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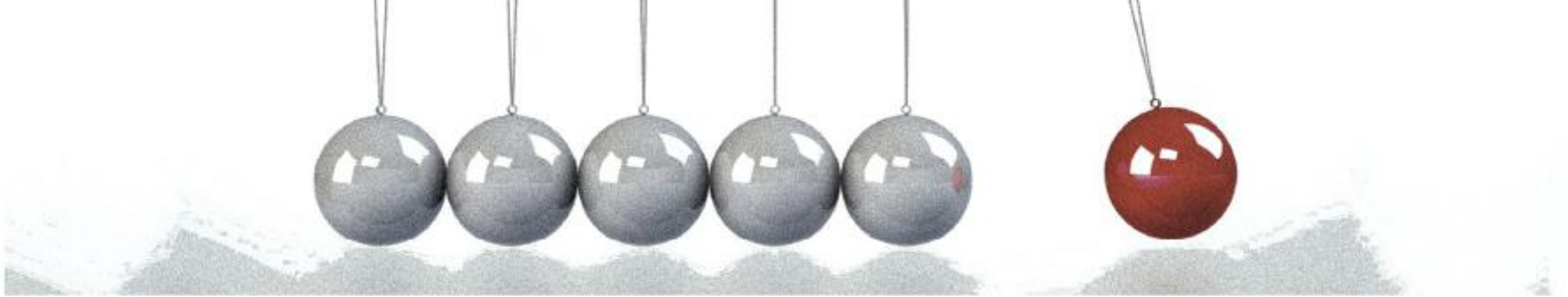
Thermoelectric material ZnSb

--- its defects and impurity band conduction

For 17th Conference of Gettering and Defect Engineering in Semiconductor Technology (GADEST 2017) Oct 01-06, 2017
Lopota Resort, Georgia



Xin Song, Postdoc
Department of Physics
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Thermoelectric material ZnSb

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--- its defects and impurity band conduction



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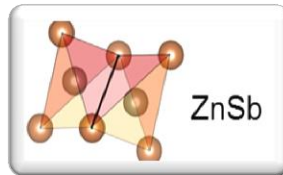


Power factor

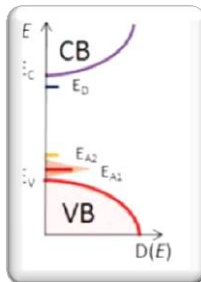
$$zT = \frac{\alpha^2 \sigma}{\kappa_e + \kappa_{ph}} T$$

Thermal conductivity from electrons and phonons

- **What matters to thermoelectrics**



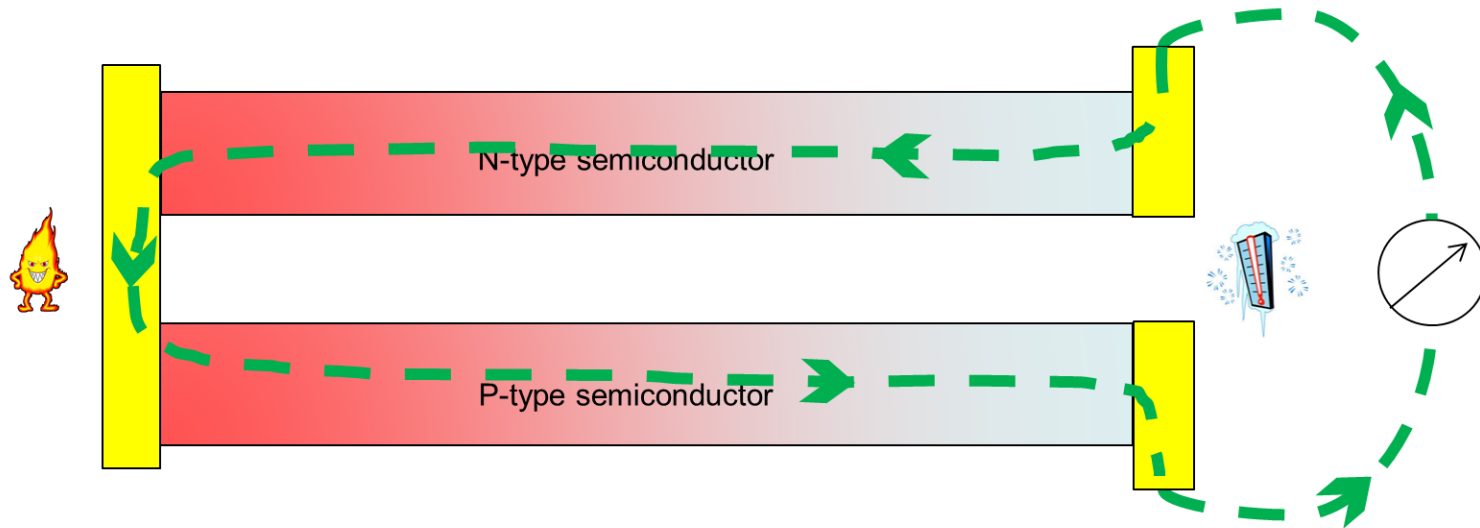
- **ZnSb and its defects**



- **Impurity band conduction and its impact on Thermoelectric performance**



Phenomenological transport of thermoelectrics

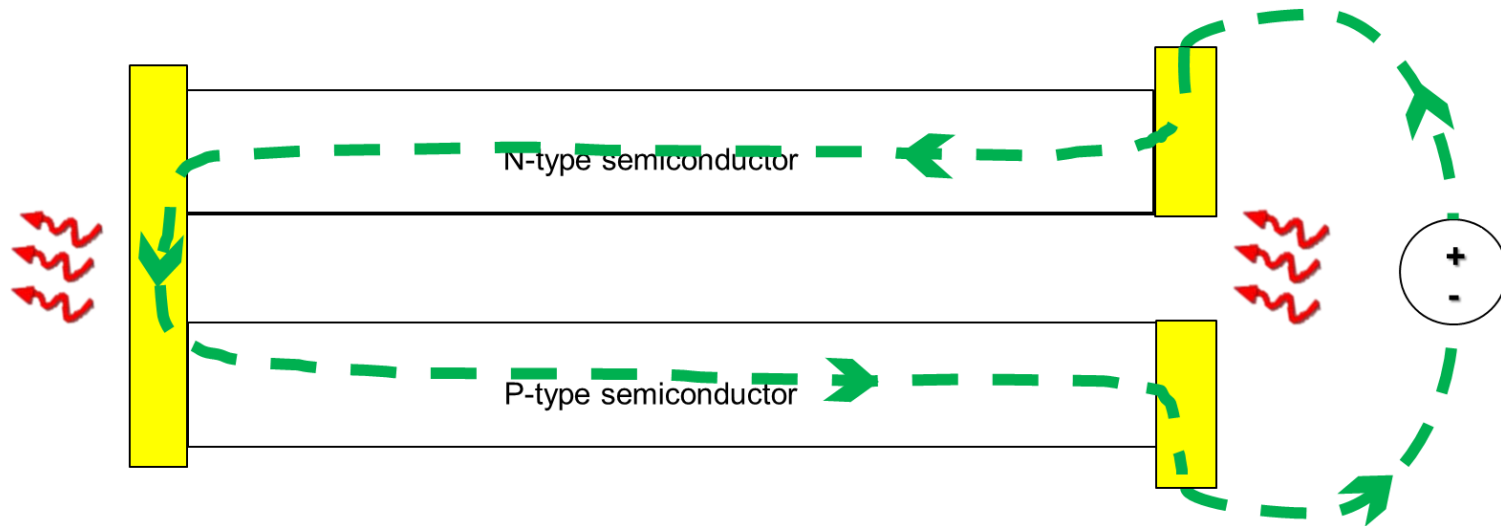


Seebeck effect

Seebeck coefficient $\alpha = \frac{\Delta V}{\Delta T}$



Phenomenological transport of thermoelectrics



Seebeck effect

Seebeck coefficient $\alpha = \frac{\Delta V}{\Delta T}$

Peltier effect

Peltier coefficient $\Pi = \frac{J_Q}{J}$



Performance of thermoelectric module and material

The coefficient of performance (COP) for a thermoelectric module

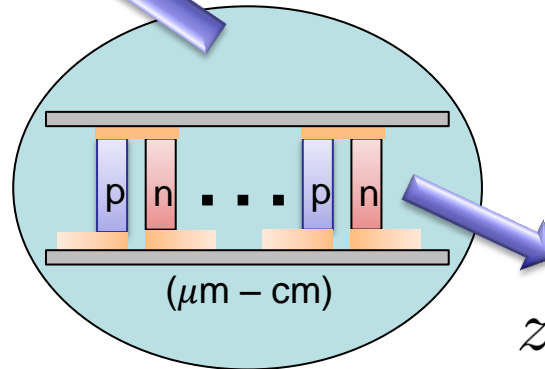
$$\eta_{COP} = \left(\frac{T_H - T_C}{T_H} \right) \left(\frac{\sqrt{1 + ZT} - 1}{\sqrt{1 + ZT} + \frac{T_C}{T_H}} \right)$$

Carnot limit

T_C -- Temperature at the cold reservoir

T_H -- Temperature at the hot reservoir

ZT -- Dimensionless thermoelectric figure of merit of module, including $(zT)_n$ and $(zT)_p$

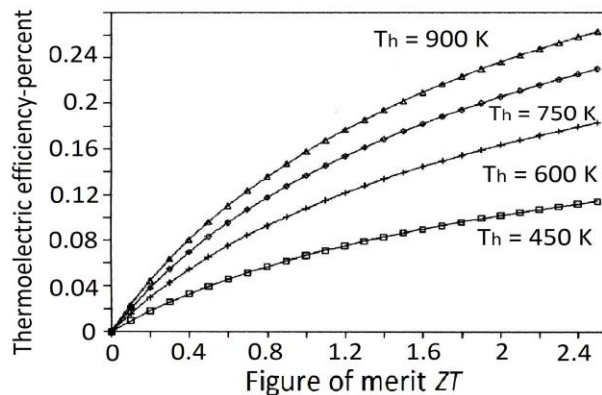


Power factor

$$zT = \frac{\alpha^2 \sigma}{\kappa_e + \kappa_{ph}} T$$

Thermal conductivity from electrons and phonons

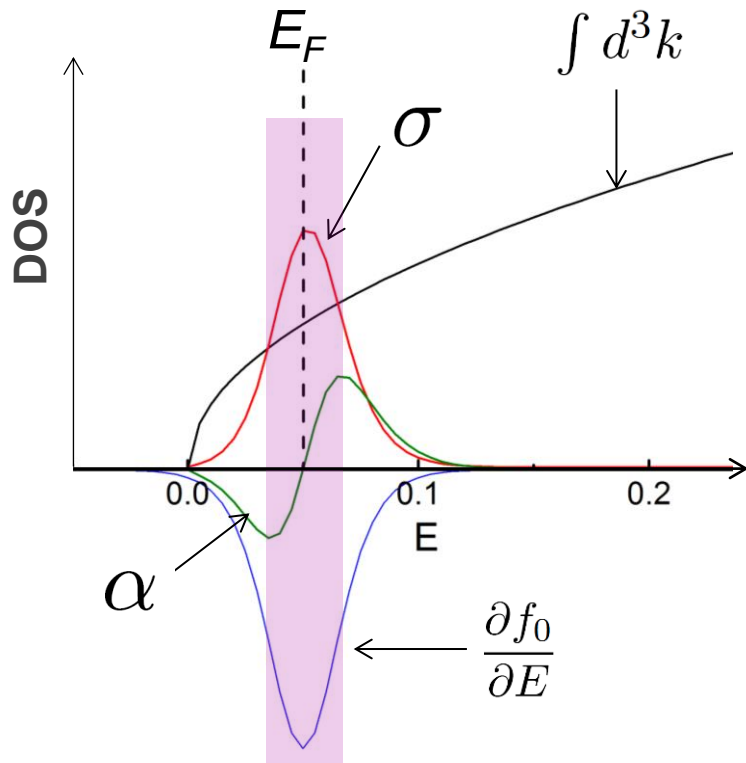
- High Seebeck coefficient α
- High electrical conductivity σ
- Low thermal conductivity κ





What matters to thermoelectrics

- Optimizing power factor



Seebeck coefficient

$$\alpha = \frac{1}{qT} \frac{\int \tau v^2 (E - E_F) \left(-\frac{\partial f_0}{\partial E}\right) d^3k}{\int \tau v^2 \left(-\frac{\partial f_0}{\partial E}\right) d^3k}$$

Electrical conductivity

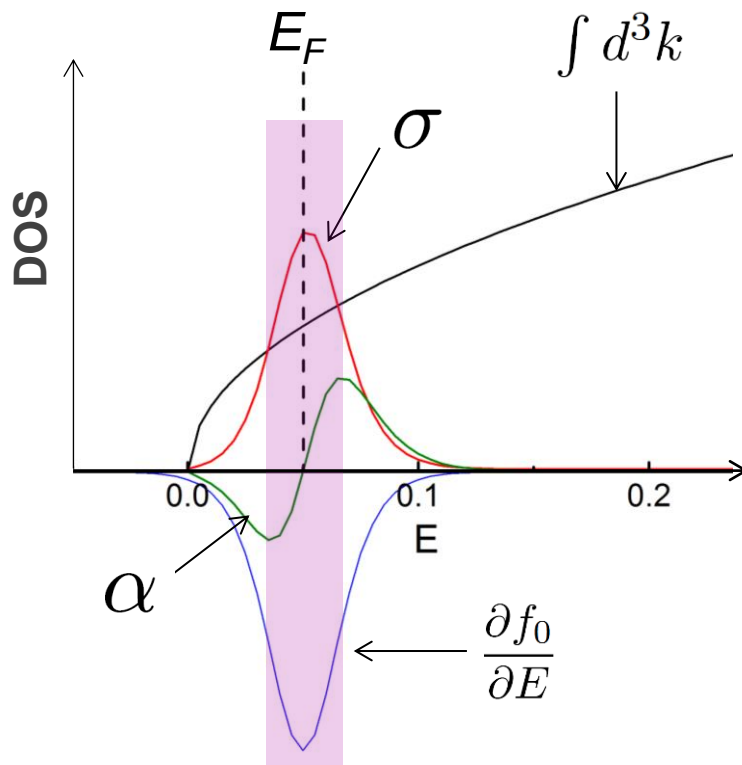
$$\sigma = \frac{q^2}{4\pi^3} \int \tau v^2 \left(-\frac{\partial f_0}{\partial E}\right) d^3k$$

$$zT = \frac{\alpha^2 \sigma}{\kappa_e + \kappa_{ph}} T$$



What matters to thermoelectrics

- Optimizing power factor



Seebeck coefficient

$$\alpha = \frac{8\pi^2 k_B^2}{3qh^2} m^* T \left(\frac{\pi}{3n} \right)^{\frac{2}{3}}$$

Electrical conductivity

$$\sigma = qn \left(\frac{q\tau}{m^*} \right)$$

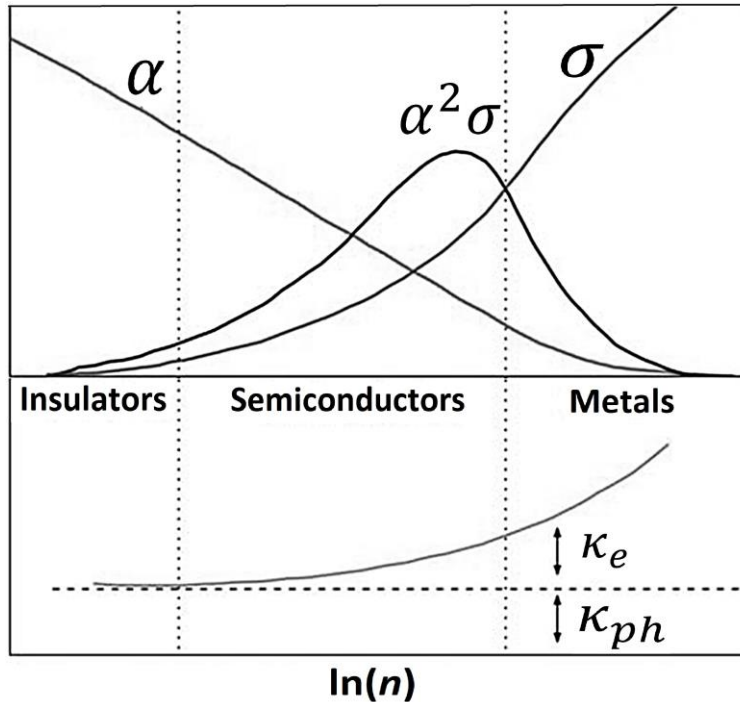
$$zT = \frac{\alpha^2 \sigma}{\kappa_e + \kappa_{ph}} T$$

- Optimize the interaction between transport coefficients
- Doping, defects scattering, energy filtering



What matters to thermoelectrics

- Optimizing power factor



Seebeck coefficient

$$\alpha = \frac{8\pi^2 k_B^2}{3qh^2} m^* T \left(\frac{\pi}{3n} \right)^{\frac{2}{3}}$$

Electrical conductivity

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$$zT = \frac{\alpha^2 \sigma}{\kappa_e + \kappa_{ph}} T$$

Wiedemann-Franz law

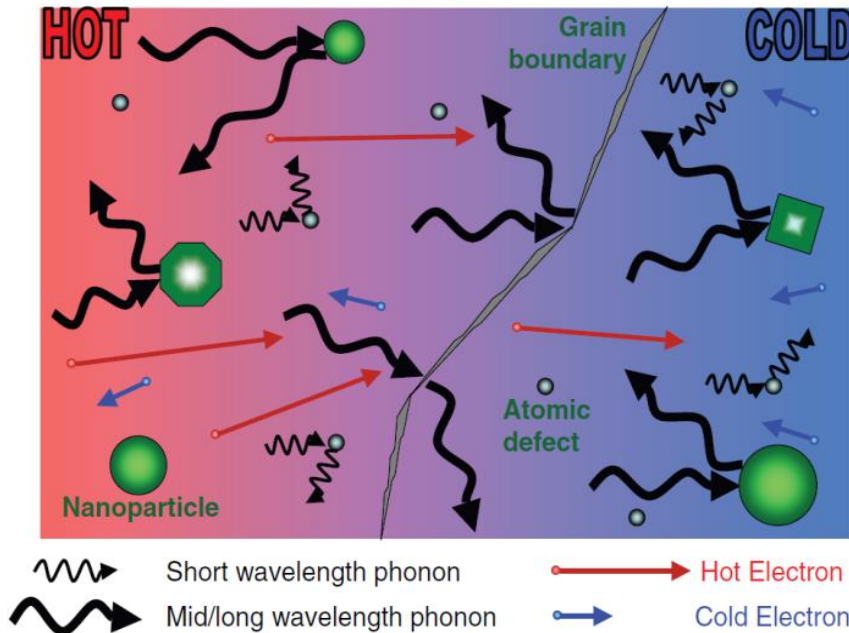
$$\kappa_e = nq\mu\mathcal{L}T$$

- Optimize the interaction between transport coefficients
- Doping, defects scattering, energy filtering



What matters to thermoelectrics

- Reducing thermal conductivity



C. J. Vineis , et al. *Adv. Mater.* 2010, 22, 3970–3980

- Phonon engineering.
- Nanostructuring, impurities.

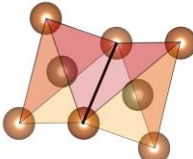
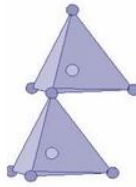
$$zT = \frac{\alpha^2 \sigma}{\kappa_e + \kappa_{ph}} T$$

Wiedemann-Franz law

$$\kappa_e = nq\mu\mathcal{L}T$$



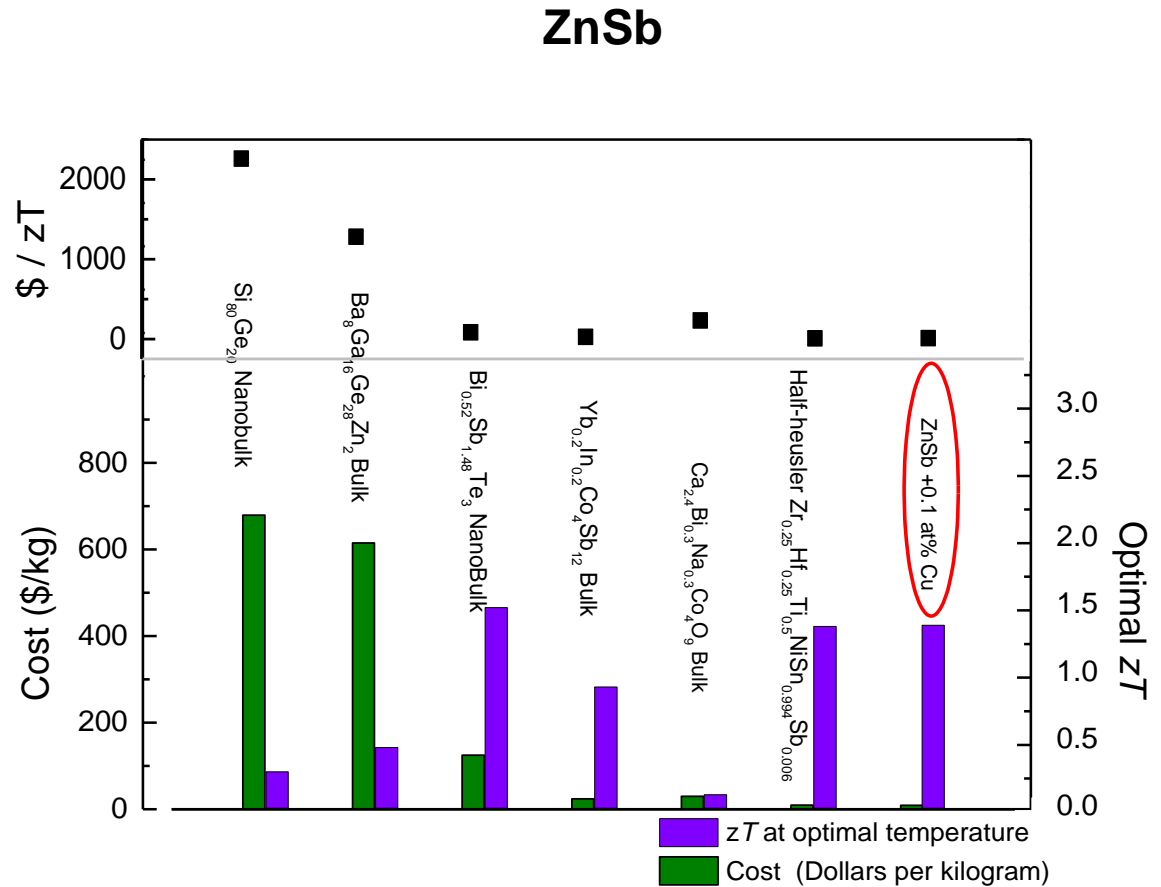
Choosing a material

Expectation	ZnSb
<ul style="list-style-type: none"> • Non-toxic • Abundant • High price-performance 	<ul style="list-style-type: none"> • Safe compound • Major industrial metals (Zn:13.7 MT; Sb 130 KT (2016)) • Estimated cost <10\$/kg
<ul style="list-style-type: none"> • Relevant operation temperature 	<ul style="list-style-type: none"> • 400K ~ 600K • Optimal temperature can be further tuned
<ul style="list-style-type: none"> • Narrow bandgap $6 \sim 10 k_B T$ (0.2~1 eV) 	<ul style="list-style-type: none"> • Band gap ~ 0.5 eV
<ul style="list-style-type: none"> • “Complex” structures 	<ul style="list-style-type: none"> • Deformed Zinc-blende <div style="display: flex; align-items: center; justify-content: center; margin-top: 10px;">  <div style="margin-left: 10px;">ZnSb</div>  <div style="margin-left: 10px;">ZnS</div> </div>



Choosing a material

Expectation
<ul style="list-style-type: none"> • Non-toxic • Abundant • High price-performance
<ul style="list-style-type: none"> • Relevant operation temperature
<ul style="list-style-type: none"> • Narrow bandgap $6 \sim 10 k_B T$ (0.2~1 eV)
<ul style="list-style-type: none"> • “Complex” structures



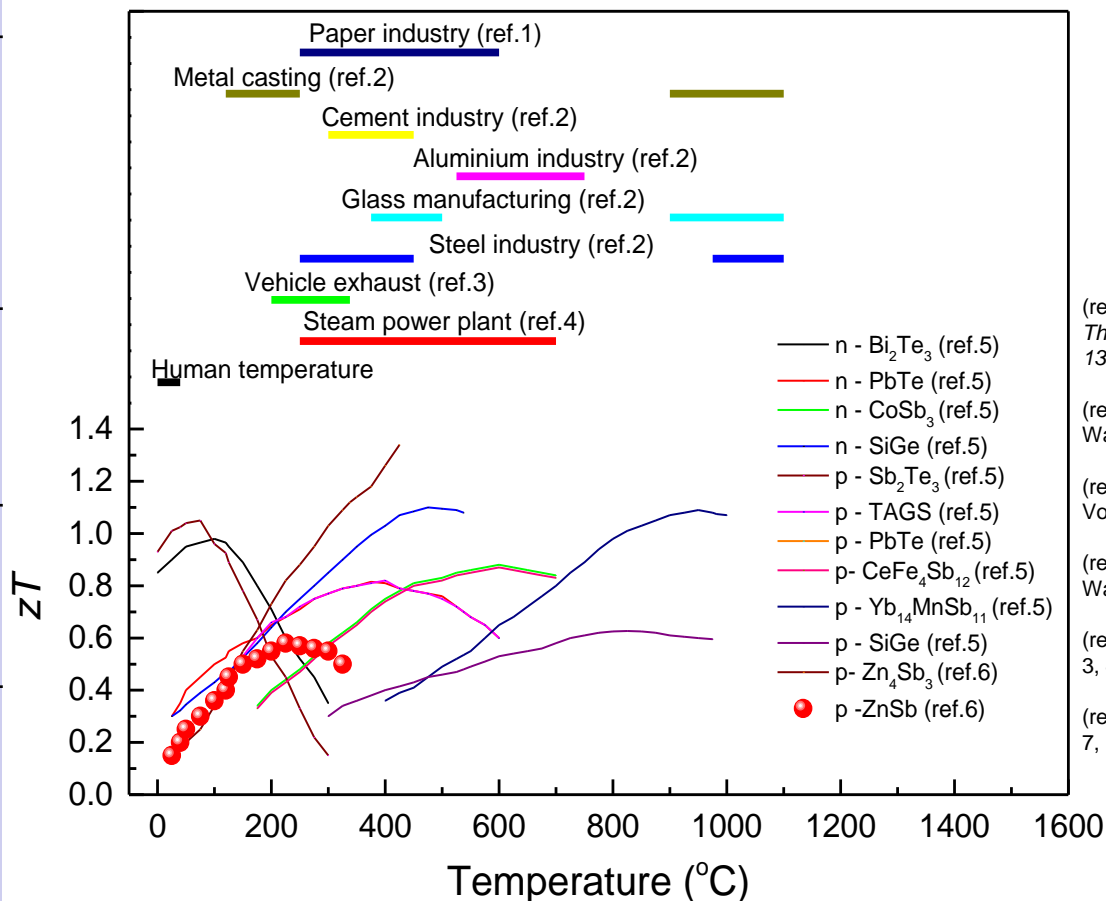
A. LeBlanc, et al. *Renewable and Sustainable Energy Reviews* 32(2014) 313.-327
* USGS, *Energy minerals*, Mar, 2015, <http://minerals.usgs.gov/minerals/pubs/commodity>



Choosing a material

Expectation
<ul style="list-style-type: none"> • Non-toxic • Abundant • High price-performance
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<ul style="list-style-type: none"> • "Complex" structures

ZnSb



(ref.1)
Thermal Science
13, 165-174 (2009)

(ref.2) EERE
Waste heat recovery, 2015

(ref.3) *J. Electron. Mater.*,
Vol. 44, No. 6, 2015

(ref.4) EPA,
Waste Heat to Power, 2012

(ref.5) *Nature Materials*
3, 458 - 463 (2004)

(ref.6) *Nature Materials*
7, 105 - 114 (2008)

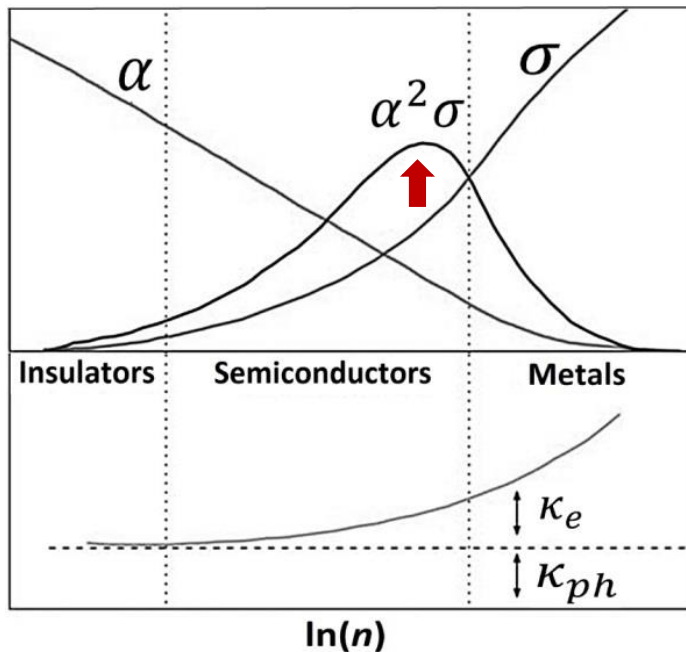


Doping defects

Given by $E_g = 0.53 \text{ eV}$ and $m^* = 0.42$ for ZnSb,

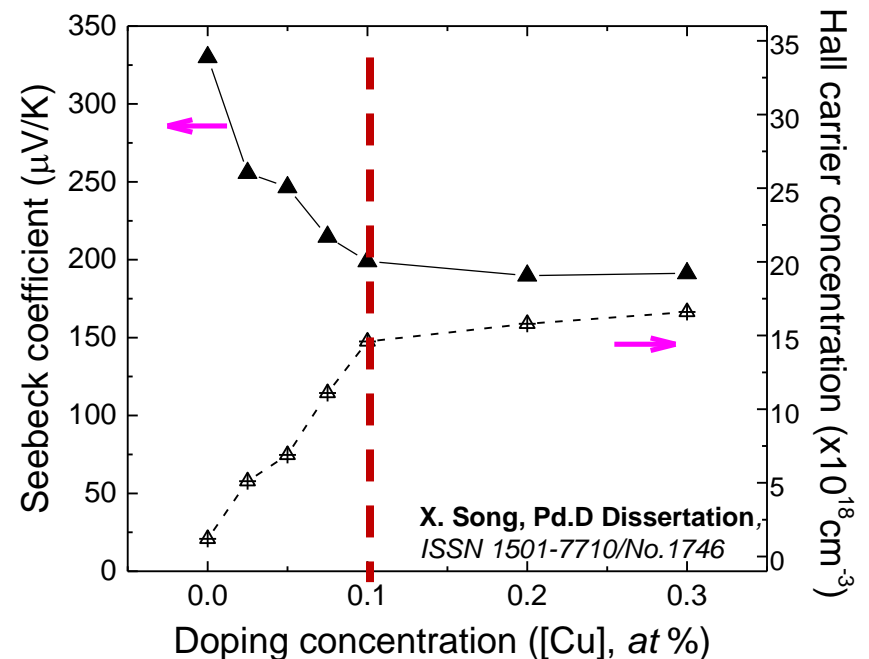
The theoretical intrinsic charge carrier concentrations at:

- Room temperature $2.41 \times 10^{14} \text{ cm}^{-3}$;
- $300 \text{ }^\circ\text{C}$: $8.42 \times 10^{16} \text{ cm}^{-3}$

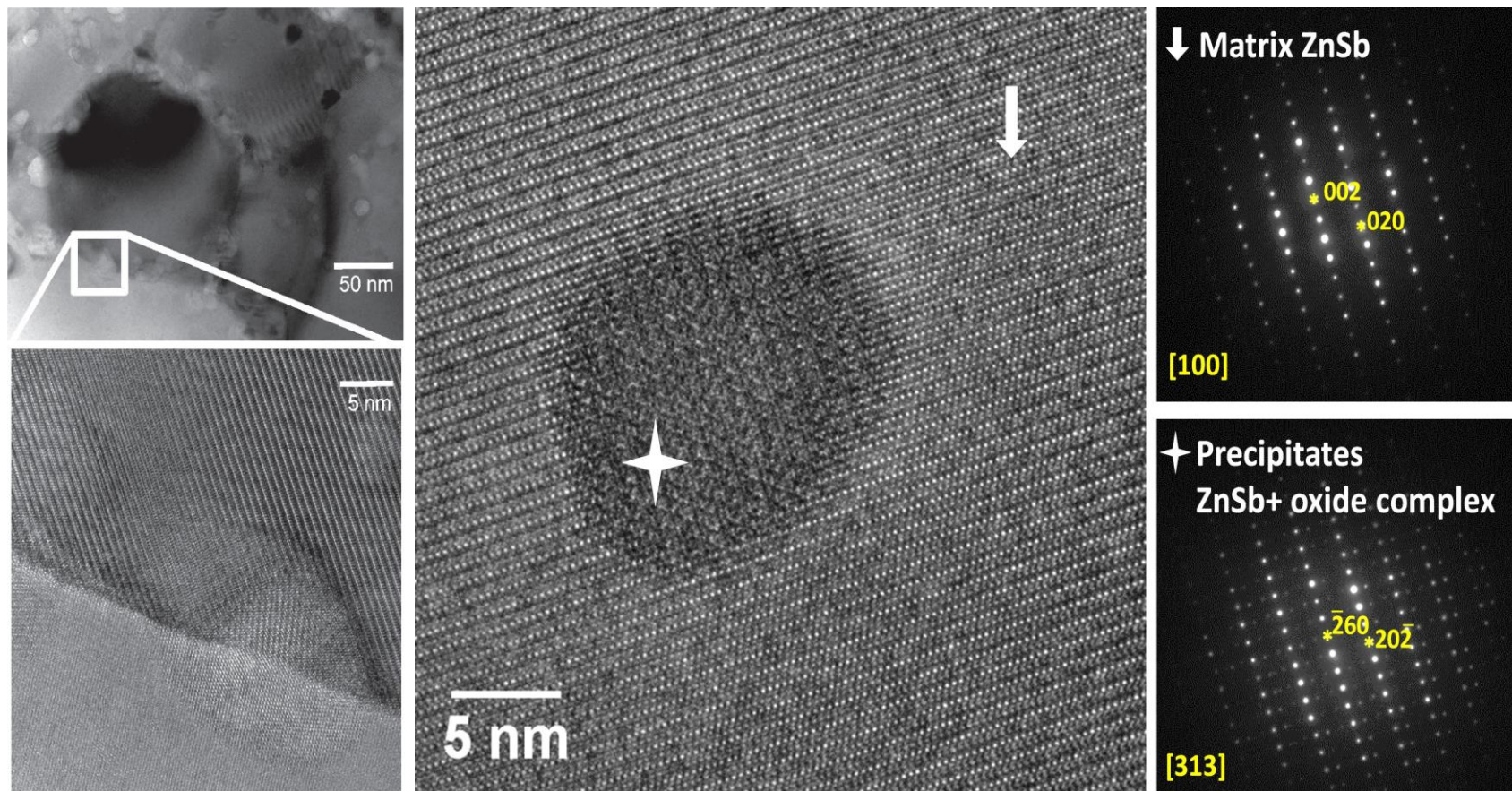


Power factor peaks at $2 \times 10^{19} \text{ cm}^{-3}$ for ZnSb

Acceptors:	Donor: (Tricky and unsuccessful)
I_{Zn} : Cu, Ag	III_{Zn} : In, Al, Ga
IV_{Sb} : Sn	VI_{Sb} : Te



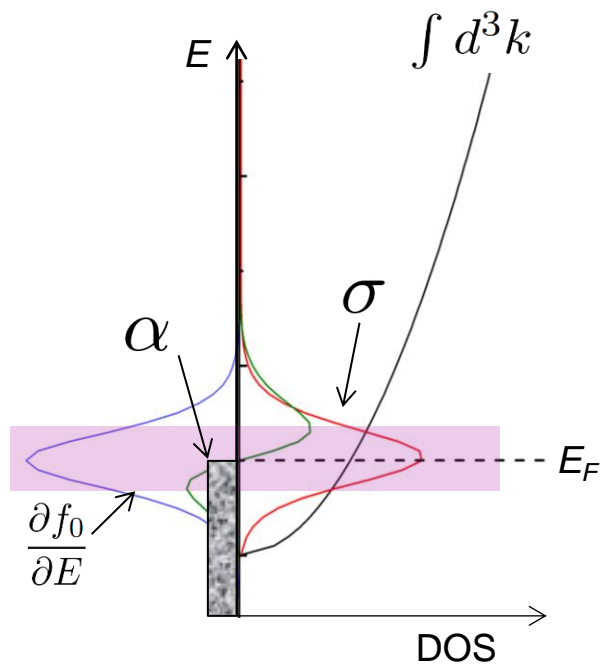
Nano-oxide “defects”



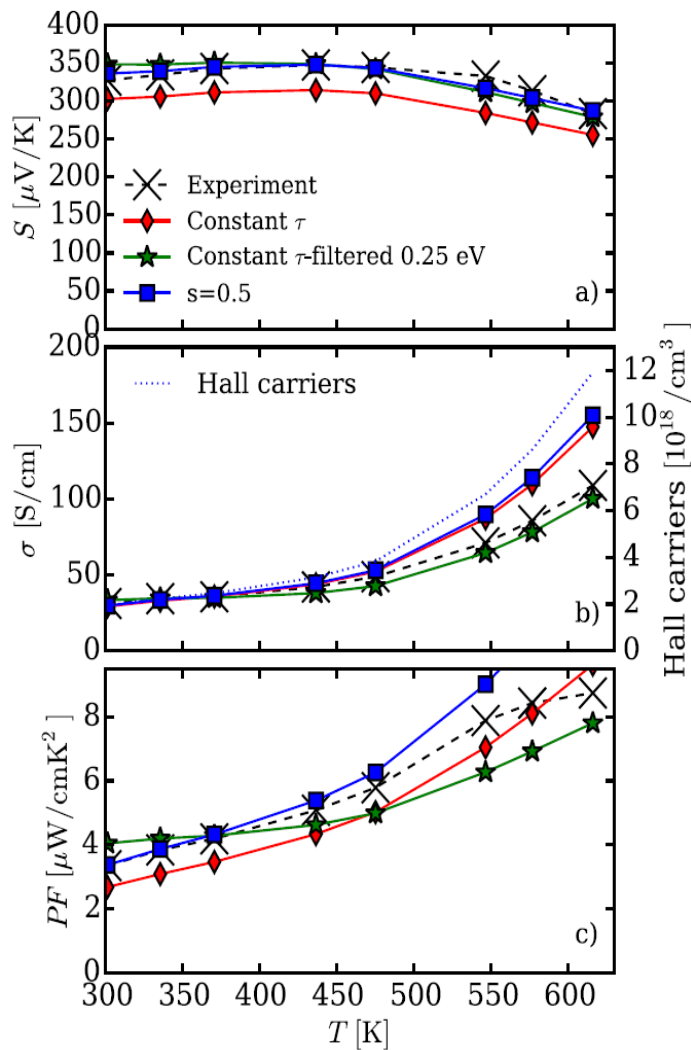
Transmission electron microscopy image of nano oxide particles in ZnSb
K. Berland, *et al*, *J. App. Phys.* 119, 125103 (2016)



Nano-oxide “defects”



Concept of energy filtering



K. Berland, et al, *J. App. Phys.* 119, 125103 (2016)



Intrinsic defects

The theoretical intrinsic charge carrier concentrations at room temperature

$$2.41 \times 10^{14} \text{ cm}^{-3} (E_g = 0.53 \text{ eV and } m^* = 0.42);$$

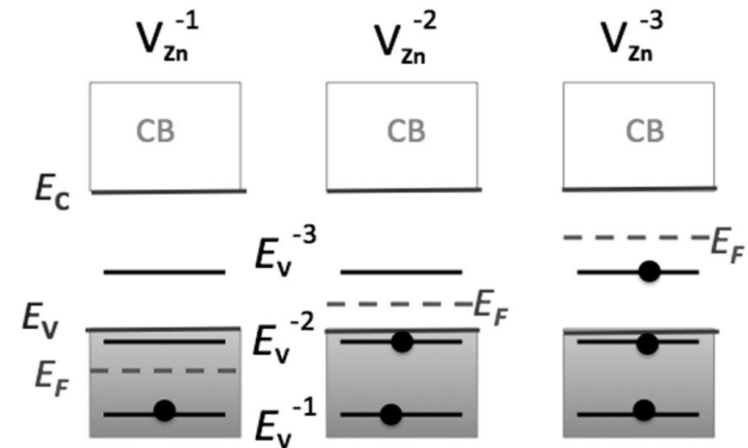
However, most of reported on single crystal undoped ZnSb at RT : $1 - 2 \times 10^{16} \text{ cm}^{-3}$;
polycrystalline undoped ZnSb at RT : $\sim 10^{18} \text{ cm}^{-3}$.

V_{Zn} V_{Sb} Zn_{Sb} Sb_{Zn} Zn_i Sb_i

Defect Formation Energies in eV for $\mu_e = \Delta\mu = 0$

q_d	V_{Zn}	V_{Sb}	Sb_{Zn}	Zn_{Sb}
-3	0.51	2.47	1.83	1.33
-2	0.30	2.05	1.50	1.26
-1	0.32	1.83	1.33	1.28
0	0.52	1.79	1.34	1.41
+1	0.88	1.70	1.39	1.67

L. Bjerg, et al. Chem. Mater. 2012, 24, 2111–2116

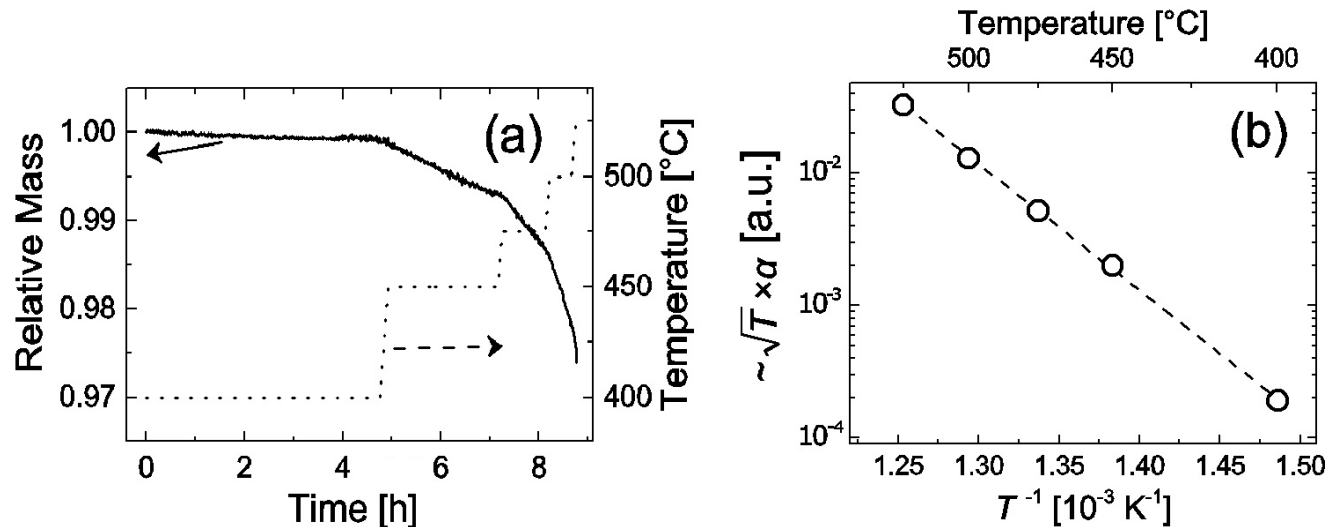
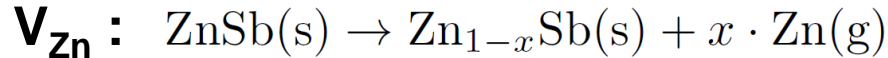


Occupancy of localized states associated with different charged V_{Zn} in ZnSb.

X. Song & T. G. Finstad, Review of Research on the Thermoelectric Material ZnSb, in Thermoelectrics for Power Generation - A Look at Trends in the Technology, Intech, (2016), Chap. 6



Intrinsic defects

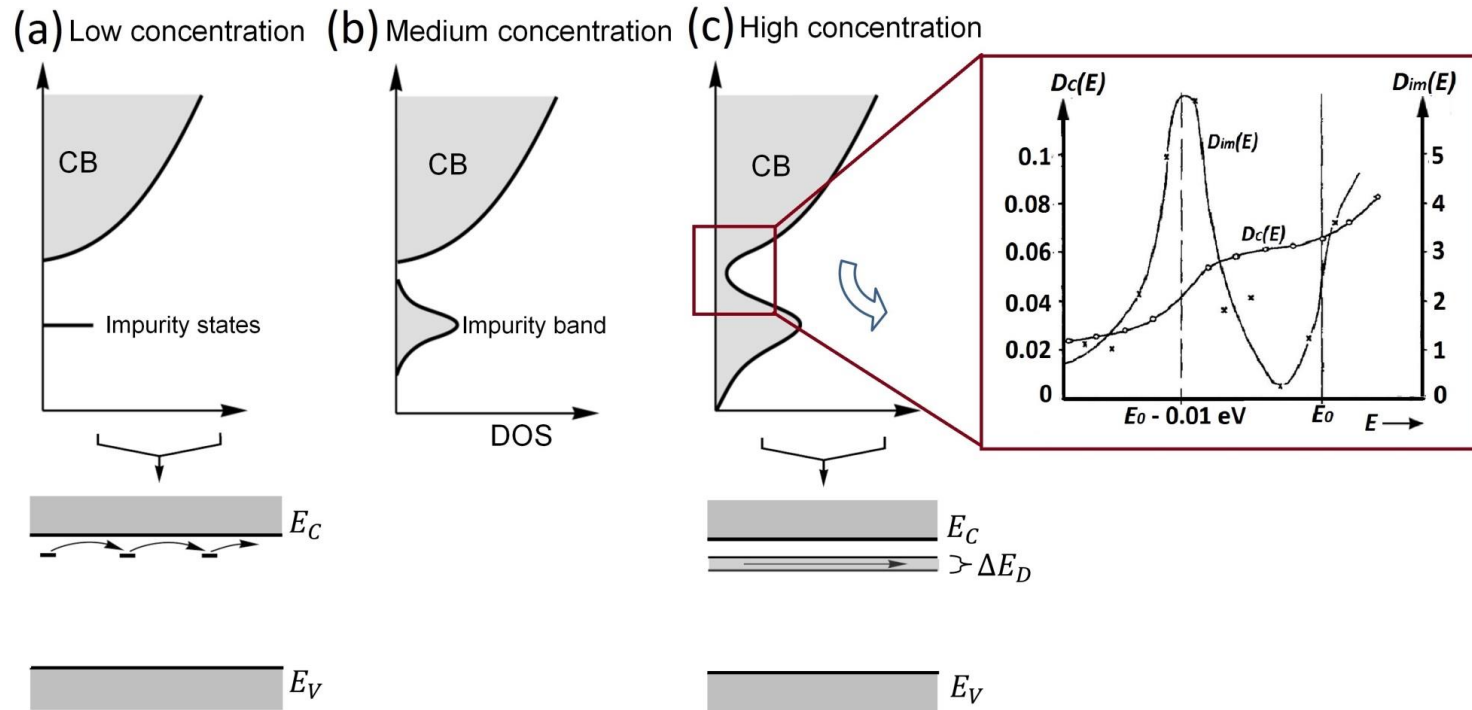


Influence of annealing at 400 °C for 30 minutes on the electronic properties of ZnSb at room temperature. The carrier concentration increases significantly for both samples, leading to an increased conductivity.

Sample	Annealing	p [$10^{18} cm^{-3}$]	σ [S cm]	μ [$cm^2 V^{-1} s^{-1}$]
# 1	Before	7.95	202	159
	After	14.1	262	116
# 2	Before	3.57	134	235
	After	10.2	360	220



Impurity band conduction

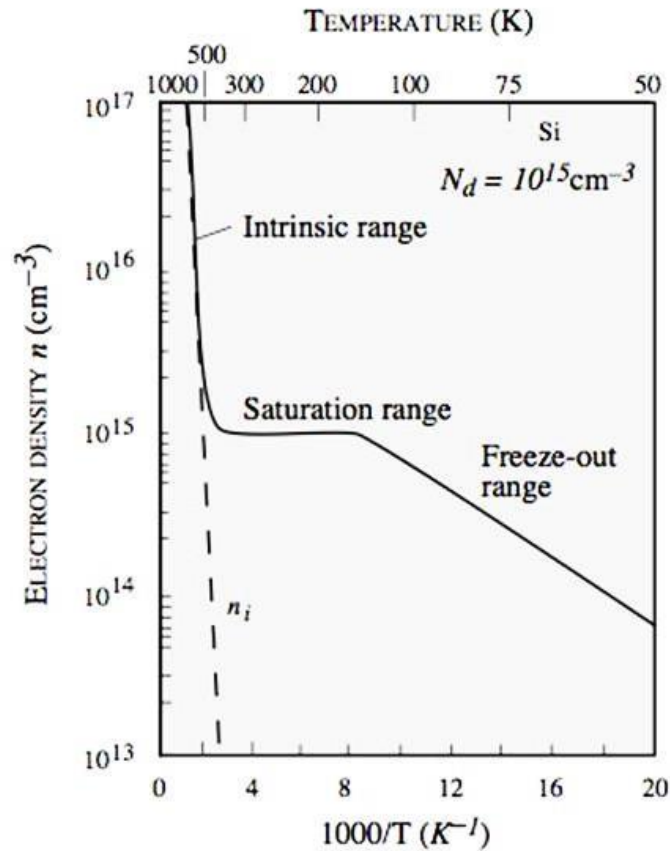


Mott transition criterion: $N_i \sim 6 \times 10^{17} \text{ cm}^{-3}$ for ZnSb

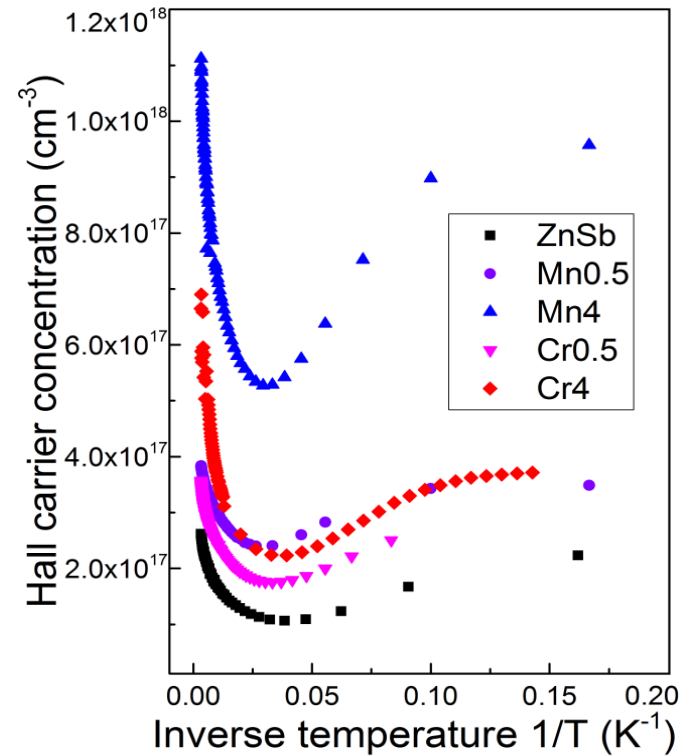


Impurity band conduction

Classical semiconductor behavior in Si



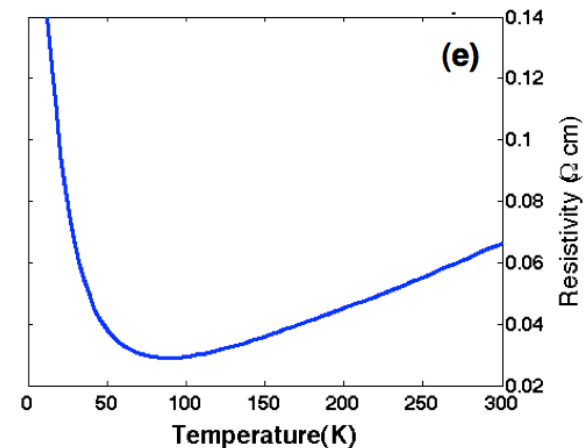
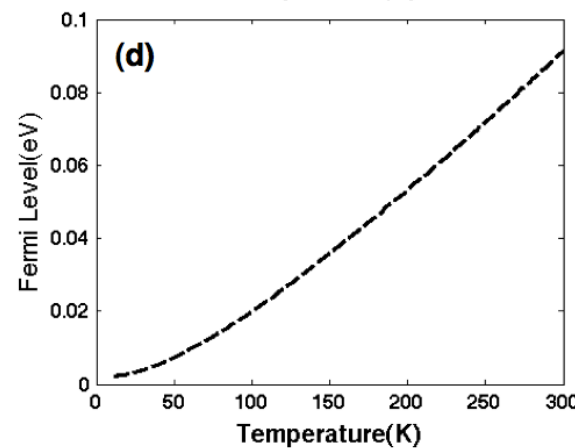
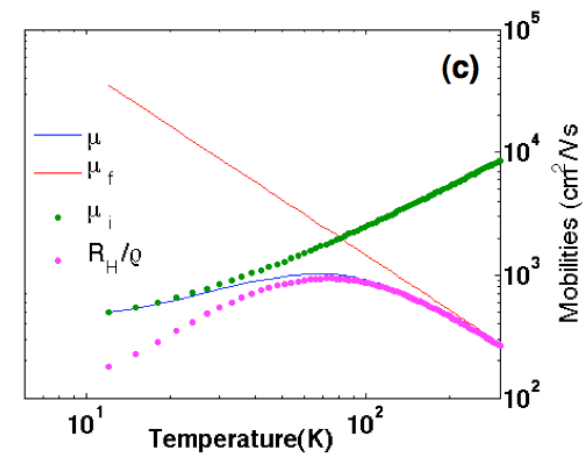
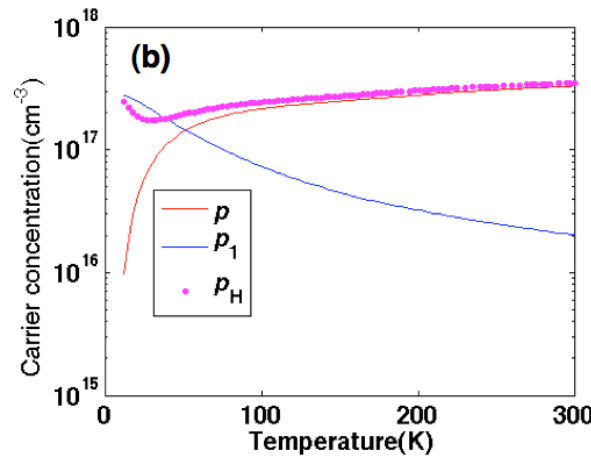
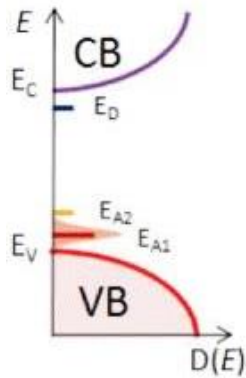
In ZnSb





Impurity band conduction

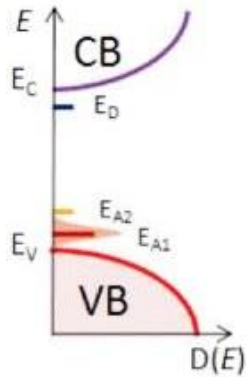
- Two types of carriers with different effective mass and mobilities
- Common in thermoelectrics, Also see in ZnSe, PdTe and BiTe, Ge, etc.





Impurity band conduction

Modelling of two carriers (VB carriers and IB carriers) conduction



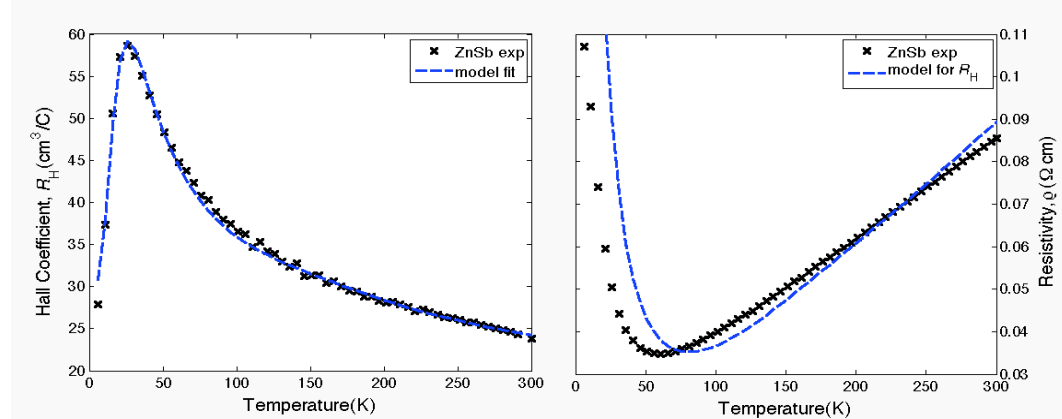
Variables: $E_{A1}, E_{A2}, N_{A1}, N_{A2}, N_D$

$$R_H = \frac{p_v \mu_v^2 + p_{im} \mu_{im}^2}{q(p_v \mu_v + p_{im} \mu_{im})^2}$$

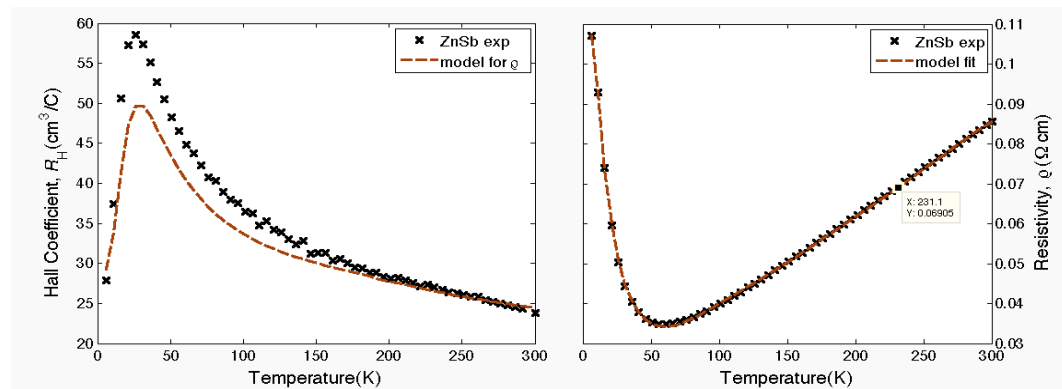
$$\rho = \frac{1}{q(p_v \mu_v + p_{im} \mu_{im})}$$

p_v, μ_v
Carrier concentration, mobility in VB
 p_{im}, μ_{im}
Carrier concentration, mobility in IB

Fitting Hall coefficient, simulating resistivity



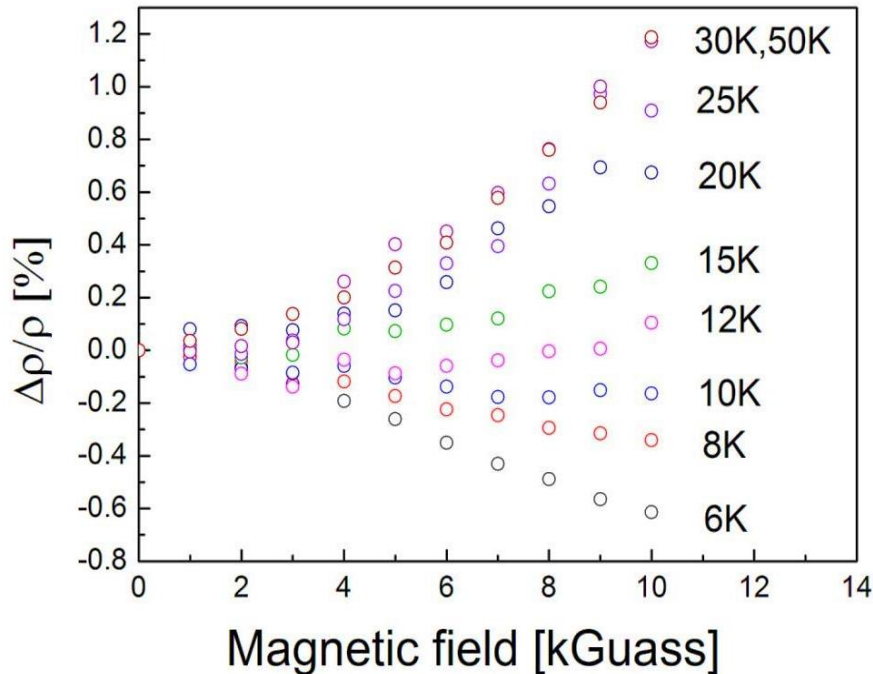
Fitting resistivity, simulating Hall coefficient





Impurity band conduction

- Magnetoresistance experimental prove of two carriers conduction: (VB carriers and IB carriers)
- Bending down indicates contribution from IB



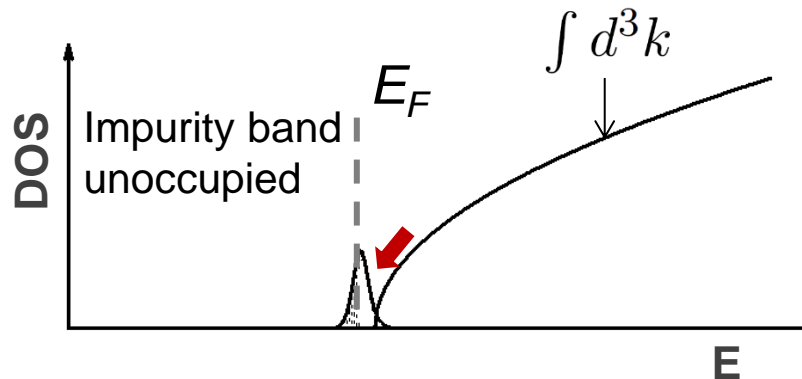
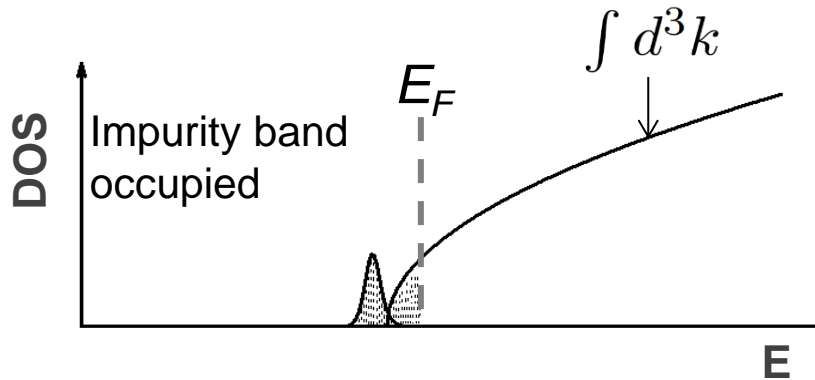
- How does it impact to thermoelectric performance?
- Can we manipulate impurity band?



Impurity band conduction

How does it impact to thermoelectric performance?

- Change shape of DOS



Seebeck coefficient

$$\alpha = \frac{1}{qT} \frac{\int \tau v^2 (E - E_F) \left(-\frac{\partial f_0}{\partial E}\right) d^3k}{\int \tau v^2 \left(-\frac{\partial f_0}{\partial E}\right) d^3k}$$

Electrical conductivity

$$\sigma = \frac{q^2}{4\pi^3} \int \tau v^2 \left(-\frac{\partial f_0}{\partial E}\right) d^3k$$

$$zT = \frac{\alpha^2 \sigma}{\kappa_e + \kappa_{ph}} T$$

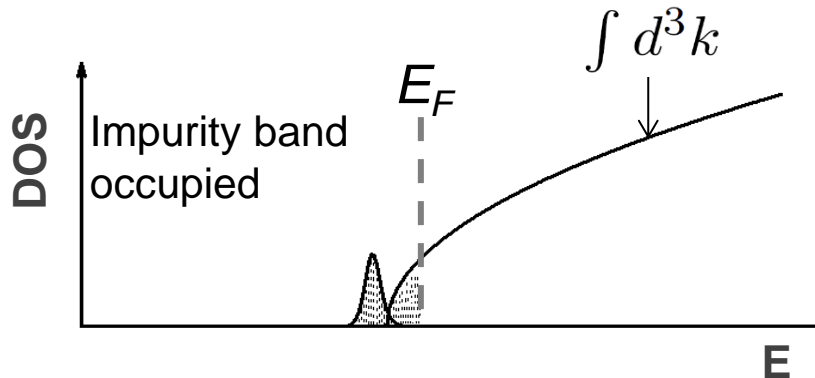




Impurity band conduction

How does it impact to thermoelectric performance?

- Change shape of DOS
- Change position of Fermi level -- Seebeck coefficient



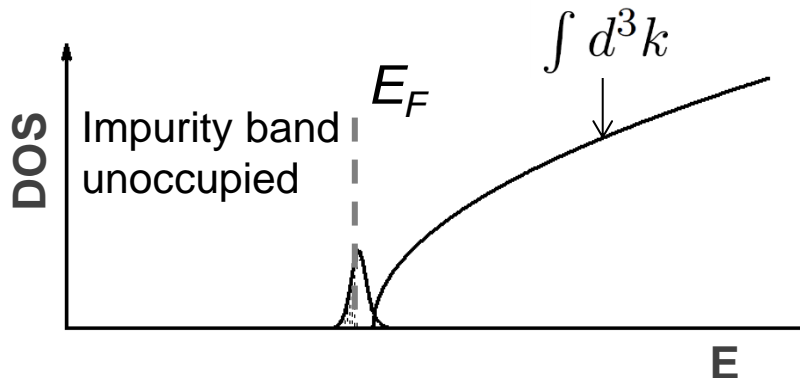
Seebeck coefficient

$$\alpha = \frac{1}{qT} \frac{\int \tau v^2 (E - E_F) \left(-\frac{\partial f_0}{\partial E}\right) d^3k}{\int \tau v^2 \left(-\frac{\partial f_0}{\partial E}\right) d^3k}$$

Position of Fermi level depends on

- Mass action law
- Charge neutrality

Fermi level goes between IB and VB at equilibrium

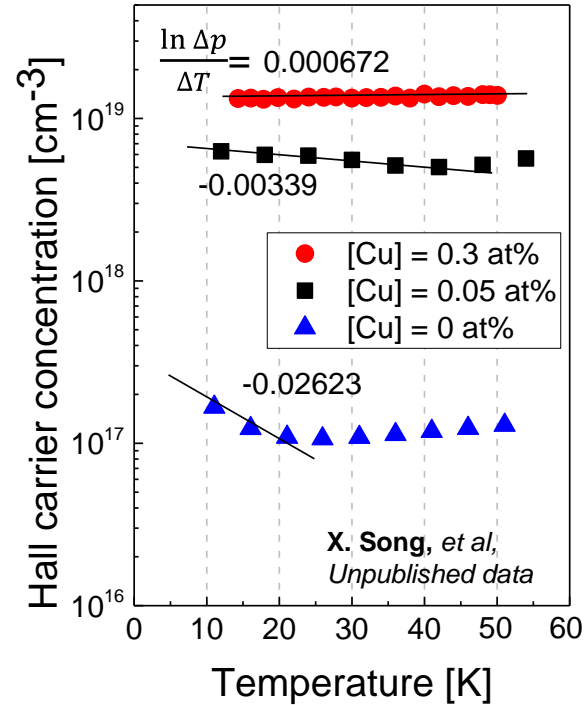
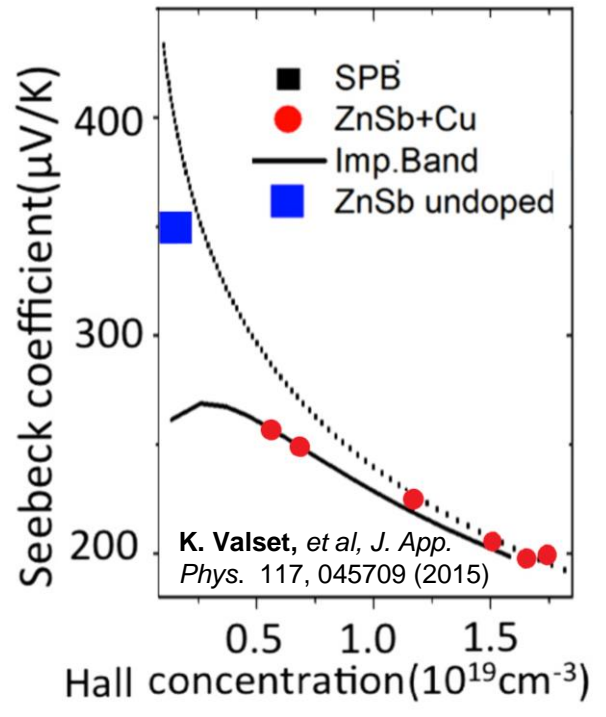




Impurity band conduction



Can we manipulate impurity band?



$$\alpha = \frac{\alpha_{VB} \cdot p_{VB} + \alpha_{IB} \cdot p_{IB}}{p_{VB} + p_{IB}}$$

Impurity band at 0.065eV above VB with concentration $2.5 \times 10^{18}\text{cm}^{-3}$

- Impurity band conduction in low concentration samples
- High concentration suppress impurity band conduction

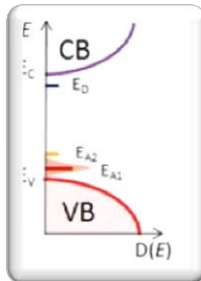
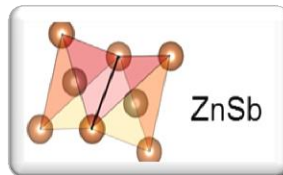


Summary

$$zT = \frac{\alpha^2 \sigma}{\kappa_e + \kappa_{ph}} T$$

Power factor

Thermal conductivity from electrons and phonons



• What matters to thermoelectrics

- Maximum zT at a particular carrier concentration
- Optimising power factor
- Reducing thermal transport

• ZnSb and its defects

- Decent thermoelectric material around 300 °C
- Doping defects – increase carrier concentration
- Nano-oxide – energy filtering
- Zn vacancies defects – intrinsic charge carriers

• Impurity band conduction and its impact on thermoelectric performance

- Elevated carrier concentration at cryo-temperature
- Two type of carriers with different effective mass and mobility
- Change DOS shape and Fermi level -- power factor