# A Finite Element Method For Nematic Liquid Crystals With Variable Degree Of Orientation

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Mathematical Models and Numerical Methods
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#### **Outline**

Ericksen's Model

Discrete Energy and Finite Element Method

 $\Gamma$ -Convergence

**Gradient Flow** 

**Numerical Experiments** 

#### **OUTLINE**

#### Ericksen's Model

Discrete Energy and Finite Element Method

 $\Gamma$ -Convergence

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Numerical Experiments

### Macroscopic Model of Liquid Crystals (De Gennes)

- Liquid crystals find everyday use in modern technology.
- LCDs, optics, etc.

Fricksen's Model



- ▶ Liquid crystal *molecules* are often *idealized* as elongated rods or ellipsoidal disks.
- Further simplify by an averaging procedure to replace local arrangements of many rods by a **few** order parameters.

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- Liquid crystal molecules are often idealized as elongated rods or ellipsoidal disks.
- Further simplify by an averaging procedure to replace local arrangements of many rods by a few order parameters.

Essentially, the state of the system is:

- ▶ The local (average) orientation of molecules.
- ► The local (average) degree of orientation.

One model that captures this is the Q-tensor approach [De Gennes 74].

Ericksen's Model

#### Model: Ericksen's model

 $lackbox{ }Q$  tensor is a symmetric, traceless  $3\times 3$  tensor. Thus,  $oldsymbol{Q}$  is of the form

$$\mathbf{Q} = -(s_1\mathbf{n}_1 \otimes \mathbf{n}_1 + s_2\mathbf{n}_2 \otimes \mathbf{n}_2) + \frac{1}{3}(s_1 + s_2)\mathbf{I}$$

with eigenvalues between -1/3 and 2/3.

- ▶ Uniaxial liquid crystal: Q tensor reduces to  $Q = s(n \otimes n \frac{1}{3}I)$ .
- Further simplification: deal with the director n.
- ▶ The equilibrium state minimizes (one-constant Ericksen's model)

$$E := \underbrace{\int_{\Omega} \kappa |\nabla s|^2 + s^2 |\nabla n|^2 dx}_{:=E_1} + \underbrace{\int_{\Omega} \psi_B(s) dx}_{:=E_2}$$

where  $\kappa > 0$  and  $\psi_B$  is a double well potential

- ▶ s is the degree of orientation (-1/2 < s < 1).
- ightharpoonup s = 1: perfect alignment with n.
- ightharpoonup s = 0: no preferred direction (isotropic). This defines the set of defects:

$$\{x \in \Omega, \ s(x) = 0\}$$

ightharpoonup s = -1/2: perpendicular to n.

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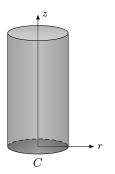
Fricksen's Model

#### Ericksen's model vs. Oseen-Frank model

▶ When  $s = s_0 > 0$ , the energy reduces to the Oseen-Frank energy:

$$E := \int_{\Omega} |\nabla \boldsymbol{n}|^2 dx.$$

Line and plane defects have infinite energy in the Oseen-Frank model:



Line defects:

$$oldsymbol{n} = rac{oldsymbol{r}}{|oldsymbol{r}|}, \quad |
abla oldsymbol{n}| = rac{2}{|oldsymbol{r}|}.$$

• Compute  $\int_C |\nabla \boldsymbol{n}|^2 dx$ :

$$\int_0^1 \frac{4}{r^2} \ r dr = \infty.$$

► Ericksen's model regularizes the defect.

▶ Goal: design a robust finite element method to capture these defects.

#### Literature Review

Ericksen's Model

```
Hardt Kinderlehrer Lin (1986), Kinderlehrer (1991), Kinderlehrer, Ou,
Walkington (1993)
Brezis (1987, 1989) Ericksen (1991)
book: Virga (1994)
Ambrosio (1990)
Lin (1989, 1991), Lin Liu (1995)
Bauman Calderer Liu Phillips (2002), Bauman Park Phillps (2012)
Ball Zarnescu (2011)
Alouges (1997)
Badia Guillén-González Gutiérrez-Santacreu (2011, 2013)
Bartels (2010)
Cohen Lin Luskin (1989)
Liu Walkington (2000)
Yang Forest Li Liu Shen Wang (2013)
Calderer, Golovaty, Lin, Liu (2002) (Time evolution for Ericksen model)
Barrett Feng Prohl (2006) (2D-FEM via regularization)
James Willman Fernández (2006) (Q tensor method)
Shin Cho Lee Yoon and Won (2008) (Q tensor method)
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#### **OUTLINE**

Ericksen's Mode

#### Discrete Energy and Finite Element Method

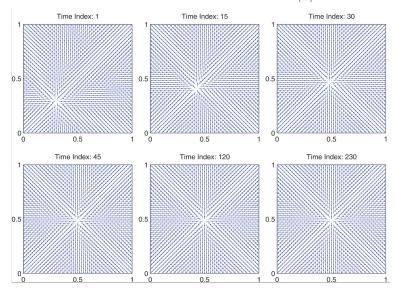
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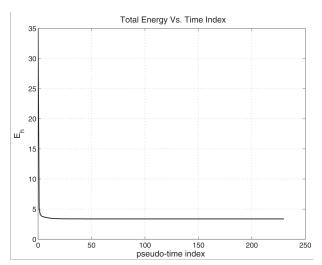
### Numerical Experiments: (point defect in 2d)

Consider the Dirichlet boundary conditions  $s=s^*, \quad \boldsymbol{n}=\frac{\boldsymbol{x}}{|\boldsymbol{x}|}.$  MOVIE



### Numerical Experiments: (point defect in 2d)

#### Energy Decrease:



#### Admissible class

• Assume  $\kappa > 0$ . We seek a **discrete energy**  $E_1^h$  to approximate

$$E_1 := \int_{\Omega} \kappa |\nabla s|^2 + s^2 |\nabla \mathbf{n}|^2 dx.$$

▶ Since  $\nabla |\mathbf{n}|^2 = 2(\nabla \mathbf{n})\mathbf{n} = \mathbf{0}$ , we have an **identity**:

$$\int_{\Omega} |\nabla \underbrace{(s\boldsymbol{n})}_{\boldsymbol{\cdot} = \boldsymbol{n}}|^2 dx = \int_{\Omega} |\nabla s \otimes \boldsymbol{n} + s \nabla \boldsymbol{n}|^2 dx = \int_{\Omega} |\nabla s|^2 + s^2 |\nabla \boldsymbol{n}|^2 dx.$$

We rewrite the energy [Ambrosio 90, Lin 91]:

$$E_1 = \int_{\Omega} (\kappa - 1) |\nabla s|^2 + |\nabla (s\boldsymbol{n})|^2 dx,$$

i.e. a simple quadratic functional, but with a negative term.

Admissible class:

$$\mathcal{A} := \{(s, \boldsymbol{n}) : s \in H^1(\Omega), \ \boldsymbol{u} = s\boldsymbol{n} \in H^1(\Omega) \text{ and } |\boldsymbol{n}| = 1 \text{ a.e. in } \Omega\}.$$

#### **Numerical Discretization**

- Let  $\mathcal{T}_h = \{T\}$  be a conforming, shape-regular triangulation of  $\Omega$ , with set of nodes (vertices) denoted by  $\mathcal{N}_h$ .
- **Exact solution**: (s, n); Discrete solution (S, N) or  $(s_h, n_h)$ .
- ▶ Piecewise linear approximation:  $S \in W_h$ ,  $N \in V_h$ :

$$\begin{split} W_h &:= \{S \in H^1(\Omega) : S|_T \text{ is affine}\}, \\ U_h &:= \{ \boldsymbol{U} \in H^1(\Omega)^d : \boldsymbol{U}|_T \text{ is affine in each component}\}, \\ V_h &:= \{ \boldsymbol{N} \in U_h : |\boldsymbol{N}(\boldsymbol{x_i})| = 1 \text{ at all nodes } x_i \in \mathcal{N}_h \}. \end{split}$$

- I.e. impose the unit length constraint at the mesh nodes.
- ▶ Denote the continuous piecewise linear "hat" basis functions by  $\{\phi_i\}$ .
- Assume the entries of the stiffness matrix  $\{k_{ij}\}$  satisfy

$$k_{ij} = -\int_{\Omega} \nabla \phi_i \cdot \nabla \phi_j dx \ge 0, \quad \text{for } i \ne j.$$

▶ If the mesh is weakly acute, then this condition is true.

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### Discretization of the Energy

- ▶ The condition |N(x)| = 1 is only at the nodes  $x = x_i$ ,
- ► This suggests to view the **discrete energy** in terms of **nodal values**.
- ▶ For a piecewise linear function *S*, we have

$$\int_{\Omega} |\nabla S|^2 dx = \frac{1}{2} \sum_{i,j} k_{ij} (S_i - S_j)^2,$$

where  $S_i = S(x_i)$  for all nodes  $x_i$ .

We approximate

$$E_1 = \int_{\Omega} \kappa |\nabla s|^2 + s^2 |\nabla n|^2 dx,$$

by a discrete energy:

$$E_1^h := \underbrace{\frac{\kappa}{2} \sum_{i,j=1}^N k_{ij} \left(S_i - S_j\right)^2}_{\text{standard}} + \underbrace{\frac{1}{2} \sum_{i,j=1}^N k_{ij} \left(\frac{S_i^2 + S_j^2}{2}\right) |\boldsymbol{N}_i - \boldsymbol{N}_j|^2}_{\text{not standard}}$$

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where  $N_i = N(x_i)$  for all nodes  $x_i$ .

- ▶ Let  $U \in U_h$  such that  $U(x_i) = S(x_i)N(x_i)$  for all nodes  $x_i$ .
- Energy inequality:

$$E_1^h[\Omega, S, \mathbf{N}] \ge (\kappa - 1) \int_{\Omega} |\nabla S|^2 dx + \int_{\Omega} |\nabla \mathbf{U}|^2 dx.$$

Moreover

$$E_1^h[\Omega, S, N] \ge (\kappa - 1) \int_{\Omega} |\nabla I_h| S ||^2 dx + \int_{\Omega} |\nabla \widetilde{\boldsymbol{U}}|^2 dx.$$

- $ightharpoonup I_h$  is the linear interpolant of |S|.
- $ightharpoonup \widetilde{U} \in U_h$ , such that  $\widetilde{U}(x_i) = |S(x_i)|N(x_i)$  for all nodes  $x_i$ .

**Remark:** eventually, we need the right-hand-side to be **convex** with respect to the **gradient**. Thus, we need |S|.

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Remark: eventually, we need the right-hand-side to be convex with respect to the **gradient**. Thus, we need |S|.

$$\text{prove this:} \quad E_1^h \geq (\kappa-1) \int_{\Omega} \left| \nabla S \right|^2 \! dx + \int_{\Omega} \left| \nabla \boldsymbol{U} \right|^2 \! dx.$$

$$\begin{split} \int_{\Omega} |\nabla \boldsymbol{U}|^2 dx &= \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} |\underbrace{\boldsymbol{S}_i \boldsymbol{N}_i}_{\boldsymbol{U}_i} - \underbrace{\boldsymbol{S}_j \boldsymbol{N}_j}_{\boldsymbol{U}_j}|^2 \\ &= \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} \left| \left( \frac{S_i + S_j}{2} \right) (\boldsymbol{N}_i - \boldsymbol{N}_j) + (S_i - S_j) \left( \frac{\boldsymbol{N}_i + \boldsymbol{N}_j}{2} \right) \right|^2. \end{split}$$

$$(oldsymbol{N}_i-oldsymbol{N}_j)\cdot(oldsymbol{N}_i+oldsymbol{N}_j)=|oldsymbol{N}_i|^2-|oldsymbol{N}_j|^2=0.$$
 Thus

$$\begin{split} \int_{\Omega} |\nabla U|^2 dx &= \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} \left( \frac{S_i + S_j}{2} \right)^2 |N_i - N_j|^2 + \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} (S_i - S_j)^2 \left| \frac{N_i + N_j}{2} \right| \\ &= \frac{1}{2} \sum_{i=1}^{N} k_{ij} \left( \frac{S_i + S_j}{2} \right)^2 |N_i - N_j|^2 + \frac{1}{2} \sum_{i=1}^{N} k_{ij} (S_i - S_j)^2 - \frac{1}{2} \mathcal{E}, \end{split}$$

FEM For Liquid Crystals

$$\text{prove this:} \quad E_1^h \geq (\kappa-1) \int_{\Omega} \left| \nabla S \right|^2 \! dx + \int_{\Omega} \left| \nabla \boldsymbol{U} \right|^2 \! dx.$$

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The unit length constraint on  ${m N}$  at the nodes implies:

$$({m N}_i - {m N}_j) \cdot ({m N}_i + {m N}_j) = |{m N}_i|^2 - |{m N}_j|^2 = 0.$$
 Thus,

$$\int_{\Omega} |\nabla U|^2 dx = \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} \left( \frac{S_i + S_j}{2} \right)^2 |\mathbf{N}_i - \mathbf{N}_j|^2 + \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} (S_i - S_j)^2 \left| \frac{\mathbf{N}_i + \mathbf{N}_j}{2} \right|^2$$

$$= \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} \left( \frac{S_i + S_j}{2} \right)^2 |\mathbf{N}_i - \mathbf{N}_j|^2 + \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} (S_i - S_j)^2 - \frac{1}{2} \mathcal{E},$$

where we used  $\left|\frac{{m N}_i + {m N}_j}{2}\right|^2 = 1 - \left|\frac{{m N}_i - {m N}_j}{2}\right|^2.$ 

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where

$$\mathcal{E} := \sum_{i,j=1}^N k_{ij} \left(S_i - S_j\right)^2 \left| rac{oldsymbol{N}_i - oldsymbol{N}_j}{2} 
ight|^2$$
 (positive term).

Next, plug in:

$$\left(\frac{S_i + S_j}{2}\right)^2 = \left(\frac{S_i^2 + S_j^2}{2}\right) - \left(\frac{S_i - S_j}{2}\right)^2,$$

to get

$$\int_{\Omega} |\nabla \boldsymbol{U}|^2 dx = \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} (S_i - S_j)^2 + \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} \left( \frac{S_i^2 + S_j^2}{2} \right) |\boldsymbol{N}_i - \boldsymbol{N}_j|^2 - \mathcal{E},$$

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prove this: 
$$E_1^h \ge (\kappa - 1) \int_{\Omega} |\nabla S|^2 dx + \int_{\Omega} |\nabla U|^2 dx$$
.

Adding  $(\kappa - 1) \int_{\Omega} |\nabla S|^2 dx$  to both sides gives

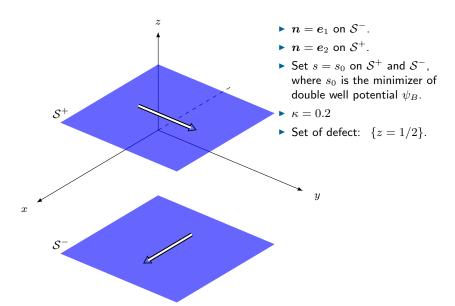
$$(\kappa - 1) \int_{\Omega} |\nabla S|^{2} dx + \int_{\Omega} |\nabla U|^{2} dx$$

$$= \frac{\kappa}{2} \sum_{i,j=1}^{N} k_{ij} (S_{i} - S_{j})^{2} + \frac{1}{2} \sum_{i,j=1}^{N} k_{ij} \left( \frac{S_{i}^{2} + S_{j}^{2}}{2} \right) |\mathbf{N}_{i} - \mathbf{N}_{j}|^{2} - \mathcal{E}$$

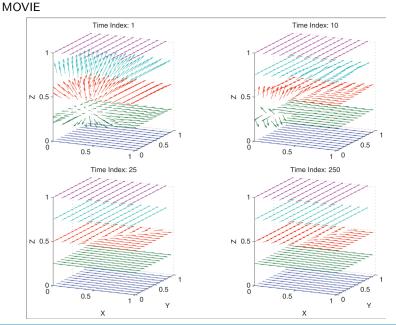
$$< E_{1}^{h}.$$

Note that  $\mathcal{E} > 0$  provided the meshes are weakly acute.

### Numerical Experiments: (plane defect in 3d) [Ambrosio Virga 1991]

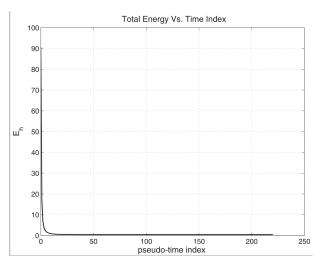


## Numerical Experiments: (plane defect in 3d)



### Numerical Experiments: (plane defect in 3d)

### Energy Decrease:



Fricksen's Mode

Discrete Energy and Finite Element Method

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#### Main Theorem

### $\Gamma$ -Convergence of the Discrete Energy $E_1^h$ :

- ▶ Let  $\{\mathcal{T}_h\}$  be a sequence of weakly acute meshes.
- ▶ lim-sup: there exists a sequence  $\{(s_h, n_h)\}$ , such that  $(s_h, n_h)$  converges to (s, n) in  $L^2$ , and

$$E_1[s, \boldsymbol{n}] \ge \limsup_{h \to 0} E_1^h[s_h, \boldsymbol{n}_h].$$

▶ **lim-inf:** for *every* sequence  $\{(s_h, n_h)\}$ , such that  $(s_h, n_h)$  converges to (s, n) in  $L^2$ , we have

$$E_1[s, \boldsymbol{n}] \leq \liminf_{h \to 0} E_1^h[s_h, \boldsymbol{n}_h].$$

#### Convergence of the Finite Element Method

- ▶  $E_1^h[s_h, n_h]$  is coercive: any sequence  $\{(s_h, n_h)\}$  with finite discrete energy is pre-compact in  $L^2$ .
- Let  $\{(s_h, n_h)\}$  be a minimizing sequence of  $E_1^h[s_h, n_h]$ .
- ▶ Then  $(s_h, n_h)$  converges to (s, n) in  $L^2$ , where (s, n) is a minimizer of  $E_{r_0}[s, n]$

### $\Gamma$ -Convergence of the Discrete Energy $E_1^h$ :

- ▶ Let  $\{\mathcal{T}_h\}$  be a sequence of weakly acute meshes.
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- ▶ Let  $\{(s_h, n_h)\}$  be a minimizing sequence of  $E_1^h[s_h, n_h]$ .
- ▶ Then  $(s_h, n_h)$  converges to (s, n) in  $L^2$ , where (s, n) is a minimizer of  $E_1[s, n]$ .

Let (s, n) be in  $\mathcal{A}$  and set u = sn. Then

$$E_1^h[s_h, \boldsymbol{n}_h] \to E_1[s, \boldsymbol{n}], \text{ as } h \to 0,$$

where  $(s_h, \boldsymbol{n}_h)$  are the Lagrange interpolants of  $(s, \boldsymbol{n})$ .

▶ Recall the identity

$$E_1^h[s_h, \boldsymbol{n}_h] = (\kappa - 1) \int_{\Omega} |\nabla s_h|^2 dx + \int_{\Omega} |\nabla \boldsymbol{u}_h|^2 dx + \mathcal{E}$$

▶ Since  $s_h \rightarrow s$ ,  $u_h \rightarrow u$ , we only need to show

$$\mathcal{E} = \sum_{i,j=1}^N k_{ij} \left(s_i - s_j
ight)^2 \left|rac{m{n}_i - m{n}_j}{2}
ight|^2 
ightarrow 0, \;\; \mathsf{as} \;\; h 
ightarrow 0.$$

where  $s_i = s_h(x_i)$ ,  $\boldsymbol{n}_i = \boldsymbol{n}_h(x_i)$ 

- $\triangleright \mathcal{E} \approx 2h^2 \int_{\Omega} |\nabla s_h|^2 dx$  if  $\boldsymbol{n}$  is smooth.
- ▶ But *n* may have a **discontinuity**

#### **Limit-Sup Inequality**

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#### Proof:

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### Consistency of the Discrete Energy

#### Proof:

We divide  $\Omega$  into two regions:

$$\mathcal{S}_{\epsilon} = \{x \in \Omega, \ |s(x)| < \epsilon\} \quad \text{ and } \quad \mathcal{K}_{\epsilon} = \overline{\Omega} \setminus \mathcal{S}_{\epsilon}.$$

Step 1, Estimate on  $\mathcal{K}_{\epsilon}$ :

$$\sum_{x_i, x_j \in \mathcal{K}_{\epsilon}} k_{ij} (s_i - s_j)^2 \left| \frac{\boldsymbol{n}_i - \boldsymbol{n}_j}{2} \right|^2 \leq C \left( \max_{\substack{x_i, x_j \in \mathcal{K}_{\epsilon}, \\ |x_i - x_j| \leq h}} |\boldsymbol{n}_i - \boldsymbol{n}_j|^2 \right) \int_{\Omega} |\nabla s_h|^2 dx.$$

Γ-Convergence

$$\sum_{\text{either } x_i \text{ or } x_j \in \mathcal{S}_{\epsilon}} k_{ij} (s_i - s_j)^2 \left| \frac{\boldsymbol{n}_i - \boldsymbol{n}_j}{2} \right|^2 \le 2 \int_{\cup \omega_i} |\nabla s_h|^2 dx \le 2 \int_{\mathcal{S}_{2\epsilon}} |\nabla s_h|^2 dx,$$

$$\mathcal{E} \leq C \underbrace{\left(\max_{\substack{x_i, x_j \in \mathcal{K}_{\epsilon}, \\ |x_i - x_j| \leq h}} |n_i - n_j|^2\right)}_{\mathcal{O}(h^2)} \int_{\mathcal{K}_{\epsilon}} |\nabla s_h|^2 dx + 2 \underbrace{\int_{\mathcal{S}_{2\epsilon}} |\nabla s_h|^2 dx}_{\to 0 \text{ as } \epsilon \to 0}.$$

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Γ-Convergence

▶ Step 2, Estimate on  $S_{\epsilon}$ :

$$\sum_{\text{either } x_i \text{ or } x_j \in \mathcal{S}_{\epsilon}} k_{ij} (s_i - s_j)^2 \left| \frac{\boldsymbol{n}_i - \boldsymbol{n}_j}{2} \right|^2 \leq 2 \int_{\cup \omega_i} |\nabla s_h|^2 dx \leq 2 \int_{\mathcal{S}_{2\epsilon}} |\nabla s_h|^2 dx,$$

where  $\cup \omega_i$  is taken over all nodes  $x_i$  in  $S_{\epsilon}$ .

▶ Combining both estimates, we have

$$\mathcal{E} \leq C \underbrace{\left( \max_{\substack{x_i, x_j \in \mathcal{K}_{\epsilon,i} \\ |x_i - x_j| \leq h}} |\boldsymbol{n}_i - \boldsymbol{n}_j|^2 \right)}_{O(h^2)} \int_{\mathcal{K}_{\epsilon}} |\nabla s_h|^2 dx + 2 \underbrace{\int_{\mathcal{S}_{2\epsilon}} |\nabla s_h|^2 dx}_{\to 0 \text{ as } \epsilon \to 0}.$$

### **Limit-Inf Inequality**

The **main issue** is dealing with the case  $0 < \kappa < 1$ .

Recall a previous result:

$$E_1^h[\Omega, s_h, \boldsymbol{n}_h] \ge (\kappa - 1) \int_{\Omega} |\nabla I_h| s_h ||^2 dx + \int_{\Omega} |\nabla \tilde{\boldsymbol{u}}_h|^2 dx,$$

where  $\tilde{\boldsymbol{u}}_h$  in  $U_h$  and  $\tilde{\boldsymbol{u}}_h(x_i) = |s_h(x_i)| \boldsymbol{n}_h(x_i)$  for all nodes  $x_i$ .

**Coercivity:** for all  $\kappa > 0$ , we have

$$E_1^h[\Omega, s_h, \boldsymbol{n}_h] \ge \min\{\kappa, 1\} \int_{\Omega} |\nabla \tilde{\boldsymbol{u}}_h|^2 dx \ge \min\{\kappa, 1\} \int_{\Omega} |\nabla I_h| s_h ||^2 dx.$$

$$\int_{\Omega} L_h(\boldsymbol{w}_h, 
abla \boldsymbol{w}_h) dx$$
, where

$$L_h(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h) := (\kappa - 1)|\nabla I_h|\boldsymbol{w}_h||^2 + |\nabla \boldsymbol{w}_h|^2,$$

$$\liminf_{h\to 0} \int_{\Omega} L_h(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h) dx \ge \int_{\Omega} (\kappa - 1) |\nabla |\boldsymbol{w}||^2 + |\nabla \boldsymbol{w}|^2 dx$$

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▶ Weak Lower Semi-continuity: Let  $w_h \in U_h$ . The energy

$$\int_{\Omega} L_h(oldsymbol{w}_h, 
abla oldsymbol{w}_h) dx,$$
 where

$$L_h(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h) := (\kappa - 1) |\nabla I_h| \boldsymbol{w}_h|^2 + |\nabla \boldsymbol{w}_h|^2,$$

is weakly lower semi-continuous, i.e. we have

$$\liminf_{h\to 0} \int_{\Omega} L_h(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h) dx \ge \int_{\Omega} (\kappa - 1) |\nabla |\boldsymbol{w}||^2 + |\nabla \boldsymbol{w}|^2 dx.$$

for any weakly convergent sequence  $w_h \rightharpoonup w$  in the  $H^1$  norm.

FEM For Liquid Crystals

# Numerical Experiment: (plane defect in 3d)

Visualization of defect formation:

**MOVIE** 

### Proof:

#### Main Goal:

Show that

$$L_h(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h) := (\kappa - 1)|\nabla I_h|\boldsymbol{w}_h||^2 + |\nabla \boldsymbol{w}_h|^2$$

Γ-Convergence

is convex with respect to  $\nabla w_h$ , even for  $0 < \kappa < 1$ .

### Rewrite energy density:

▶ Suppose dimension is d = 2. Let T be a triangle in  $\mathcal{T}_h$  with vertices  $x_0, x_1, x_2$ . Define

$$e_i := x_i - x_0$$
, for  $i = 1, 2$ ,  $w_i := w_h(x_i)$  for  $i = 0, 1, 2$ .

A simple calculation gives:

$$abla oldsymbol{w}_h = (oldsymbol{w}_1 - oldsymbol{w}_0) \otimes oldsymbol{e}_1^* + (oldsymbol{w}_2 - oldsymbol{w}_0) \otimes oldsymbol{e}_2^*, 
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where  $e_i \cdot e_i^* = \delta_{ii}$  (dual basis).

$$|w_i| - |w_0| = \frac{w_i + w_0}{|w_i| + |w_0|} \cdot (w_i - w_0)$$

Γ-Convergence

# Weak Lower Semi-continuity

### Proof:

#### Main Goal:

Show that

$$L_h(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h) := (\kappa - 1)|\nabla I_h|\boldsymbol{w}_h||^2 + |\nabla \boldsymbol{w}_h|^2$$

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Note:

$$|w_i| - |w_0| = \frac{w_i + w_0}{|w_i| + |w_0|} \cdot (w_i - w_0).$$

### **Proof:**

► Therefore,

$$\nabla I_h |\boldsymbol{w}_h| = G_h(\boldsymbol{w}_h) : \nabla \boldsymbol{w}_h,$$

where  $G_h$  is a 3-tensor:

$$G_h(oldsymbol{w}_h) := rac{oldsymbol{w}_1 + oldsymbol{w}_0}{|oldsymbol{w}_1| + |oldsymbol{w}_0|} \otimes oldsymbol{e}_1 \otimes oldsymbol{e}_1 \otimes oldsymbol{e}_1^* + rac{oldsymbol{w}_2 + oldsymbol{w}_0}{|oldsymbol{w}_2| + |oldsymbol{w}_0|} \otimes oldsymbol{e}_2 \otimes oldsymbol{e}_2^*.$$

- ▶ We define  $(g_1 \otimes g_2 \otimes g_3) : (m_1 \otimes m_2) = (g_1 \cdot m_1)(g_2 \cdot m_2)g_3$ .
- Hence,

$$L_h(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h) = |\nabla \boldsymbol{w}_h|^2 + (\kappa - 1)|G_h : \nabla \boldsymbol{w}_h|^2.$$

#### Note:

- $e_1 \otimes e_1^* + e_2 \otimes e_2^* = I$ , i.e. the identity matrix.
- ▶ If  $w_h \to w$  a.e., then, for a.e. x such that  $w(x) \neq 0$ ,

$$\frac{\boldsymbol{w}_i + \boldsymbol{w}_0}{|\boldsymbol{w}_i| + |\boldsymbol{w}_0|} \to \frac{\boldsymbol{w}}{|\boldsymbol{w}|}, \quad \text{for } i = 1, 2 \quad \Rightarrow \quad G_h \to G(\boldsymbol{w}) := \frac{\boldsymbol{w}}{|\boldsymbol{w}|} \otimes I.$$

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#### Proof:

► Claim: the energy density

$$L(\mathbf{w}, M) := |M|^2 + (\kappa - 1)|G(\mathbf{w}) : M|^2$$

Γ-Convergence

is convex with respect to any matrix M.

#### Proof:

► Claim: the energy density

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- ▶ Indeed, L(w, M) is quadratic in M, so we only need to show that L(w, M) > 0 for any M.
- ▶ This is equivalent to showing  $|G:M| \le |M|$ , which follows by simple inequalities.
- A similar argument shows that  $L_h(\boldsymbol{w}_h, M) > 0$  for any matrix M, and so also convex.

For the remainder of the proof, letting  $w_h$  be a weakly convergent sequence in  $H^1(\Omega)$ , and applying standard limiting arguments, we get the assertion.

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# **Limit-Inf Inequality**

#### Proof:

- Let  $\{(s_h, n_h)\}$  be any sequence converging strongly to (s, n) in  $L^2$ .
- We know that the  $H^1$  norms of  $I_h|s_h|$  and  $\tilde{\boldsymbol{u}}_h$  are bounded.
- Extract a subsequence  $\{(I_h|s_h|, \tilde{\boldsymbol{u}}_h)\}$  converging weakly in  $H^1$  and strongly in  $L^2$  to  $(|s|, \tilde{\boldsymbol{u}})$ .
- Moreover, one can show  $\tilde{\boldsymbol{u}}(x) = |s(x)|\boldsymbol{n}(x)$ , with  $|\boldsymbol{n}(x)| = 1$ , for a.e. x.

$$\liminf_{h \to 0} E_1^h[s_h, \boldsymbol{n}_h] \ge \int_{\Omega} (\kappa - 1) |\nabla |\tilde{\boldsymbol{u}}||^2 + |\nabla \tilde{\boldsymbol{u}}|^2 dx$$

$$= \int_{\Omega} \kappa |\nabla |s||^2 + |s|^2 |\nabla \boldsymbol{n}|^2 dx$$

$$= \int_{\Omega} \kappa |\nabla s|^2 + s^2 |\nabla \boldsymbol{n}|^2 dx$$

$$= E_1[s, \boldsymbol{n}].$$

#### Proof:

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- ▶ Moreover, one can show  $\tilde{\boldsymbol{u}}(x) = |s(x)|\boldsymbol{n}(x)$ , with  $|\boldsymbol{n}(x)| = 1$ , for a.e. x.

By previous results and weak lower semi-continuity:

$$\begin{aligned} \liminf_{h \to 0} E_1^h[s_h, \boldsymbol{n}_h] &\geq \int_{\Omega} (\kappa - 1) |\nabla |\tilde{\boldsymbol{u}}||^2 + |\nabla \tilde{\boldsymbol{u}}|^2 dx \\ &= \int_{\Omega} \kappa |\nabla |s||^2 + |s|^2 |\nabla \boldsymbol{n}|^2 dx \\ &= \int_{\Omega} \kappa |\nabla s|^2 + s^2 |\nabla \boldsymbol{n}|^2 dx \\ &= E_1[s, \boldsymbol{n}]. \end{aligned}$$

### **OUTLINE**

Ericksen's Mode

Discrete Energy and Finite Element Method

 $\Gamma$ -Convergence

**Gradient Flow** 

Numerical Experiments

- We design a gradient flow to seek a minimizer (S, N) of the discrete energy  $E^h[S, N]$ .

$$\mathbb{T}_h^k = \{ \boldsymbol{T} \in H^1(\Omega), \boldsymbol{T}|_T \text{ is affine, and } \boldsymbol{T}_i \cdot \boldsymbol{N}_i^k = 0 \text{ for all nodes } x_i \}.$$

▶ Step (a): find  $T^k$  in  $\mathbb{T}^k_h$  such that for any V in  $\mathbb{T}^k_h$  we have

$$\delta_{\boldsymbol{N}} E_1^h[S^k, \boldsymbol{N}^k + \boldsymbol{T}^k; \boldsymbol{V}] = 0.$$

► Step (b): normalize:

$$oldsymbol{N}_i^{k+1} := rac{oldsymbol{N}_i^k + oldsymbol{T}_i^k}{|oldsymbol{N}_i^k + oldsymbol{T}_i^k|}, \quad ext{at all nodes}$$

▶ Step (c): find  $S^{k+1}$  in  $W_h$  such that for any  $Z_h \in W_h$  we have

$$\int_{S} \frac{S_h^{k+1} - S_h^k}{\delta t} Z_h = -\delta_S E_1^h [S_h^{k+1}, N_h^{k+1}; Z_h] - \delta_S E_2^h [S_h^{k+1}; Z_h]$$

- We design a gradient flow to seek a minimizer  $(S, \mathbf{N})$  of the discrete energy  $E^h[S, \mathbf{N}]$ .
- ▶ Given the k-th iteration  $(S^k, N^k)$ , to respect the unit length constraint for  $N^{k+1}$  at all nodes, we consider a first order variation with respect to  $N^k$  in the **discrete tangent space** [Alouges 97, Bartels 10]:

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$$oldsymbol{N}_i^{k+1} := rac{oldsymbol{N}_i^k + oldsymbol{T}_i^k}{|oldsymbol{N}_i^k + oldsymbol{T}_i^k|}, \quad ext{at all nodes}.$$

▶ Step (c): find  $S^{k+1}$  in  $W_h$  such that for any  $Z_h \in W_h$  we have

$$\int_{\Omega} \frac{S_h^{k+1} - S_h^k}{\delta t} Z_h = -\delta_S E_1^h [S_h^{k+1}, \boldsymbol{N}_h^{k+1}; Z_h] - \delta_S E_2^h [S_h^{k+1}; Z_h].$$

# **Monotone Energy Decrease**

Energy decrease of the gradient flow: Given a pair  $(N_h^k, S_h^k)$ , let  $(N_h^{k+1}, S_h^{k+1})$  be the discrete gradient flow obtained by the algorithm above. Then

$$E^{h}[S_{h}^{k+1}, N_{h}^{k+1}] \le E^{h}[S_{h}^{k}, N_{h}^{k}].$$

Equality holds if and only if the flow  $(N_h^k, S_h^k)$  attains an equilibrium state, that is,

$$(N_h^{k+1}, S_h^{k+1}) = (N_h^k, S_h^k).$$

 Proof is essentially based on linear arguments, except for step (b) which uses an argument from [Alouges 97, Bartels 10].

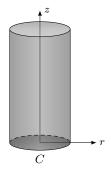
Fricksen's Mode

Discrete Energy and Finite Element Method

 $\Gamma$ -Convergence

Gradient Flow

**Numerical Experiments** 



- We neglect the double well potential  $\psi_B$ .
- Consider the minimizers of

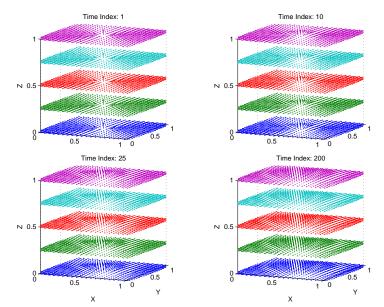
$$E = \int_C \kappa |\nabla s|^2 + s^2 |\nabla \mathbf{n}|^2 dx.$$

Boundary condition:

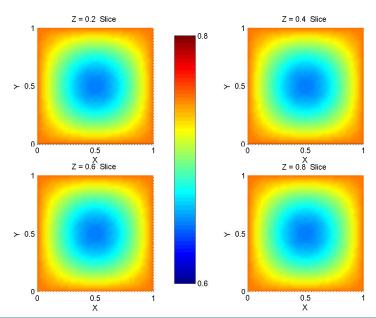
$$s|_{\mathcal{S}_0}=s_0>0 \quad \text{ and } \boldsymbol{n}|_{\mathcal{S}_0}=\frac{\boldsymbol{r}}{|\boldsymbol{r}|}.$$

► Theorem [Characterization of singular set, Mizel Roccato Virga1991]: If (s, n) is the minimizer of energy E, then either the singular set S is empty or  $S = \{ |r| = 0 \}.$ 

# Numerical Experiments: $\kappa = 10$ Fluting effect



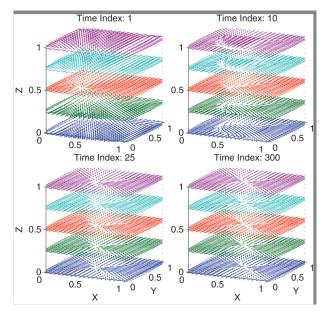
# Numerical Experiments: $\kappa = 10$ Fluting effect



Numerical Experiments

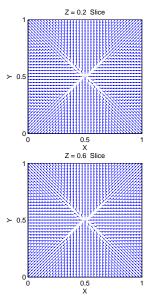
# Numerical Experiments: $\kappa = 0.1$ "Propeller" Defect

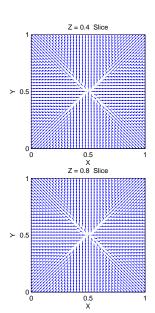
### MOVIE



# Numerical Experiments: $\kappa = 0.1$ "Propeller" Defect



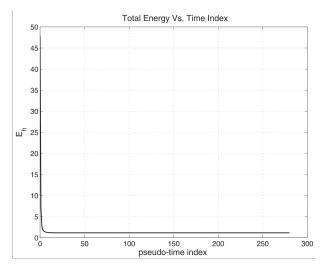




Numerical Experiments

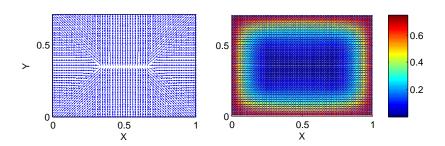
# Numerical Experiments: $\kappa=0.1$ "Propeller" Defect

### Energy Decrease:



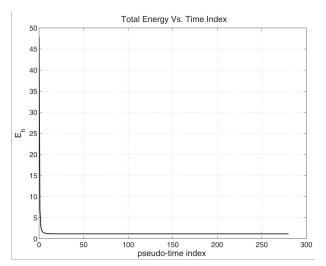
MOVIE: Director Field

MOVIE: Defect Evolution



# Numerical Experiments: $\kappa=0.1$ Defect In A Rectangular Box

### Energy Decrease:



Numerical Experiments

▶ We have a finite element method (FEM) for the one constant Ericksen's model:

$$E := \int_{\Omega} \kappa |\nabla s|^2 + s^2 |\nabla \boldsymbol{n}|^2 dx + \int_{\Omega} \psi_B(s) dx.$$

- $\blacktriangleright$  We have  $\Gamma$ -convergence of the FEM.
- ▶ We have monotone energy decrease of the gradient flow.
- Our FEM is capable of capturing high dimensional defects.

### Future Work:

- ▶ Include magnetic (electric) field to manipulate the liquid crystal.
- Investigate the Q-tensor method.
- Study flows of liquid crystals (couple Ericksen's model with Stokes flow).