

Equilibrium states on right LCM semigroup C*-algebras revisited

Nadia S. Larsen

University of Oslo

4 December 2017

joint with N. Brownlowe, J. Ramagge and N. Stammeier

The KMS condition for finite systems

Finite quantum systems: a time evolution on $M_n(\mathbb{C})$ is given by a one-parameter group of automorphisms

$$\sigma_t(a) = e^{itH} a e^{-itH},$$

where $t \in \mathbb{R}$, $a \in M_n(\mathbb{C})$ and H is a self-adjoint matrix.

The KMS condition for finite systems

Finite quantum systems: a time evolution on $M_n(\mathbb{C})$ is given by a one-parameter group of automorphisms

$$\sigma_t(a) = e^{itH} a e^{-itH},$$

where $t \in \mathbb{R}$, $a \in M_n(\mathbb{C})$ and H is a self-adjoint matrix.

The *Gibbs state* at $\beta > 0$ is $\varphi_G(a) = \frac{\text{Tr}(a e^{-\beta H})}{\text{Tr}(e^{-\beta H})}$. It satisfies

$$\varphi_G(ab) = \varphi_G(b \sigma_{i\beta}(a)), \quad (1)$$

for $a, b \in M_n(\mathbb{C})$ analytic, i.e. $t \mapsto \sigma_t(a)$ extends to an entire function on \mathbb{C} .

The KMS condition for finite systems

Finite quantum systems: a time evolution on $M_n(\mathbb{C})$ is given by a one-parameter group of automorphisms

$$\sigma_t(a) = e^{itH} a e^{-itH},$$

where $t \in \mathbb{R}$, $a \in M_n(\mathbb{C})$ and H is a self-adjoint matrix.

The *Gibbs state* at $\beta > 0$ is $\varphi_G(a) = \frac{\text{Tr}(a e^{-\beta H})}{\text{Tr}(e^{-\beta H})}$. It satisfies

$$\varphi_G(ab) = \varphi_G(b \sigma_{i\beta}(a)), \quad (1)$$

for $a, b \in M_n(\mathbb{C})$ analytic, i.e. $t \mapsto \sigma_t(a)$ extends to an entire function on \mathbb{C} .

Partition function of $(M_n(\mathbb{C}), \sigma)$ is $\beta \mapsto \text{Tr}(e^{-\beta H})$.

The KMS condition for finite systems

Finite quantum systems: a time evolution on $M_n(\mathbb{C})$ is given by a one-parameter group of automorphisms

$$\sigma_t(a) = e^{itH} a e^{-itH},$$

where $t \in \mathbb{R}$, $a \in M_n(\mathbb{C})$ and H is a self-adjoint matrix.

The *Gibbs state* at $\beta > 0$ is $\varphi_G(a) = \frac{\text{Tr}(ae^{-\beta H})}{\text{Tr}(e^{-\beta H})}$. It satisfies

$$\varphi_G(ab) = \varphi_G(b\sigma_{i\beta}(a)), \quad (1)$$

for $a, b \in M_n(\mathbb{C})$ analytic, i.e. $t \mapsto \sigma_t(a)$ extends to an entire function on \mathbb{C} .

Partition function of $(M_n(\mathbb{C}), \sigma)$ is $\beta \mapsto \text{Tr}(e^{-\beta H})$.

(1) - the *KMS condition*, cf. Haag-Hugenholtz-Winnick (1967): equilibrium for a state on a C*-algebra with time evolution.

KMS states

By analogy with finite systems and the Gibbs state, extend the notions of KMS_β state, partition function, inverse temperature.

KMS states

By analogy with finite systems and the Gibbs state, extend the notions of KMS_β state, partition function, inverse temperature.

\mathcal{A} C*-algebra, $\sigma : \mathbb{R} \rightarrow \text{Aut}(\mathcal{A})$ time evolution, φ a state on \mathcal{A} .

① φ is KMS_β (at inverse temperature $\beta \in [0, \infty)$) if

$$\varphi(ab) = \varphi(b\sigma_{i\beta}(a))$$

for all $a, b \in \mathcal{A}^a$, the dense *-subalgebra of *analytic* elements.

KMS states

By analogy with finite systems and the Gibbs state, extend the notions of KMS_β state, partition function, inverse temperature.

\mathcal{A} C*-algebra, $\sigma : \mathbb{R} \rightarrow \text{Aut}(\mathcal{A})$ time evolution, φ a state on \mathcal{A} .

- ① φ is KMS_β (at inverse temperature $\beta \in [0, \infty)$) if

$$\varphi(ab) = \varphi(b\sigma_{i\beta}(a))$$

for all $a, b \in \mathcal{A}^a$, the dense *-subalgebra of *analytic* elements.

- ② A state φ is a *ground state* if for all a, b with b analytic, the function $z \rightarrow \varphi(a\sigma_z(b))$ is bounded in the upper-half plane.

KMS states

By analogy with finite systems and the Gibbs state, extend the notions of KMS_β state, partition function, inverse temperature.

\mathcal{A} C*-algebra, $\sigma : \mathbb{R} \rightarrow \text{Aut}(\mathcal{A})$ time evolution, φ a state on \mathcal{A} .

- 1 φ is KMS_β (at inverse temperature $\beta \in [0, \infty)$) if

$$\varphi(ab) = \varphi(b\sigma_{i\beta}(a))$$

for all $a, b \in \mathcal{A}^a$, the dense *-subalgebra of *analytic* elements.

- 2 A state φ is a *ground state* if for all a, b with b analytic, the function $z \rightarrow \varphi(a\sigma_z(b))$ is bounded in the upper-half plane.
- 3 KMS_∞ if $\varphi = w^* \lim \varphi_n$ as $\beta_n \rightarrow \infty$ and φ_n is KMS_{β_n} .

KMS states

By analogy with finite systems and the Gibbs state, extend the notions of KMS_β state, partition function, inverse temperature.

\mathcal{A} C*-algebra, $\sigma : \mathbb{R} \rightarrow \text{Aut}(\mathcal{A})$ time evolution, φ a state on \mathcal{A} .

- 1 φ is KMS_β (at inverse temperature $\beta \in [0, \infty)$) if

$$\varphi(ab) = \varphi(b\sigma_{i\beta}(a))$$

for all $a, b \in \mathcal{A}^a$, the dense *-subalgebra of *analytic* elements.

- 2 A state φ is a *ground state* if for all a, b with b analytic, the function $z \rightarrow \varphi(a\sigma_z(b))$ is bounded in the upper-half plane.
- 3 KMS_∞ if $\varphi = w^* \lim \varphi_n$ as $\beta_n \rightarrow \infty$ and φ_n is KMS_{β_n} .

References: Bratteli-Robinson, Pedersen, Connes-Marcolli.

KMS_∞ strictly subset of ground states

Theorem (Laca-Raeburn (2010))

$C^*(\mathbb{N} \rtimes \mathbb{N}^\times)$ is the universal C^* -algebra generated by isometries s and $\{v_p \mid p \text{ prime}\}$, subject to the relations

- 1 $v_p s = s^p v_p$;
- 2 $v_p v_q = v_q v_p$,
- 3 $v_p^* v_q = v_q v_p^*$ when $p \neq q$,
- 4 $s^* v_p = s^{p-1} v_p s^*$, and
- 5 $v_p^* s^k v_p = 0$ for $1 \leq k < p$.

Dynamics: $\sigma_t(s) = s$ and $\sigma_t(v_p) = p^{it} v_p$.

KMS_∞ strictly subset of ground states

Theorem (Laca-Raeburn (2010))

$C^*(\mathbb{N} \rtimes \mathbb{N}^\times)$ is the universal C*-algebra generated by isometries s and $\{v_p \mid p \text{ prime}\}$, subject to the relations

- 1 $v_p s = s^p v_p$;
- 2 $v_p v_q = v_q v_p$,
- 3 $v_p^* v_q = v_q v_p^*$ when $p \neq q$,
- 4 $s^* v_p = s^{p-1} v_p s^*$, and
- 5 $v_p^* s^k v_p = 0$ for $1 \leq k < p$.

Dynamics: $\sigma_t(s) = s$ and $\sigma_t(v_p) = p^{it} v_p$. Then, for $\beta < 1$, there are no KMS states, if $\beta \in [1, 2]$, there is a unique KMS_β state;

KMS $_{\infty}$ strictly subset of ground states

Theorem (Laca-Raeburn (2010))

$C^*(\mathbb{N} \rtimes \mathbb{N}^{\times})$ is the universal C*-algebra generated by isometries s and $\{v_p \mid p \text{ prime}\}$, subject to the relations

- 1 $v_p s = s^p v_p$;
- 2 $v_p v_q = v_q v_p$,
- 3 $v_p^* v_q = v_q v_p^*$ when $p \neq q$,
- 4 $s^* v_p = s^{p-1} v_p s^*$, and
- 5 $v_p^* s^k v_p = 0$ for $1 \leq k < p$.

Dynamics: $\sigma_t(s) = s$ and $\sigma_t(v_p) = p^{it} v_p$. Then, for $\beta < 1$, there are no KMS states, if $\beta \in [1, 2]$, there is a unique KMS $_{\beta}$ state; if $\beta \in (2, \infty]$, the KMS $_{\beta}$ states are parametrised by probability measures on \mathbb{T} while the ground states are parametrised by states on the Toeplitz C*-algebra generated by a single isometry.

KMS_∞ strictly subset of ground states

Theorem (Afsar-Brownlowe-L-Stammeier (2016))

Let S be a right LCM monoid and $N : S \rightarrow \mathbb{N}^\times$ homomorphism such that S is *admissible*. Consider the time evolution $\sigma_t(v_s) = N_s^{it} v_s$. If $\beta_c \in \mathbb{R}$ is such that the function

$$\zeta_N(\beta) := \sum_{n \in \text{Irr}(N(S))} n^{-(\beta-1)},$$

converges for $\beta \geq \beta_c$, then for $(C^*(S), \mathbb{R}, \sigma)$ we have

KMS_∞ strictly subset of ground states

Theorem (Afsar-Brownlowe-L-Stammeier (2016))

Let S be a right LCM monoid and $N : S \rightarrow \mathbb{N}^\times$ homomorphism such that S is *admissible*. Consider the time evolution $\sigma_t(v_s) = N_s^{it} v_s$. If $\beta_c \in \mathbb{R}$ is such that the function

$$\zeta_N(\beta) := \sum_{n \in \text{Irr}(N(S))} n^{-(\beta-1)},$$

converges for $\beta \geq \beta_c$, then for $(C^*(S), \mathbb{R}, \sigma)$ we have

- 1 $\beta \in [0, 1)$: no KMS_β state;

KMS_∞ strictly subset of ground states

Theorem (Afsar-Brownlowe-L-Stammeier (2016))

Let S be a right LCM monoid and $N : S \rightarrow \mathbb{N}^\times$ homomorphism such that S is *admissible*. Consider the time evolution $\sigma_t(v_s) = N_s^{it} v_s$. If $\beta_c \in \mathbb{R}$ is such that the function

$$\zeta_N(\beta) := \sum_{n \in \text{Irr}(N(S))} n^{-(\beta-1)},$$

converges for $\beta \geq \beta_c$, then for $(C^*(S), \mathbb{R}, \sigma)$ we have

- 1 $\beta \in [0, 1)$: no KMS_β state;
- 2 $\beta \in [1, \beta_c]$: unique KMS_β if action $S_c \curvearrowright S/S_c$ essentially free, where $S_c \subset S$ subsemigroup of elements having LCM with any t in S ;

KMS_∞ strictly subset of ground states

Theorem (Afsar-Brownlowe-L-Stammeier (2016))

Let S be a right LCM monoid and $N : S \rightarrow \mathbb{N}^\times$ homomorphism such that S is *admissible*. Consider the time evolution $\sigma_t(v_s) = N_s^{it} v_s$. If $\beta_c \in \mathbb{R}$ is such that the function

$$\zeta_N(\beta) := \sum_{n \in \text{Irr}(N(S))} n^{-(\beta-1)},$$

converges for $\beta \geq \beta_c$, then for $(C^*(S), \mathbb{R}, \sigma)$ we have

- 1 $\beta \in [0, 1)$: no KMS_β state;
- 2 $\beta \in [1, \beta_c]$: unique KMS_β if action $S_c \curvearrowright S/S_c$ essentially free, where $S_c \subset S$ subsemigroup of elements having LCM with any t in S ;
- 3 $\beta \in (\beta_c, \infty]$: KMS_β states parametrised by normalised traces on $C^*(S_c)$;

KMS_∞ strictly subset of ground states

Theorem (Afsar-Brownlowe-L-Stammeier (2016))

Let S be a right LCM monoid and $N : S \rightarrow \mathbb{N}^\times$ homomorphism such that S is *admissible*. Consider the time evolution $\sigma_t(v_s) = N_s^{it} v_s$. If $\beta_c \in \mathbb{R}$ is such that the function

$$\zeta_N(\beta) := \sum_{n \in \text{Irr}(N(S))} n^{-(\beta-1)},$$

converges for $\beta \geq \beta_c$, then for $(C^*(S), \mathbb{R}, \sigma)$ we have

- 1 $\beta \in [0, 1)$: no KMS_β state;
- 2 $\beta \in [1, \beta_c]$: unique KMS_β if action $S_c \curvearrowright S/S_c$ essentially free, where $S_c \subset S$ subsemigroup of elements having LCM with any t in S ;
- 3 $\beta \in (\beta_c, \infty]$: KMS_β states parametrised by normalised traces on $C^*(S_c)$;
- 4 Ground states: parametrised by states on $C^*(S_c)$.