

The Price of Tiny Kinetic Mixing

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Based on Tony Gherghetta, JK, Keith Olive, and Maxim Pospelov,
[arXiv:1909.00696](https://arxiv.org/abs/1909.00696)

1 Introduction

2 Bottom-Up Models

3 Top-Down Models

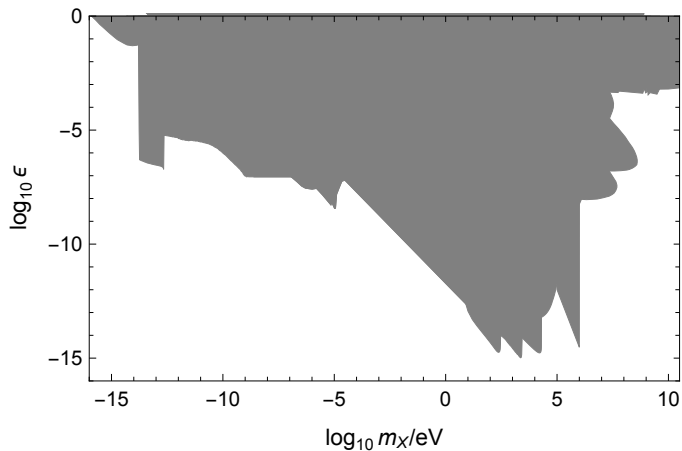
Dark Photons with Kinetic Mixing

- No new physics at LHC \rightsquigarrow hiding at **low energies**?
- One candidate: gauge boson X^μ of new $U(1)_X$ (**dark photon**)
- Mass m_X from Brout-Englert-Higgs or Stückelberg mechanism
- Applications: dark matter candidate, mediator of dark matter self-interactions, $(g-2)_\mu, \dots$
- Possibly part of **dark sector**
- Simplest way to couple to Standard Model: **kinetic mixing** with $U(1)_Y$ gauge boson B^μ

$$\mathcal{L}_{\text{km}} = -\frac{1}{2} \epsilon B_{\mu\nu} X^{\mu\nu}$$

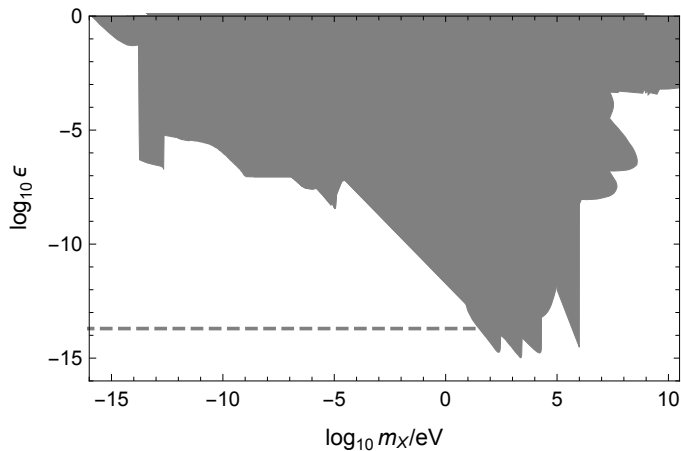
\rightsquigarrow kinetic mixing with photon ($\epsilon \cos \theta_W$) and Z ($\epsilon \sin \theta_W$)

Constraints on Kinetic Mixing



Redondo, personal communication; Beacham et al., [arXiv:1901.09966](https://arxiv.org/abs/1901.09966)

Constraints on Kinetic Mixing

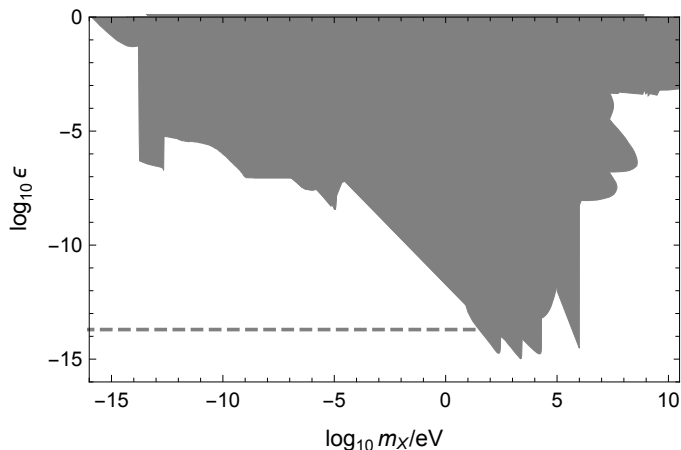


Redondo, personal communication; Beacham et al., [arXiv:1901.09966](https://arxiv.org/abs/1901.09966)

Dark photon mass from dark Higgs: $g_X \epsilon \lesssim 10^{-14}$ for $m_X \lesssim 10 \text{ keV}$

Ahlers, Jaeckel, Redondo, Ringwald, PRD **78** (2008)

Constraints on Kinetic Mixing



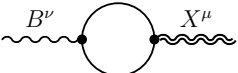
Redondo, personal communication; Beacham et al., arXiv:1901.09966

$\rightsquigarrow \epsilon \lesssim 10^{-15} \dots 10^{-7}$ for $\mu\text{eV} \lesssim m_X \lesssim 10 \text{ MeV}$

\rightsquigarrow How to get such a small number from a model?

Generic One-Loop Expectation

- Vanishing kinetic mixing at high scale
- Field with mass M charged under $U(1)_X$ and $U(1)_Y$


$$\rightsquigarrow \epsilon \sim \frac{g' g_X}{8\pi^2} \ln \frac{M}{\mu} \sim 10^{-3} \dots 10^{-1} \gg 10^{-7}$$

Holdom, PLB 166 (1986)

- Non-decoupling effect \rightsquigarrow large M does not help

Known Suppression Mechanisms

- **Fine-tuned cancellation** between tree level and loop correction
- $g_x \ll 1$ from LARGE volume string compactifications
Burgess et al., JHEP **07** (2008)
Cicoli, Goodsell, Jaeckel, Ringwald, JHEP **07** (2011)
- Cancellation due to **mass degeneracy**
 - SUSY, string theory Dienes, Kolda, March-Russell, NPB **492** (1997)
 - GUTs Arkani-Hamed, Weiner, JHEP **12** (2008)
- **Type-II string theory** with warped extra dimensions and fluxes
Abel et al., JHEP **07** (2008)
- **4-loop suppression** in mirror-symmetric **twin Higgs** model
Chacko, Goh, Harnik, PRL **96** (2006)
Koren, McGehee, arXiv:1908.03559
- 2 additional options to be mentioned later

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↪ Additional ways to get $\epsilon \lesssim 10^{-7}$ for $g_X \sim g' \sim 1$?

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Higher Loop Order

- Additional gauge group $U(1)_M$, broken above TeV-scale
- Heavy fermions ψ and χ

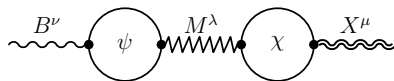
	Charge		
	$U(1)_Y$	$U(1)_M$	$U(1)_X$
ψ	1	1	0
χ	0	1	1

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\rightsquigarrow Kinetic mixing at 2-loop order?



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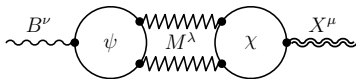
↪ Kinetic mixing at 2-loop order?

$$\sim \Pi_{YM}^{\nu\rho}(k^2) D_M^{\rho\sigma} \Pi_{MX}^{\sigma\mu}(k^2) \sim \frac{k^4}{m_M^2}$$

- ↪ Operator with **derivatives** of $B^{\mu\nu}$ and $X^{\mu\nu}$
- ↪ Does **not contribute** to kinetic mixing

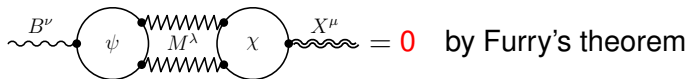
Higher Loop Order

↪ Kinetic mixing at 3-loop order?



Higher Loop Order

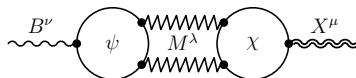
↪ Kinetic mixing at 3-loop order?



$= 0$ by Furry's theorem

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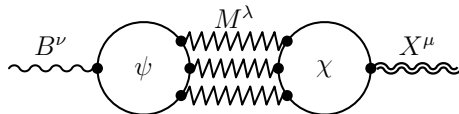
↪ Kinetic mixing at 3-loop order?



A Feynman diagram showing two fermion loops, labeled ψ and χ , connected by two internal mass insertions M^λ . The ψ loop is on the left, connected to an external wavy line B^ν . The χ loop is on the right, connected to an external wavy line X^μ . The two mass insertions M^λ connect the two loops. The diagram is followed by an equals sign and a red zero, and the text "by Furry's theorem".

$$= 0 \quad \text{by Furry's theorem}$$

↪ Kinetic mixing at 4-loop order

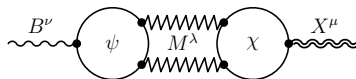


A Feynman diagram showing two fermion loops, labeled ψ and χ , connected by four internal mass insertions M^λ . The ψ loop is on the left, connected to an external wavy line B^ν . The χ loop is on the right, connected to an external wavy line X^μ . The four mass insertions M^λ connect the two loops. The diagram is followed by an approximation symbol and the expression $\epsilon \sim \left(\frac{1}{16\pi^2}\right)^4 \sim 10^{-9}$.

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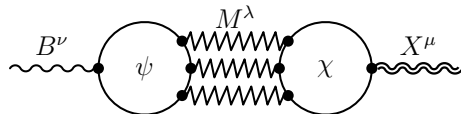
Higher Loop Order

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↪ Kinetic mixing at 4-loop order



$\rightsquigarrow \epsilon \sim \left(\frac{1}{16\pi^2} \right)^4 \sim 10^{-9}$

Similar model with $U(1)_M \rightarrow SU(3)_c$: Dunsky, Hall, Harigaya, JHEP 07 (2019)

More Gauge Groups

Gauged Clockwork

Giudice, McCullough, JHEP 02 (2017); Lee, PLB 778 (2018)

- $N + 1$ gauge symmetries $U(1)_i$, $i = 0, \dots, N$
- Equal gauge coupling g
- Corresponding gauge fields A_μ^i
- N Higgs fields ϕ_j , $j = 0, \dots, N - 1$, each with charges $(1, -q)$ under $U(1)_j \times U(1)_{j+1}$ (and charge 0 under the other groups)
- $\langle \phi_j \rangle = f$ for all $j \rightsquigarrow U(1)^{N+1} \rightarrow U(1)_X$

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 \rightsquigarrow zero mode = dark photon = linear combination of all A_μ^i
- Field charged only under $U(1)_N$:
coupling to dark photon $g_{\text{eff}} \sim \frac{g}{q^N}$ exponentially suppressed

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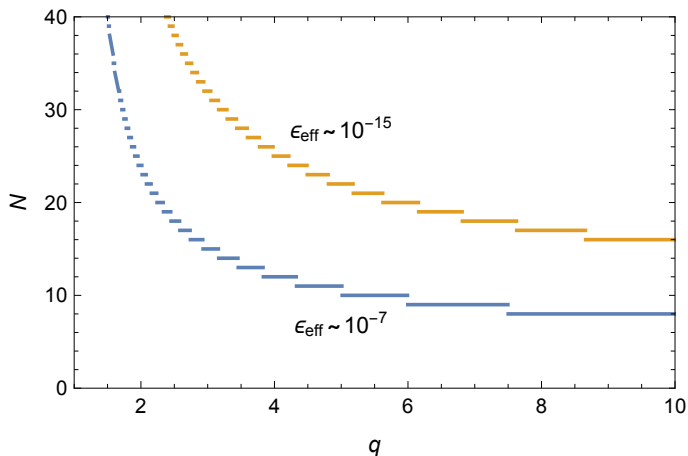
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- Continuum limit $N \rightarrow \infty$: equivalent to 5D theory, $g_{\text{eff}} \sim e^{-kR}$

Clockwork-Suppressed Kinetic Mixing

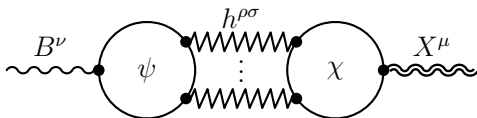
Kinetic mixing of B_μ only with A_μ^N

\rightsquigarrow mixing with dark photon $\epsilon_{\text{eff}} \sim \frac{\epsilon}{q^N}$ can be **tiny** even for $\epsilon \sim 1$



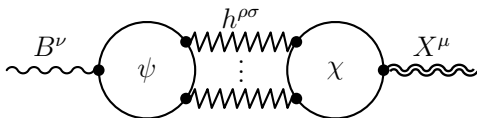
Gravity-Mediated Kinetic Mixing

- Visible and dark sector only connected via gravity
- **No** separate **C** conservation in dark sector
- Divergent graviton loops $\rightsquigarrow \epsilon \sim \left(\frac{\Lambda}{M_{\text{Pl}}}\right)^n$ **unsuppressed** if $\Lambda \sim M_{\text{Pl}}$



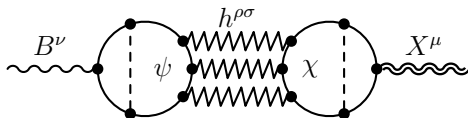
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- $\sum Y_i = \sum Y_i^3 = 0 \rightsquigarrow$ need non-universal couplings in loops



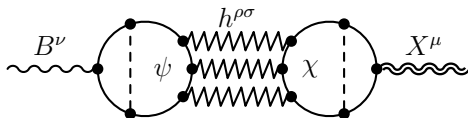
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$$\rightsquigarrow \epsilon \sim \left(\frac{1}{16\pi^2}\right)^6 \left(\frac{\Lambda}{M_{\text{Pl}}}\right)^6 \sim 10^{-13}$$

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Embedding Light and Dark Photons in a Single Group

- $G_{\text{SM}} \times U(1)_X \subset$ non-Abelian group G
 - $\rightsquigarrow \epsilon = 0$ at breaking scale of G
 - \rightsquigarrow Stückelberg mass excluded $\rightsquigarrow g_X \epsilon \lesssim 10^{-14}$ or $m_X \gtrsim 10 \text{ keV}$

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- **Light particles** charged under both $U(1)_Y$ and $U(1)_X$
 - \rightsquigarrow generic $\epsilon \sim 10^{-3} \dots 10^{-1}$
- **Heavy particles** charged under both $U(1)_Y$ and $U(1)_X$ fill complete GUT multiplets
 - $\rightsquigarrow \epsilon = 0$ for exact mass degeneracy
 - \rightsquigarrow Expect $\epsilon \sim 10^{-6} \dots 10^{-4}$
 - Arkani-Hamed, Weiner, JHEP 12 (2008)
 - \rightsquigarrow Can be ok for $m_X \gtrsim 1 \text{ MeV}$

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 - \rightsquigarrow Can be ok for $m_X \gtrsim 1$ MeV
- **Light particles** charged under both $U(1)$'s **hard to avoid**
 - Example:** $SO(10) \rightarrow SU(5) \times U(1)_X$
 - Standard Model fields in $\mathbf{16} = (\bar{\mathbf{5}}, \mathbf{3}) + (\mathbf{10}, -\mathbf{1}) + (\mathbf{1}, -\mathbf{5})$
- ... but **not impossible** for sufficiently large groups
 - Example:** $E_8 \rightarrow E_6 \times SU(3) \rightarrow E_6 \times U(1)_X$

Only Standard Model Embedded in Simple Group

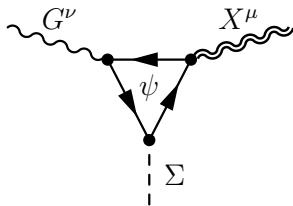
- $G_{\text{SM}} \times U(1)_X \subset G_{\text{GUT}} \times U(1)_X \rightsquigarrow G^{\mu\nu} X_{\mu\nu}$ not gauge-invariant
- **Effective operator** $\frac{1}{\Lambda} \Sigma G^{\mu\nu} X_{\mu\nu}$ with scalar Σ can be gauge-invariant
 $\rightsquigarrow \epsilon \sim \frac{\langle \Sigma \rangle}{\Lambda} \ll 1$

Arkani-Hamed, Weiner, JHEP 12 (2008)

- Generated via loops with heavy particles (mass Λ)
- Alternative: embed $U(1)_X$ in non-Abelian group

Example 1: Adjoint Scalar

- Scalar Σ in **adjoint** representation of **GUT group**, no $U(1)_X$ charge
- Vector-like fermion charged under **both** groups, mass Λ

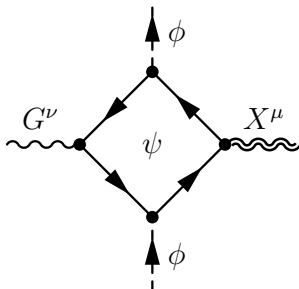


$$\rightsquigarrow \epsilon \sim \frac{y_\Sigma}{16\pi^2} \frac{\langle \Sigma \rangle}{\Lambda} \gtrsim \frac{y_\Sigma}{16\pi^2} \frac{M_{\text{GUT}}}{M_{\text{Pl}}} \sim 10^{-4} y_\Sigma$$

\rightsquigarrow Ok for **moderately small** Yukawa coupling y_Σ

Example 2: Non-Adjoint Scalar in $SU(5)$

- Scalar $\phi \sim (75, 0)$ under $SU(5) \times U(1)_X$
- Vector-like fermion $\psi = \chi \sim (10, q)$, mass Λ

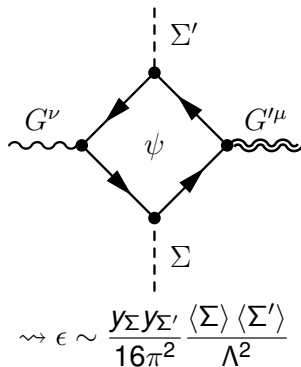


$$\rightsquigarrow \epsilon \sim \frac{y_\phi^2}{16\pi^2} \frac{\langle \phi \rangle^2}{\Lambda^2} \sim \frac{y_\phi^2}{16\pi^2} \left(\frac{M_{\text{GUT}}}{M_{\text{Pl}}} \right)^2 \sim 10^{-6} y_\phi^2$$

\rightsquigarrow Ok

Non-Abelian Groups in Both Sectors

- $G_{\text{SM}} \times U(1)_X \subset G_{\text{GUT}} \times G'$
- Adjoint scalars Σ, Σ'



- $\langle \Sigma' \rangle \ll M_{\text{GUT}} \rightsquigarrow$ tiny ϵ

Goldberg, Hall, PLB 174 (1986)

Conclusions

Kinetic mixing between visible and **dark photon** severely constrained:

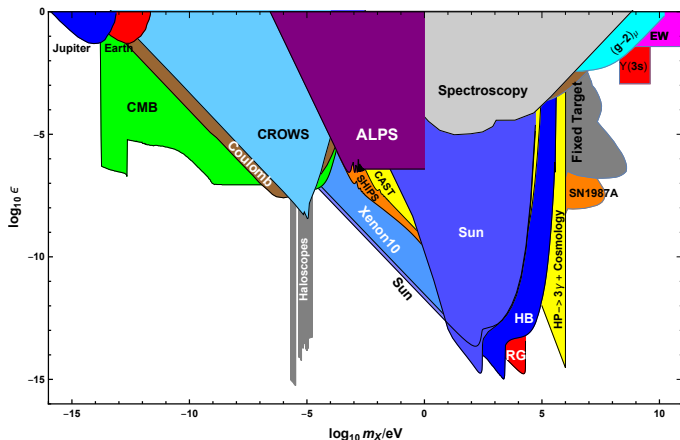
$$\epsilon \lesssim 10^{-15} \dots 10^{-7}$$

↪ **Scenarios to explain such a small mixing**

- Fine-tuning
- String theory
- Generation at high loop order $\rightsquigarrow \epsilon \sim 10^{-13} \dots 10^{-9}$
- Suppression by gauged clockwork
- Gravity mediation $\rightsquigarrow \epsilon \sim 10^{-13}$
- Embedding of both $U(1)$'s in common group $\rightsquigarrow \epsilon \sim 10^{-6} \dots 10^{-4}$
- Effective operators in GUT $\rightsquigarrow \epsilon \sim 10^{-28} \dots 10^{-4}$

↪ **Tiny kinetic mixing possible, but not for free**

Constraints on Kinetic Mixing



Redondo, personal communication (2018)

Collider limits (BaBar etc.) not included