simon D. M. White,³ Richard Bower,² Robert A. Crain,⁴ Michelle Furlong,² Matthieu Schaller,² Joop Schaye⁵ and Tom Theuns²



BORROWED FROM A WAY-BETTER SPEAKER :)

Dwarf Spheroidal Galaxies (dSphs)



~10 (pre-SDSS) brightest ones (Classicals)

After SDSS and DES surveying the Sky, ~ **50** DM dominated MW satellites found.

Fairly close to Us (tens-100s kpc from Sun)

among the Way's Milley Way ites Milley satellites

No ongoing star formation (old stellar populations)

Faint stellar component (gas below exp sensitivity)

Large Mass-to-Light ratios $M/L \sim \mathcal{O}(10^{2-3}) M_{\odot}/L_{\odot}$

CLOSE TO IDEAL DM ASTRO-LABORATORY!

CDM-only predicts many more ... A "MISSING SATELLITES" (MS) PROBLEM ?! Moore et al. '99, Klypin et al. '99 ... but see, e.g., Kim et al.'17 for a recent criticism

Collisionless Boltzmann equation:

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla_{\vec{x}} f - \nabla_{\vec{x}} \phi \cdot \nabla_{\vec{v}} f = 0$$

- I) DYNAMICAL EQUILIBRIUM
- 2) SPHERICAL SYMMETRY

2nd MOMENT OF THE EQ. :

Evolution of phase space density of star in the galaxy, tracing the total gravitational potential.

$$(\rho_\star \sigma_r^2)' + 2\,\beta\,\rho_\star \sigma_r^2 = -\rho_\star \Phi_{\rm tot}'$$

where
$$' \equiv d/d\log r$$
 .

Collisionless Boltzmann equation:

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$$(\rho_\star \sigma_r^2)' + 2\beta \rho_\star \sigma_r^2 = -\rho_\star \Phi_{\rm tot}'$$

 $\rho_{\star}(r)$ is the **stellar density** of the system,

to be matched to photometric measurements —> connected to surface brightness, i.e.

$$\Sigma_{\star}(R) = \int_{R^2}^{\infty} \frac{dr^2}{\sqrt{r^2 - R^2}} \rho_{\star}$$

(hereafter : r = 3D physical radius, R = 2D projected radius)



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$$(\rho_\star \sigma_r^2)' + 2\beta \rho_\star \sigma_r^2 = -\rho_\star \Phi_{\rm tot}'$$

 $\sigma_{r(t)}(r)$ is the **radial (tangential)** component of the stellar velocity dispersion.

THE STELLAR **ORBITAL ANISOTROPY** IS DEFINED AS :

$$-\infty < \beta(r) \equiv 1 - \sigma_t^2 / \sigma_r^2 \le 1$$

circular limit

 $\beta = 0$: isotropic motion

radial orbits

Observational info on the stellar velocity dispersions only along line of sight (l.o.s.) from spectroscopy.

$$\Rightarrow \sigma_{los}^2(R) = \Sigma_{\star}^{-1} \int_{R^2}^{\infty} \frac{dr^2}{\sqrt{r^2 - R^2}} \left(1 - \beta \frac{R^2}{r^2}\right) \rho_{\star} \sigma_r^2$$

Collisionless Boltzmann equation:

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla_{\vec{x}} f - \nabla_{\vec{x}} \phi \cdot \nabla_{\vec{v}} f = 0$$

- I) DYNAMICAL EQUILIBRIUM
- 2) SPHERICAL SYMMETRY

2nd MOMENT OF THE EQ. :

$$\begin{array}{l} (\rho_{\star}\sigma_{r}^{2})' + 2\,\beta\,\rho_{\star}\sigma_{r}^{2} = -\rho_{\star}\Phi_{\mathrm{tot}}' \\ \underline{\mathbf{Obs}} \text{. tot} \sim \mathrm{DM} \end{array}$$

L.o.s. projection + Poisson's eq. leave us with 2 unknowns for 1 observable!



$$\Rightarrow \sigma_{los}(R) = f(\beta, \mathcal{M})(R)$$

DEGENERACY PROBLEM

In the spherical Jeans analysis, the total mass profile must be determined together with the orbital anisotropy function.

EXISTENCE OF A "MASS ESTIMATOR" I.E., APPROXIMATELY W/O DEGENERACY!

There is a *mass estimator* for dispersion-supported systems like MW dSphs.

FROM THE JEANS EQ.: $\mathcal{M}(r;\beta) - \mathcal{M}(r;\beta=0) \propto \gamma_{\rho_{\star}} + \gamma_{\sigma_r^2} + \gamma_{\beta} - 3$

Carina 10⁷ M(r) [M_©] 10⁶ r_{limit} 10^{5} 0.0 0.2 0.6 0.8 0.4 3D Physical Radius [kpc]

MNRAS 406 (2010) 1220, Wolf, J. et al.

Numerical evidence for mass difference to be 0 close by **3D half-light radius**, i.e. when density slope approaches 3.

where $\gamma_O \equiv -O'/O$.

 $\mathcal{M}(r_{1/2};\beta) \simeq \mathcal{M}(r_{1/2};\beta=0)$

$$\mathcal{M}_{1/2} \simeq 3 r_{1/2} \langle \sigma_{los}^2 \rangle / G_N$$

Similar formula, but w/o analytical insights, originally proposed in *Walker et al.* **'09**.

P.Ullio & M.V. JCAP 1607 (2016) 025 THEORY BIAS ON THE MASS ESTIMATOR



EXISTENCE OF A "MASS ESTIMATOR" I.E., APPROXIMATELY W/O DEGENERACY!

Recently, dedicated simulations for dSph-like systems confirm estimator.

Campbell et al.'16, Gonzalez-Samaniego et al.'17

... BUT OBSERVATIONAL ERRORS MAY NOT BE WELL UNDER CONTROL !!!

See, e.g., Fattahi, A. et al. '16

MW DWARFS INTERNAL DYNAMICS IN TENSION WITH CDM PREDICTIONS !

 $dSphV_{circ} @ r_{1/2}$ from stellar kinematics

$$\mathcal{M}(r_{1/2}) \Rightarrow V_{circ}(r_{1/2}) = \sqrt{3\langle \sigma_{los}^2 \rangle}$$

TOO-BIG-TO-FAIL (TBTF) PARADOX

Most massive subhalos in CDM seem to be too dense to host the observed brightest MW satellites.

On other hand, typically easier to trigger star formation in deep potential wells ...

... SO, WHERE ARE THEY?

M. Boylan-Kolchin, J.S. Bullock & M. Kaplinghat MNRAS 415 (2011) L40, MNRAS 422 (2012) 1203

- Dependence on the mass of host galaxy -> Wang, J. et al., MNRAS 424 (2012) 2715
- MW tidal field effects + Baryonic physics -> Sawala, T. et al., MNRAS 457 (2016) 1931

INNER CORES IN MW DWARFS PROFILES?

Chemo-distinct stellar populations can trace reliably the same gravity potential well, but @ different $r_{1/2}$.

-> MEASURE dSph MASS SLOPE!

ApJ 681 (2008) L13, Battaglia, G. et al. MNRAS 406 (2010) 1220, Amorisco & Evans ApJ 742 (2011) 20, Walker & Penarubbia Sculptor & Fornax host inner core

MNRAS 429 (2013) L89 , Amorisco, N. et al. Fornax core is large, of order kpc

ApJ 838 (2017) 123, Strigari + FW Generalized PSDF approach showing Sculptor may be fine also with NFW

Consensus not completely unanimous ANSWER MAY (LIKELY) BE "YES"!

CORE VS CUSP + DIVERSITY + TBTF : NEW PARADIGM BEYOND CDM?

Self-Interacting Dark Matter (SIDM), D.Spergel & P.Steinhardt '00

SELF-INTERACTIONS SHARED BY DM PARTICLES ALLOW FOR HEAT TRANSPORT. THIS ESTABLISHES A THERMALIZATION REGIME IN THE INNER GALACTIC HALO.

FORMATION OF INNER DENSITY CORES!

WARNING: upper-limit on self-scattering x-section per unit mass do exist ! Most recent merging cluster analyses typically agree on: $\sigma/m \lesssim cm^2/g$ Stellar kinematics of cluster brightest central galaxies gives even stronger limit. See. e.g., S.Tulin & H.B.Yu `17 for an extensive discussion!

Astro-Object Size <—> Different Collider Energy

Medium energies (v/c ~ 10⁻³)

High energies (v/c ~ 10⁻²)

PRL 116 (2016) 041302 , M.Kaplinghat, S.Tulin & H.B.Yu

Radius

SEMI-ANALYTIC SIDM HALO MODEL

$$\Gamma_{\text{scatt.}}\Big|_{r=r_1} \simeq t_{\text{age}}^{-1}, \ \Gamma_{\text{scatt.}} = \frac{\langle \sigma v \rangle}{m} \rho(r)$$

SIDM —> kinetic eq. for $r < r_1$. Therefore:

$$\nabla p = -\rho \, \nabla \phi_{tot} \ , \ p = \sigma_0^2 \, \rho \ .$$

ISOTHERMAL CORED PROFILE

$$\rho_{\text{SIDM}}(r) = \begin{cases} \rho_{\text{ISO}}(r) & \text{if } r \leq r_1 \\ \rho_{\text{NFW}}(r) & \text{if } r \geq r_1 \end{cases}$$

+ same matching on the mass profile.

Solving the Diversity Problem

More recent Ultra-faint dSphs not included already not-easy life with Classicals' data!

COLLIDERS @ kpc: SIDM IN MW dSphs M.V. & H.B.Yu, 1711.03502

Cauchy problem for SIDM profile in dSphs:

$$\frac{d^2h}{dr^2} = -\frac{2}{r}\frac{dh}{dr} - \frac{4\pi G_N \rho_0}{\sigma_0^2} \exp(h),$$

$$h \equiv \ln(\rho/\rho_0), \ h(0) = 1, \ h'(0) = 0.$$

NFW matching at r_I :
$$\begin{cases} \mathcal{M}_{\rm NFW}(r_1) = \mathcal{M}_{\rm ISO}(r_1) \\ \rho_{\rm NFW}(r_1) = \rho_{\rm ISO}(r_1). \end{cases}$$

$$\rho_{\star}(r) \propto \frac{3}{4R_{1/2}} \left(1 + (r/R_{1/2})^2\right)^{-5/2}$$

Plummer model provides good fit of available photometry (see, e.g., *Irwin & Hatzidimitriou '95*).

Orbital anisotropy from Baes & Van Hese '07,

$$\beta(r) = \frac{\beta_0 + \beta_\infty (r/r_\beta)^\eta}{1 + (r/r_\beta)^\eta}$$

M.V. & H.B.Yu, **1711.03502**

IN OUR ANALYSIS FURTHER CONSTRAINTS FROM CDM-ONLY N-BODY SIMULATIONS.

Matching conditions @ r_1 allows for:

$$\rho_0, \sigma_0) \leftrightarrow (V_{max}, R_{max})$$

 $R_{max} = 2.16 r_s$

 $V_{max} = 0.465 \sqrt{4\pi G_N \rho_n r_s^2}$

To get the "right cosmology" at large scales, we implement the CDM-only mass-concentration relation:

$$\log_{10}\left(\frac{R_{max}}{\text{kpc}}\right) = A + B \,\log_{10}\left(\frac{V_{max}}{km/s}\right)$$

$$A = 0.48, B = 0.18$$

$$\sigma_{\log_{10}(R_{max})} = 0.2$$

$$30 \lesssim V_{max} \lesssim 60$$

M.V. & H.B.Yu, **1711.03502**

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PRACTICALLY, USE OF GAUSSIAN WEIGHT:

$$\chi^2_{\rm CDM} = \frac{1}{2} \left(\frac{\log_{10}(R_{max}) - 0.48 - 0.18 \, \log_{10}(V_{max})}{0.2} \right)^2$$

WE ASSUME SUCH RELATION TO BE VALID UP TO HALO SIZE ~ (10 kpc)

+ PENALTY FACTORS FOR V_{MAX} RANGE :

 $V_{\rm circ}(r = 10 \,{\rm kpc}) < 25$, > 60 km/s $V_{\rm circ}^{\rm NFW}(r = 0.5 \,{\rm kpc}) < 19 \,{\rm km/s}$

Dot-Dashed: CDM fit (6 params)

Dashed: SIDM fit (4 params) Continuous: SIDM fit (7 params)

M.V. & H.B.Yu, **1711.03502**

Good fit of kinematic data while respecting cosmo-constraint!

<u>Obs</u>.

Theory condition from the stellar phase-space density non-negativity:

$$\gamma_{\rho_{\star}}\big|_{r\to 0} \ge 2\beta\big|_{r\to 0}$$

J.An & N.W.Evans ApJ 642 (2006) 752

M.V. & H.B.Yu, **1711.03502**

TBTF & SIDM FROM MCMC ANALYSIS

- NICE MATCH TO CDM COSMOLOGY
 ROUGHLY AT SCALES ≳ kpc.
- MATCH ENCLOSED MASS @ r_{1/2} WITHIN ~2 SIGMA LEVEL
- MASS ESTIMATOR SHOWN WITH QUITE "CONSERVATIVE" ERRORS
- FIT TO FULL DATASET USEFUL: MASS ESTIMATOR INDUCE SOME BIAS, DUE TO ANISOTROPY!

M.V. & H.B.Yu, **1711.03502 ABOUT SIDM X-SEC:** A DATA-DRIVEN ESTIMATE $\frac{\langle \sigma v \rangle}{m} \simeq \frac{1}{\rho(r_1) t_{aae}} \Rightarrow \frac{\sigma}{m} \simeq \frac{\sqrt{\pi}}{4\sigma_0} \frac{1}{\rho(r_1) t_{aae}} ,$ under (motivated) Maxwellian approximation, i.e.: $\langle v \rangle \simeq 4\sigma_0 / \sqrt{\pi} , \ \langle \sigma v \rangle \simeq \sigma \langle v \rangle .$ **FORNAX** 00 5 $? \lesssim \sigma/m \lesssim ?$ $\langle v \rangle ~ [{
m km~s^{-1}}]$ 60 NS NS Marginalization over the age of the system using the flat prior: 20 $8 \lesssim t_{\rm age} \, [{\rm Gyr}] \lesssim 12$ $\sigma/m ~[\mathrm{cm^2~g^{-1}}]$ 80 A. W. McConnachie ApJ 144 (2012) 4 ⁰0 estimate of SIDM 0 2500 2000 1500 5000 Ň 60 15 ω_{o} 20 % 0 20 x-sec from data! $\langle v
angle ~ [{
m km~s^{-1}}] ~ \sigma/m ~ [{
m cm^2~g^{-1}}]$ $\left<\sigma v\right>/m~[{
m cm}^3~{
m g}^{-1}]$

COLLIDERS @ kpc: SIDM IN MW dSphs M.V. & H.B.Yu, 1711.03502

X-SEC PROBED BY MW DWARF SATELLITES SPANS 2 ORDERS OF MAGNITUDE !

HOWEVER, 6 OUT OF 8 GALAXIES CAN BE CONSISTENT @ 68% WITH 0.5 – 3 cm² g⁻¹.

Same SIDM ballpark to address "Core VS Cusp" in other kpc-sized systems.
 Kaplinghat, M. et al. '16, Kamada, A. et al. '17

III) Two main "outliers" pointing to possible importance of environmental effects.

-> SIDM ameliorates TBTF problem! Vogelsberger, M. et al. '16 (ETHOS)

M.V. & H.B.Yu, **1711.03502**

Our study on SIDM halo in MW dSphs shows:

X-sec range in agreement with current indications from N-body simulations.
 Zavala, J. et al. '13, Elbert, O. et al. '15

N.B. We did not account for core collapse ...

NEXT STEPS: A TBTF PROBLEM IN THE FIELD!

dSph FIELD GALAXIES MAY REPRESENT EXTREMELY IMPORTANT TESTS!

SIDM SIMULATIONS FOR FIELD DWARFS CONFIRM TREND SEEN FOR MW ONES ... -> see, for instance, **Elbert et al. '14** ... DEEPER INVESTIGATION NEEDED!

Measured kinematics for field galaxies near MW & M31, in the Local Group, point once again to less massive subhalos than in CDM studies.

Garrison-Kimmel et al. '14, Papastergis et al. '14

"field galaxies" should not be influenced by (otherwise possibly relevant) external tides!

NEXT STEPS: MODEL-BUILDING NEED GUIDELINES...

Strongly Interacting Massive Particles

PRL 113 (2014) 171301, Hochberg,Y. et al. **PRL 115 (2015) 021301,** Hochberg,Y. et al.

Self-Interactions with Light Mediators

PRL 104 (2009) 151301, J. Feng et al.
PRD 81 (2010) 083522, M.R.Buckley & P.J.Fox
PRL 106 (2011) 171303, A.Loeb & N.Weiner
PRL 110 (2013) 111301, S.Tulin et al.

@ strong coupling, strong scale emerges:

$$m_{DM} \sim \alpha_{eff} (T_{eq}^2 M_{Pl})^{1/3}$$

"Simple" realizations involve non-Abelian dark sector with QCD-like chiral symmetry breaking

Dominant 3, 4 —> 2 annihilations, dark sector cannot be completely secluded from SM

ApJ 398 (1992) 43 , E.D. Carlson, M. E. Machacek & L.J.Hall

Large self-scattering points to MeV mediators for weak-scale DM (perturbative regime):

$$g^4 \frac{m_{\chi}^2}{m_{\phi}^4} \sim 10^{14} \, \frac{\alpha_{EW}^2}{m_{\chi}^2} \ \Rightarrow \ \frac{m_{\phi}}{m_{\chi}} \sim \left(\frac{g}{0.1}\right)^4 \, 10^{-4}$$

 \longrightarrow U(1)_D coupled to SM through U(1)_Y small mixing

LIGHT MEDIATOR MODELS ALLOW FOR DM V-DEPENDENT SELF-SCATTERING X-SEC!

NEXT STEPS: MODEL-BUILDING NEED GUIDELINES...

Self-Interactions with Light Mediators

PRL 104 (2009) 151301, J. Feng et al.
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LIGHT MEDIATOR MODELS ALLOW FOR DM V-DEPENDENT SELF-SCATTERING X-SEC!

... ESPECIALLY IN LIGHT OF PRESENT CONSTRAINTS ...

... BUT "NIGHTMARE SCENARIOS" MAY BE (STILL) INTERESTING!

 $\mathcal{L}_{\rm int} = -ig_{\chi}\bar{\chi}\gamma^{\mu}\chi\phi_{\mu} + m_{\chi}\bar{\chi}\chi + \frac{1}{2}m_{\phi}^{2}\phi^{\mu}\phi_{\mu} - ig_{f}\bar{f}\gamma^{\mu}f\phi_{\mu}$

I.E. ADD DARK RADIATION SPECIES (E.G. STERILE V)

LATE KINETIC DECOUPL. —> **TBTF** + **MS IN 1 SHOT!** van den Aarsen et al., PRL '12, Bringmann et al., PRD '14

UPDATED LYMAN- α CONSTRAINTS **MAY POINT TO T_{KD} > 1 keV** ... IF SO, NOT GOOD NEWS FOR MS!

"NIGHTMARE SCENARIO" STILL PROVIDES X-SEC GOOD FOR TBTF PROBLEM!

- —> INTERESTING SIGNATURES FOR NEXT-GEN COSMO-SURVEYS
- --> NEED MORE PRECISE SCRUTINY ON ITS GALACTIC PREDICTIONS

NEXT STEPS: SIDM VS NP SMALL-SCALE ALTERNATIVES

Few N-body simulation studies so far ... H.-Y.Schive et al. '14 , P.Mocz et al. '17

NOT ONLY BOSONS: Degenerate Fermi gas

V. Domcke & A.Urbano '15, L.Randall, J.Scholtz & J.Unwin '17

$$f_{\rm FD} = \begin{cases} 1 & p \leqslant p_{\rm F} \\ 0 & p > p_{\rm F} \end{cases}$$

Fermi quantum pressure can balance gravity, yielding cored density profiles.

PERSONAL RECAP

The SIDM proposal offers a very compelling DM paradigm:

- TO SOLVE COREVS CUSP (LIKELY PRESENT ALSO IN DSPHS)
- TO ADDRESS (ELEGANTLY) THE DIVERSITY PROBLEM

(TO LOOSE MEMORY ON BARYONIC FEEDBACK !!!)

- TO PARTIALLY ADDRESS ''TBTF PROBLEMS'' AS THE ONE SHOWN BY THE INTERNAL DYNAMICS IN MW DSPHS
- -> NEW MODEL-BUILDING AVENUES TO BE FURTHER EXPLORED IN RELATION TO COREVS CUSP, DIVERSITY, TBTF, MS & "NIGHTMARE SCENARIOS"
- -> CLOSE COMPARISON WITH OTHER SMALL-SCALE NP SOLUTIONS MISSING !

Tying SIDM to Baryons

• SIDM may follow the stellar distribution; halo morphology

- The SIDM distribution is sensitive to the final baryon distribution
- But, it is not sensitive to the formation history

credit to Hai-Bo Yu, UC Riverside

Strong Feedback vs. SIDM

Vogelsberger, M. et al., MNRAS 460 (2016) 1399

Name	$lpha_{\chi}$	$lpha_{ u}$	$m_{\phi} \ [{ m MeV}{ m c}^{-2}]$	$m_{\chi} \ [{ m GeV}{ m c}^{-2}]$	$r_{ m DAO} \ [h^{-1}{ m Mpc}]$	$r_{ m SD} \ [h^{-1}{ m Mpc}]$	$a_4 \ [h{ m Mpc}^{-1}]$	$\langle \sigma_T angle_{30}/m_\chi \ [m cm^2g^{-1}]$	$\langle \sigma_T angle_{200}/m_\chi \ [m cm^2g^{-1}]$	$\langle \sigma_T angle_{1000}/m_\chi \ [m cm^2g^{-1}]$
CDM	_	_	_	_	_	_	_	_	_	
ETHOS-1	0.071	0.123	0.723	2000	0.362	0.225	14095.65	4.98	0.072	0.0030
ETHOS-2	0.016	0.03	0.83	500	0.217	0.113	1784.05	9.0	0.197	0.00097
ETHOS-3	0.006	0.018	1.15	178	0.141	0.063	305.94	16.9	0.48	0.0028
ETHOS-4 (tuned)	0.5	1.5	5.0	3700	0.138	0.0615	286.09	0.16	0.022	0.00075
CDM:	has MS	and TB'	TF problems	ETHOS-1	:	over-solves MS	S and TBTF pro	blems ETHOS	S-2: over-solves	TBTF problem

ETHOS-4 (tuned):

alleviates MS and TBTF problems

